

# Ultrarelativistic boost of a black hole in the magnetic universe of Levi-Civita–Bertotti–Robinson (LCBR)\*

Marcello Ortaggio

Institute of Mathematics of the Czech Academy of Sciences

Wojanów – September 24th, 2018

---

\* Joint work with M. Astorino, Phys. Rev. D 97, 104052 (2018)

# Contents

- 1 The AS  $pp$ -wave
- 2 The Alekseev-Garcia solution
- 3 Ultrarelativistic limit of the AG solution

# The AS $pp$ -wave

Aichelburg-Sexl'70,'71:

- 1 take the Schwarzschild metric
- 2 *boost* it to the speed  $v$
- 3 take the limit  $v \rightarrow 1$  with  $m = p\sqrt{1-v^2}$   
use the identity  $\lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} f\left(\frac{x}{\epsilon}\right) = \delta(x) \int_{-\infty}^{\infty} f(t) dt$

gives

# The AS $pp$ -wave

Aichelburg-Sexl'70,'71:

- ① take the Schwarzschild metric
- ② *boost* it to the speed  $v$
- ③ take the limit  $v \rightarrow 1$  with  $m = p\sqrt{1-v^2}$   
use the identity  $\lim_{\epsilon \rightarrow 0} \frac{1}{\epsilon} f\left(\frac{x}{\epsilon}\right) = \delta(x) \int_{-\infty}^{\infty} f(t) dt$

gives

$$ds^2 = 2dudv + dx^2 + dy^2 - 4\sqrt{2}p \ln(x^2 + y^2) \delta(u) du^2,$$

with  $T_{uu} \propto p \delta(x) \delta(y) \delta(u)$ .

- it's a  $pp$ -wave ( $\partial_v$ )

- it's a  $pp$ -wave ( $\partial_v$ )
- all curvature invariants vanish

- it's a  $pp$ -wave ( $\partial_v$ )
- all curvature invariants vanish
- immune to higher-order corrections  $L = L(\text{Riem}, \nabla \text{Riem}, \dots)$   
[Amati-Klimčík'89, Horowitz-Steif'90]

- it's a  $pp$ -wave ( $\partial_v$ )
- all curvature invariants vanish
- immune to higher-order corrections  $L = L(\text{Riem}, \nabla \text{Riem}, \dots)$   
[Amati-Klimčík'89, Horowitz-Steif'90]
- used as “incoming state”:  
high-speed black hole encounters [Curtis'78, D'Eath'78, ...]  
semi-classical approaches to Planckian scattering and BH  
production “in the lab” [’t Hooft'87, Eardley-Giddings'02, ...]



Other constructions of impulsive waves:

- Penrose's “cut-and-paste” method
- limits of sandwich waves.

Other constructions of impulsive waves:

- Penrose's “cut-and-paste” method
- limits of sandwich waves.

Generalizations to various backgrounds . . . (and dimensions)

Backgrounds	cut-and-paste	sandwich waves	AS boost
Minkowski	Penrose'68,'72 Dray-'t Hooft'85 (*) Sfetsos'95	Brinkmann'25 Peres'60, Kundt'61 Bondi <i>et al.</i> '62,... Bonnor'69	Aichelburg-Sexl'70 Loustó-Sanchez'89,'90 Balasin-Nachbagauer'95 Barrabés-Hogan'01,...
(A)dS	Hotta-Tanaka'93 Podolský-Griffiths'99 Horowitz-Itzhaki'99	García D.-Plebański'81 Ozsváth <i>et al.</i> '83 Siklos'85	Hotta-Tanaka'93
Bonnor-Melvin	(cf. *)	Garfinkle-Melvin'92	M.O.'04
LCBR, Nariai	(cf. *)	García D.-Alvarez C.'84 Khlebnikov'86 Lewandowski'92 Podolský-M.O.'03	??

# The Alekseev-Garcia solution

$$ds^2 = -e^{2\psi} \cosh^2 \frac{z}{b} dt^2 + e^{2\gamma} (dz^2 + d\rho^2) + e^{-2\psi} b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$

coordinates  $(t, z, \rho, \phi)$ ,  $\rho \in [0, \pi b]$ ,  $\phi \in [0, 2\pi]$ ; 3 parameters  $m, b, l$

# The Alekseev-Garcia solution

$$ds^2 = -e^{2\psi} \cosh^2 \frac{z}{b} dt^2 + e^{2\gamma} (dz^2 + d\rho^2) + e^{-2\psi} b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$

coordinates  $(t, z, \rho, \phi)$ ,  $\rho \in [0, \pi b]$ ,  $\phi \in [0, 2\pi]$ ; 3 parameters  $m, b, l$

$$e^{2\psi} = \frac{(R_+ + R_- - 2m \cos \frac{\rho}{b})^2}{(R_+ + R_-)^2 - 4m^2},$$

$$e^{2\gamma} = \frac{(R_+ + R_- - 2m \cos \frac{\rho}{b})^2}{4R_+ R_-} \left[ \frac{R_+ - b \sinh \frac{z}{b} + (l + m) \cos \frac{\rho}{b}}{R_- - b \sinh \frac{z}{b} + (l - m) \cos \frac{\rho}{b}} \right]^2,$$

$$R_{\pm}^2 = \left( l \pm m - b \sinh \frac{z}{b} \cos \frac{\rho}{b} \right)^2 + b^2 \cosh^2 \frac{z}{b} \sin^2 \frac{\rho}{b},$$

$$A_{\phi} = -b \frac{R_+ + R_- + 2m}{R_+ + R_- - 2m \cos \frac{\rho}{b}} \left( 1 - \cos \frac{\rho}{b} \right).$$

# The Alekseev-Garcia solution

$$ds^2 = -e^{2\psi} \cosh^2 \frac{z}{b} dt^2 + e^{2\gamma} (dz^2 + d\rho^2) + e^{-2\psi} b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$

coordinates  $(t, z, \rho, \phi)$ ,  $\rho \in [0, \pi b]$ ,  $\phi \in [0, 2\pi]$ ; 3 parameters  $m, b, l$

$$e^{2\psi} = \frac{(R_+ + R_- - 2m \cos \frac{\rho}{b})^2}{(R_+ + R_-)^2 - 4m^2},$$

$$e^{2\gamma} = \frac{(R_+ + R_- - 2m \cos \frac{\rho}{b})^2}{4R_+ R_-} \left[ \frac{R_+ - b \sinh \frac{z}{b} + (l + m) \cos \frac{\rho}{b}}{R_- - b \sinh \frac{z}{b} + (l - m) \cos \frac{\rho}{b}} \right]^2,$$

$$R_{\pm}^2 = \left( l \pm m - b \sinh \frac{z}{b} \cos \frac{\rho}{b} \right)^2 + b^2 \cosh^2 \frac{z}{b} \sin^2 \frac{\rho}{b},$$

$$A_{\phi} = -b \frac{R_+ + R_- + 2m}{R_+ + R_- - 2m \cos \frac{\rho}{b}} \left( 1 - \cos \frac{\rho}{b} \right).$$

Horizon at:  $\rho = 0$ ,  $l - m \leq b \sinh \frac{z}{b} \leq l + m$ .

$$m = 0:$$

LCBR solution  $\text{AdS}_2 \times S_2$  [Levi-Civita'17, Bertotti'59, Robinson'59]

$$ds^2 = -\cosh^2 \frac{z}{b} dt^2 + dz^2 + d\rho^2 + b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$
$$A_\phi = -b \left( 1 - \cos \frac{\rho}{b} \right).$$

$$m = 0:$$

LCBR solution  $\text{AdS}_2 \times S_2$  [Levi-Civita'17, Bertotti'59, Robinson'59]

$$ds^2 = -\cosh^2 \frac{z}{b} dt^2 + dz^2 + d\rho^2 + b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$
$$A_\phi = -b \left( 1 - \cos \frac{\rho}{b} \right).$$

Same metric in the asymptotic region  $z \rightarrow -\infty$ .



$$m = 0:$$

LCBR solution  $\text{AdS}_2 \times S_2$  [Levi-Civita'17, Bertotti'59, Robinson'59]

$$ds^2 = -\cosh^2 \frac{z}{b} dt^2 + dz^2 + d\rho^2 + b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$

$$A_\phi = -b \left( 1 - \cos \frac{\rho}{b} \right).$$

Same metric in the asymptotic region  $z \rightarrow -\infty$ .

$$b \rightarrow +\infty: \text{ (exterior) Schwarzschild in Weyl coordinates}$$

$$ds^2 = -e^{2\psi} dt^2 + e^{2\gamma} (d\tilde{z}^2 + d\rho^2) + e^{-2\psi} \rho^2 d\phi^2,$$

$$e^{2\psi} = \frac{R_+ + R_- - 2m}{R_+ + R_- + 2m}, \quad e^{2\gamma} = \frac{(R_+ + R_- + 2m)^2}{4R_+ R_-},$$

$$R_\pm^2 = (\pm m - \tilde{z})^2 + \rho^2 \quad (\tilde{z} = z - l).$$

# Ultrarelativistic limit of the AG solution

Idea: decompose

$$ds_{AG}^2 = (ds_{AdS_2}^2 + ds_{S_2}^2) + \Delta \quad (\Delta = \dots)$$

# Ultrarelativistic limit of the AG solution

Idea: decompose

$$ds_{AG}^2 = (ds_{AdS_2}^2 + ds_{S_2}^2) + \Delta \quad (\Delta = \dots)$$

in embedding coordinates

$$ds_{AdS_2}^2 = 2dUdV - dZ_2^2, \quad -2UV + Z_2^2 = b^2$$

# Ultrarelativistic limit of the AG solution

Idea: decompose

$$ds_{AG}^2 = (ds_{AdS_2}^2 + ds_{S_2}^2) + \Delta \quad (\Delta = \dots)$$

in embedding coordinates

$$ds_{AdS_2}^2 = 2dUdV - dZ_2^2, \quad -2UV + Z_2^2 = b^2$$

the AS boost is:

$$U \mapsto \epsilon^{-1}U, \quad V \mapsto \epsilon V,$$

$$m \mapsto \epsilon p.$$

# Ultrarelativistic limit of the AG solution

Idea: decompose

$$ds_{AG}^2 = (ds_{AdS_2}^2 + ds_{S_2}^2) + \Delta \quad (\Delta = \dots)$$

in embedding coordinates

$$ds_{AdS_2}^2 = 2dUdV - dZ_2^2, \quad -2UV + Z_2^2 = b^2$$

the AS boost is:

$$U \mapsto \epsilon^{-1}U, \quad V \mapsto \epsilon V,$$

$$m \mapsto \epsilon p.$$

Study  $\lim_{\epsilon \rightarrow 0} \Delta \dots$  (the rest is invariant).

# Ultrarelativistic limit of the AG solution

Idea: decompose

$$ds_{AG}^2 = (ds_{AdS_2}^2 + ds_{S_2}^2) + \Delta \quad (\Delta = \dots)$$

in embedding coordinates

$$ds_{AdS_2}^2 = 2dUdV - dZ_2^2, \quad -2UV + Z_2^2 = b^2$$

the AS boost is:

$$U \mapsto \epsilon^{-1}U, \quad V \mapsto \epsilon V,$$

$$m \mapsto \epsilon p.$$

Study  $\lim_{\epsilon \rightarrow 0} \Delta \dots$  (the rest is invariant).

(Note  $U = 0 \Rightarrow Z_2 = \pm b$ .)

Eventually:

$$ds^2 = \frac{2dudv + 2H\delta(u)du^2}{\Omega^2} + d\rho^2 + b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$
$$A_\phi = -b \left(1 - \cos \frac{\rho}{b}\right),$$

Eventually:

$$ds^2 = \frac{2dudv + 2H\delta(u)du^2}{\Omega^2} + d\rho^2 + b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$

$$A_\phi = -b \left(1 - \cos \frac{\rho}{b}\right),$$

$$\Omega = 1 - \frac{1}{2}b^{-2}uv, \quad H = \sqrt{2}p \frac{b^2}{b^2 + l^2} \ln \frac{1 + \cos \frac{\rho}{b}}{1 - \cos \frac{\rho}{b}},$$

with  $T_{uu} \propto p\delta\left(1 + \cos \frac{\rho}{b}\right)\delta\left(1 - \cos \frac{\rho}{b}\right)\delta(u)$ .



Eventually:

$$ds^2 = \frac{2dudv + 2H\delta(u)du^2}{\Omega^2} + d\rho^2 + b^2 \sin^2 \frac{\rho}{b} d\phi^2,$$

$$A_\phi = -b \left(1 - \cos \frac{\rho}{b}\right),$$

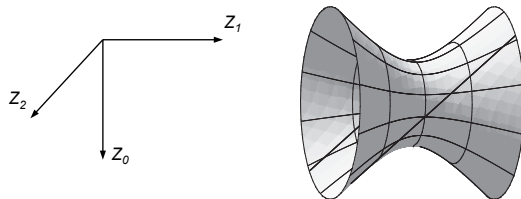
$$\Omega = 1 - \frac{1}{2}b^{-2}uv, \quad H = \sqrt{2}p \frac{b^2}{b^2 + l^2} \ln \frac{1 + \cos \frac{\rho}{b}}{1 - \cos \frac{\rho}{b}},$$

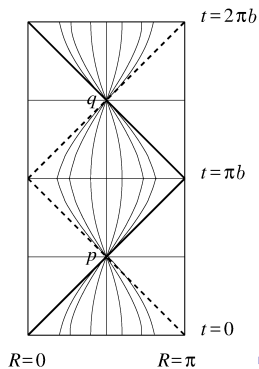
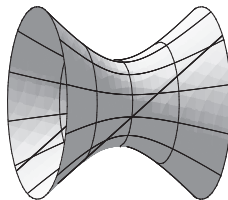
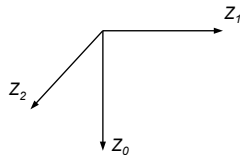
with  $T_{uu} \propto p\delta\left(1 + \cos \frac{\rho}{b}\right)\delta\left(1 - \cos \frac{\rho}{b}\right)\delta(u)$ .

LCBR coordinates given by:

$$b \cosh \frac{z}{b} \cos \frac{t}{b} = \frac{u - v}{\sqrt{2}\Omega}, \quad b \sinh \frac{z}{b} = \frac{u + v}{\sqrt{2}\Omega},$$

$$b \cosh \frac{z}{b} \sin \frac{t}{b} = (\pm)b \left(\frac{2}{\Omega} - 1\right).$$





Kundt form: under  $\tilde{v} = \frac{v}{\Omega}$ ,  $\zeta = \sqrt{2}be^{i\phi} \tan \frac{\rho}{2b}$

$$ds^2 = 2dud\tilde{v} + \left[ -b^{-2}\tilde{v}^2 + 2\tilde{H}(u, \zeta, \bar{\zeta}) \right] du^2 + \frac{2d\zeta d\bar{\zeta}}{\left(1 + \frac{1}{2}b^{-2}\zeta\bar{\zeta}\right)^2},$$

$$\tilde{H} = -\sqrt{2}p \frac{b^2}{b^2 + l^2} \delta(u) \ln \frac{\zeta\bar{\zeta}}{2b^2}.$$

Kundt form: under  $\tilde{v} = \frac{v}{\Omega}$ ,  $\zeta = \sqrt{2}be^{i\phi} \tan \frac{\rho}{2b}$

$$ds^2 = 2dud\tilde{v} + \left[ -b^{-2}\tilde{v}^2 + 2\tilde{H}(u, \zeta, \bar{\zeta}) \right] du^2 + \frac{2d\zeta d\bar{\zeta}}{\left(1 + \frac{1}{2}b^{-2}\zeta\bar{\zeta}\right)^2},$$

$$\tilde{H} = -\sqrt{2}p \frac{b^2}{b^2 + l^2} \delta(u) \ln \frac{\zeta\bar{\zeta}}{2b^2}.$$

- (degenerate) Kundt, recurrent  $\partial_{\tilde{v}}$
- CSI [Coley-Hervik-Pelavas'10]

Kundt form: under  $\tilde{v} = \frac{v}{\Omega}$ ,  $\zeta = \sqrt{2}be^{i\phi} \tan \frac{\rho}{2b}$

$$ds^2 = 2dud\tilde{v} + \left[ -b^{-2}\tilde{v}^2 + 2\tilde{H}(u, \zeta, \bar{\zeta}) \right] du^2 + \frac{2d\zeta d\bar{\zeta}}{\left(1 + \frac{1}{2}b^{-2}\zeta\bar{\zeta}\right)^2},$$

$$\tilde{H} = -\sqrt{2}p \frac{b^2}{b^2 + l^2} \delta(u) \ln \frac{\zeta\bar{\zeta}}{2b^2}.$$

- (degenerate) Kundt, recurrent  $\partial_{\tilde{v}}$
- CSI [Coley-Hervik-Pelavas'10]
- general family  $\tilde{H} = f(u, \zeta) + \bar{f}(u, \bar{\zeta})$  [García D.-Alvarez C.'84]
- Petrov type N, non-null Maxwell field

Kundt form: under  $\tilde{v} = \frac{v}{\Omega}$ ,  $\zeta = \sqrt{2}be^{i\phi} \tan \frac{\rho}{2b}$

$$ds^2 = 2dud\tilde{v} + \left[ -b^{-2}\tilde{v}^2 + 2\tilde{H}(u, \zeta, \bar{\zeta}) \right] du^2 + \frac{2d\zeta d\bar{\zeta}}{\left(1 + \frac{1}{2}b^{-2}\zeta\bar{\zeta}\right)^2},$$

$$\tilde{H} = -\sqrt{2}p \frac{b^2}{b^2 + l^2} \delta(u) \ln \frac{\zeta\bar{\zeta}}{2b^2}.$$

- (degenerate) Kundt, recurrent  $\partial_{\tilde{v}}$
- CSI [Coley-Hervik-Pelavas'10]
- general family  $\tilde{H} = f(u, \zeta) + \bar{f}(u, \bar{\zeta})$  [García D.-Alvarez C.'84]
- Petrov type N, non-null Maxwell field
- $pp$ -wave for  $b \rightarrow \infty$ .