

1st . General Introduction

- ✓ **Why multi-messengers (inc. GW)?**
- ✓ **Basics of GW Physics and Detection**
- ✓ **First detection of GW150914**

2nd . **Core-collapse supernova theory:** **how to solve “numerically”** **the space-time evolution of dying stars** **(40 min)**

3rd . **GW signatures from core-collapse supernovae:** what we can learn from future GW observation ? **(40 min)**

Gravitational Waves (GWs) from Stellar Collapse

(see reviews in Ott (2009), Fryer & New (2011), Kotake (2013), Kotake and Kuroda (2016) in "Handbook of Supernovae")

GW amplitude from the quadrupole formula

$$h_{ij} = \frac{2G}{c^4 R} \frac{\partial^2}{\partial t^2} \underline{Q_{ij}} \sim \frac{R_s}{R} \left(\frac{v}{c}\right)^2$$

Quadrupole moment

Typical values at the formation of Neutron Star (NS)

$$R_s = 3 \text{ km} \left(\frac{M}{M_\odot}\right) \quad v/c = 0.1 \quad \underline{R = 10 \text{ kpc}}$$

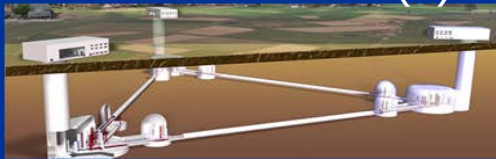
$$h \sim 10^{-20}$$

Good news ! (Future)

10 km long:

Einstein Telescope (ET)

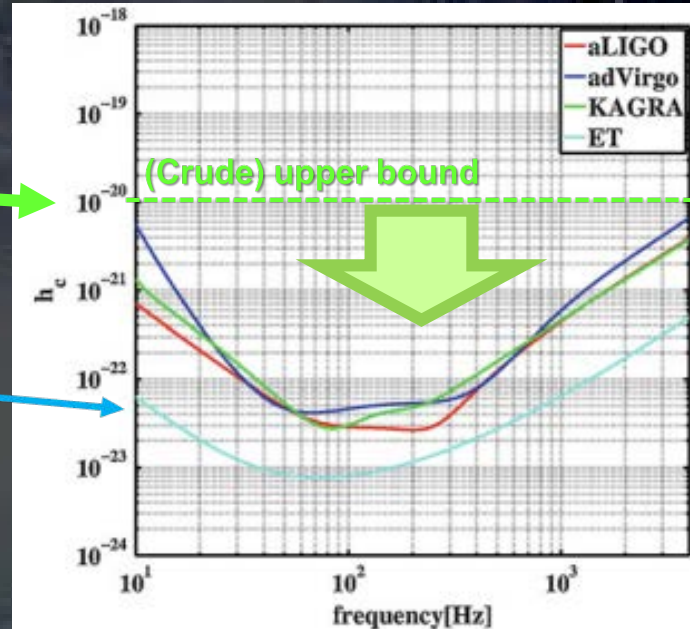
could start ~2025(?)



40 km long:

Cosmic Explore (CE)

could operate ~2035(?)



✓ CCSN event in our galaxy (several/century) is primary target !

More correctly...

$$h_{ij} = \epsilon \frac{R_s}{R} \left(\frac{v}{c}\right)^2$$

ϵ

: the degree of anisotropy.

If collapse proceeds spherically, $\epsilon = 0$ no GWs !

What makes the SN-dynamics deviate from spherical symmetry is essential for the GW emission mechanism !

Two candidate mechanisms of core-collapse supernovae

(Lecture by T. Foglizzo, reviews in Janka ('17), Müller ('16), Foglizzo+('15), Burrows('13), Kotake+ ('12))

Neutrino mechanism

Non- or slowing- rotating star
($\Omega_0 < \sim 0.1$ rad/s)

✓ Turbulent Convection and SASI

(e.g., Kazeroni, Guilet, Foglizzo, (2017))

✓ Precollapse Inhomogeneities/structures

(e.g., B.Mueller et al. ('17), Suwa & Mueller ('16))

✓ Novel microphysics: Bollig+('17), Fischer+('18)

Main players

MHD mechanism

Rapidly rotation with strong B
($\Omega_0 > \sim \pi$ rad/s, $B_0 > \sim 10^{11}$ G)

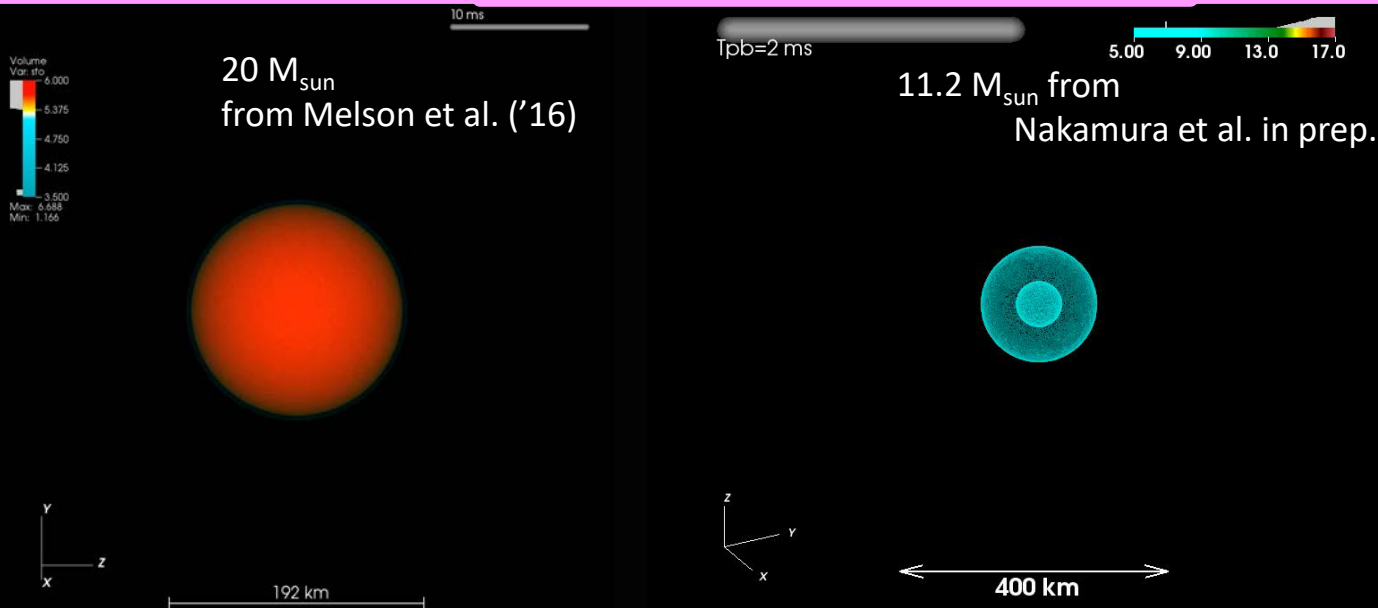
✓ Field winding and the MRI

(e.g., Obergaulinger & Aloy (2017), Rembiasz et al. (2016), Moesta et al. (2016), Masada + (2015))

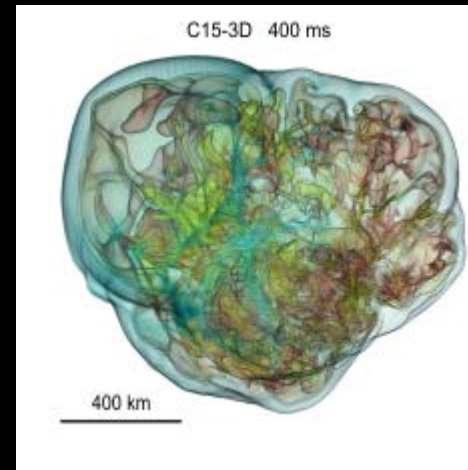
✓ Non-Axisymmetric instabilities

(e.g., Takiwaki, et al. (2016), Summa et al. (2017))

$\sim < 1\%$ (Woosley & Heger (07), ApJ):
(hypothetical link to magnetar, collapsar)



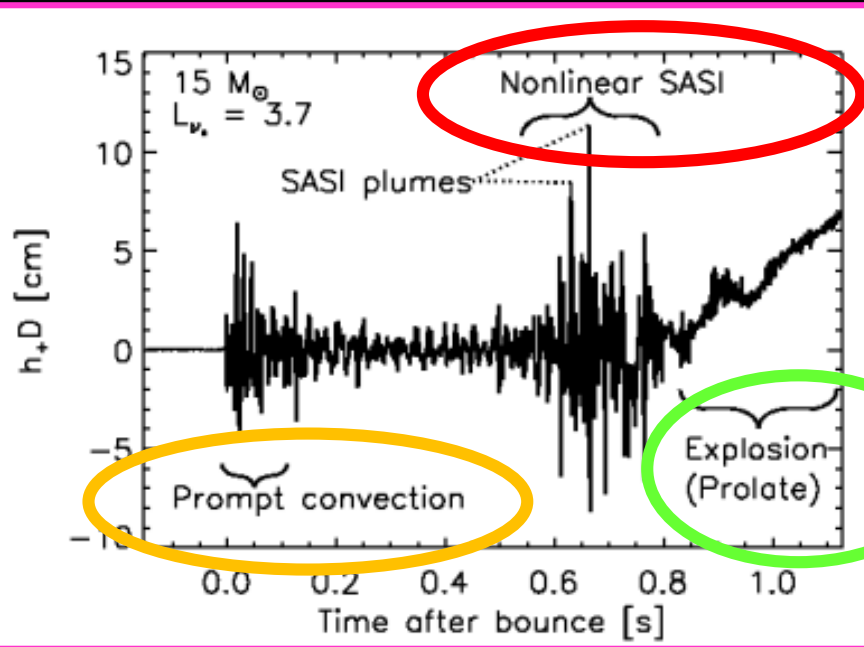
15 M_{sun} star
from Lentz et al. ('15)



(see also, Burrows et al. ('17), Melson et al. ('15), Lentz et al. ('15), Roberts et al. ('16), B. Mueller ('15), Takiwaki et al. ('16))

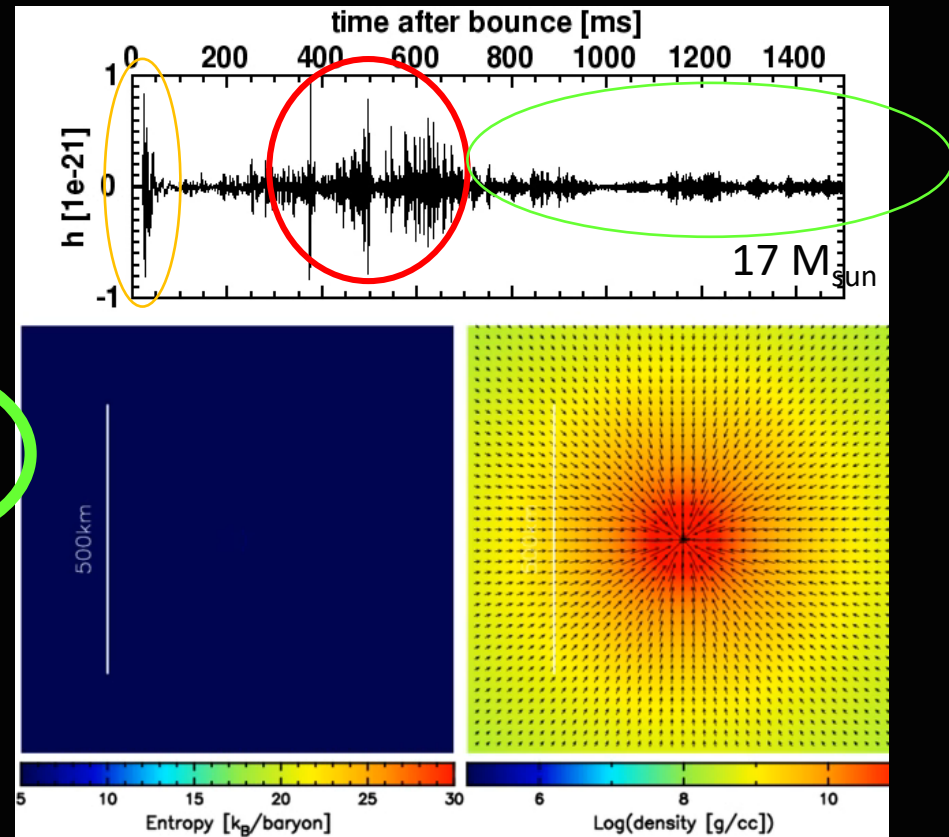
GW signatures from 2D neutrino-driven explosion (1/2)

Waveform from Murphy et al. (2009) ApJ



(Later confirmed by B. Mueller et al. ('13), ApJ, Yakunin et al. (2015), PRD)

Waveform from Nakamura et al. ('16) MNRAS



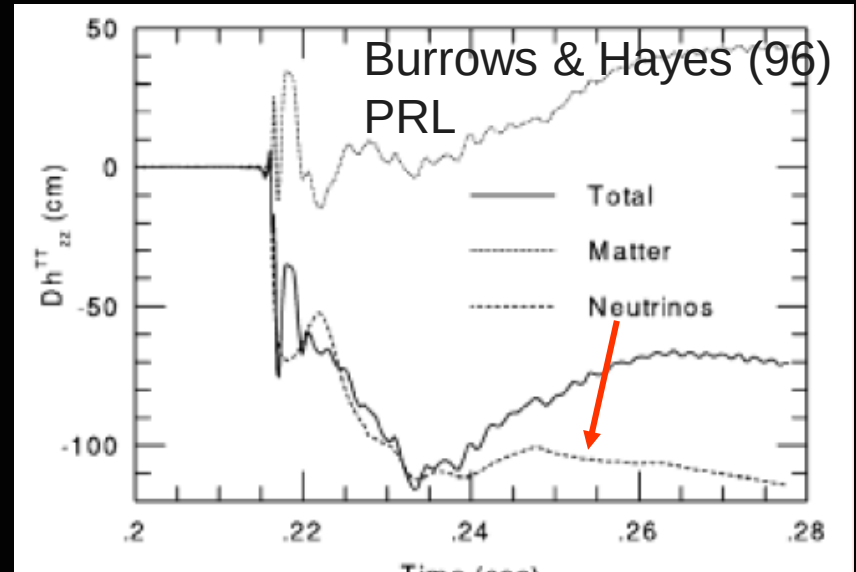
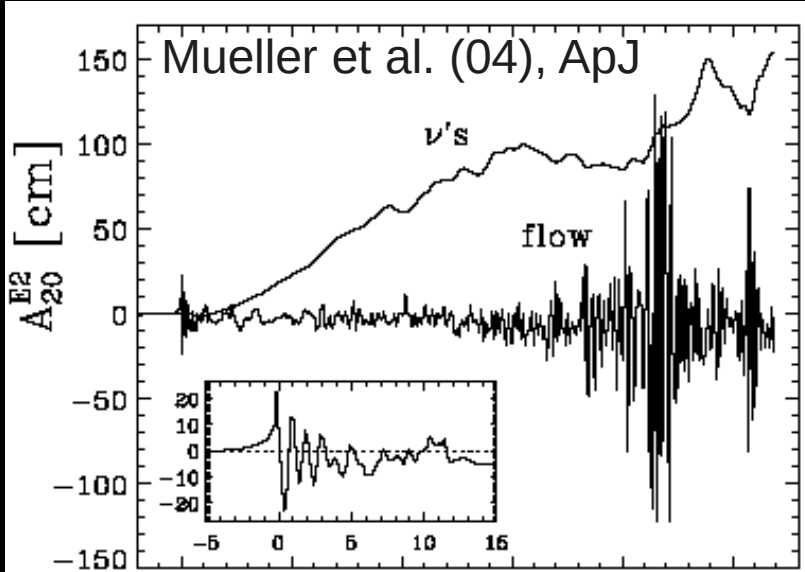
✓ Three generic phases in neutrino-driven models:

1. **Prompt-convection phase** : within ~50 ms post-bounce
2. **Non-linear phase (Convection/SASI)** : Downflows hit the PNS surface
3. **Explosion phase** : Long-lasting signal but terminates if BH forms

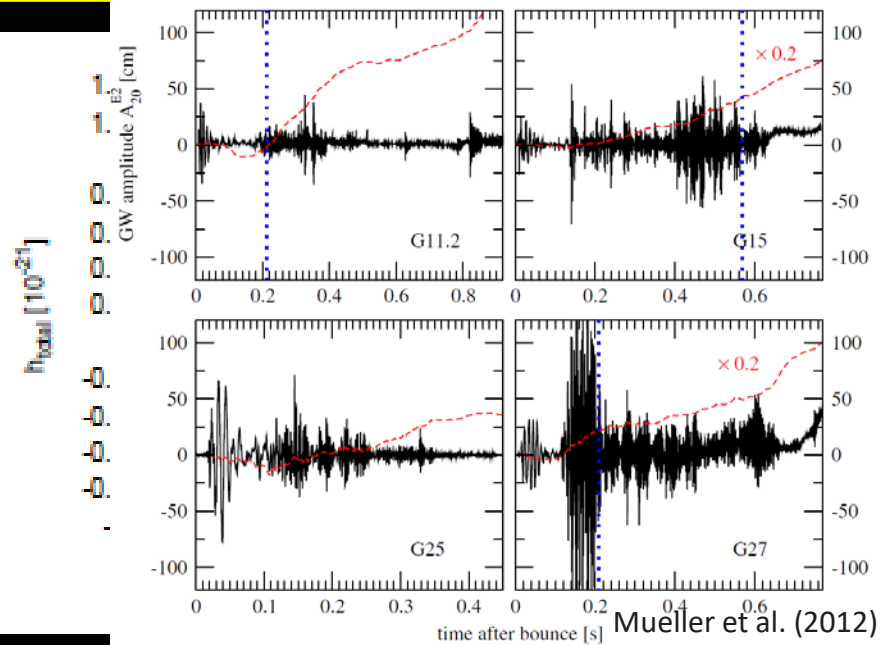
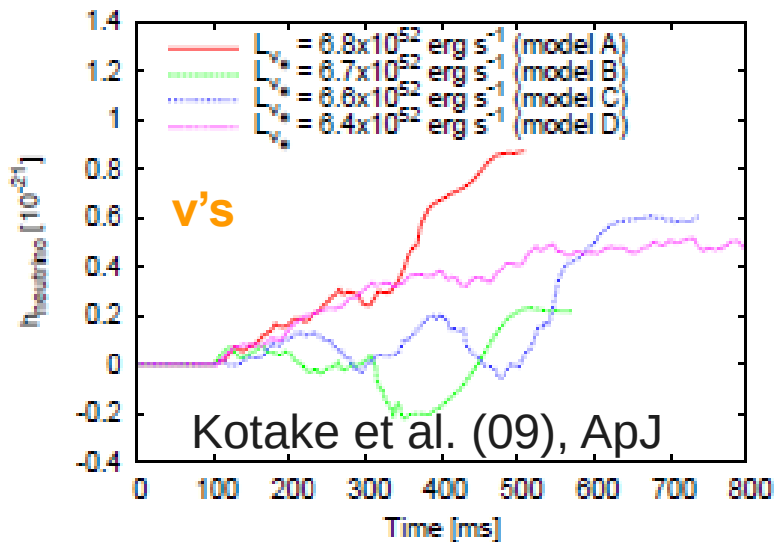
(Müller et al. (2004, ApJ), Cerda-Duran et al. (2013, ApJ))

✓ Waveforms have no template character: stochastic explosion processes.

GW signatures from 2D neutrino-driven explosion (2/2)



✓ GWs by anisotropic neutrino emission ~ bigger than the matter contribution !



GWs from anisotropic neutrino emission

$$h^{\mu\nu}(t, \mathbf{x}) = 4 \int \frac{T^{\mu\nu}(t - |\mathbf{x} - \mathbf{x}'|, \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} d^3x'$$

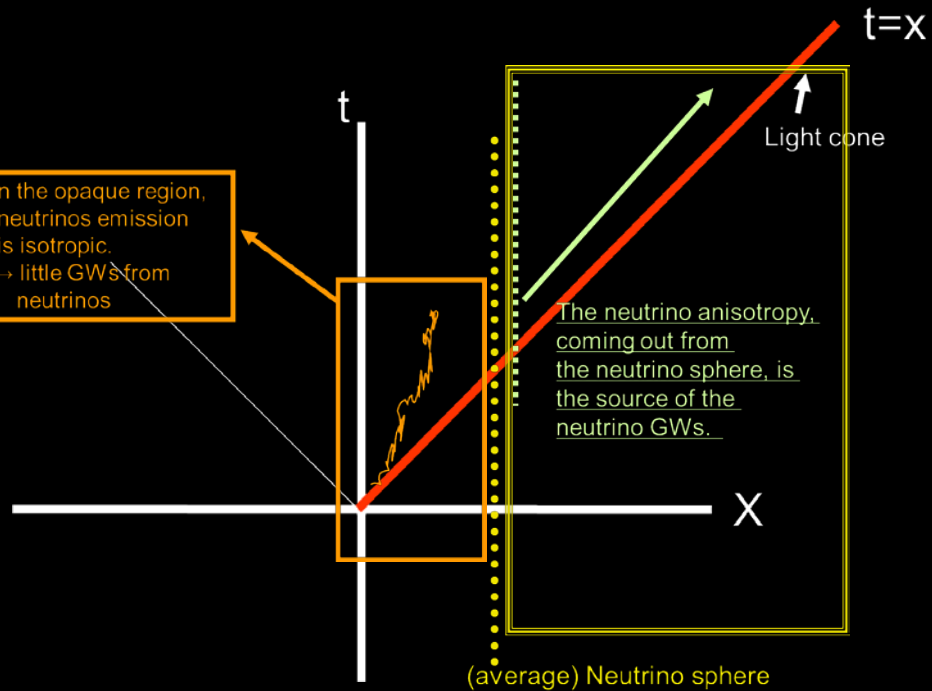
$$T^{\mu\nu} = T^{\mu\nu}_{\text{matter}} + \underline{T^{\mu\nu}_{\text{neutrino}}}$$

Epstein(78), Mueller & Janka (97)

$$h_{\nu}(t) = \frac{2G}{c^4 R} \int_0^t dt L_{\nu}(t') \alpha(t')$$

Neutrino anisotropy: degree of anisotropic neutrino emission (zero if spherical)

In 2D,
$$\alpha(t) = \frac{1}{L_{\nu}(t)} \int_{4\pi} d\Omega' \Phi(\theta') \frac{dl_{\nu}(\Omega', t')}{d\Omega'}$$



GWs from anisotropic neutrino emission

$$h^{\mu\nu}(t, \mathbf{x}) = 4 \int \frac{T^{\mu\nu}(t - |\mathbf{x} - \mathbf{x}'|, \mathbf{x}')}{|\mathbf{x} - \mathbf{x}'|} d^3x'$$

$$T^{\mu\nu} = T^{\mu\nu}_{\text{matter}} + \underline{T^{\mu\nu}_{\text{neutrino}}}$$

Epstein(78), Mueller & Janka (97)

$$h_{\nu}(t) = \frac{2G}{c^4 R} \int_0^t dt L_{\nu}(t') \alpha(t')$$

Neutrino anisotropy: degree of anisotropic neutrino emission (zero if spherical)

In 2D,
$$\alpha(t) = \frac{1}{L_{\nu}(t)} \int_{4\pi} d\Omega' \Phi(\theta') \frac{dL_{\nu}(\Omega', t')}{d\Omega'}$$

$\Delta t = x$

Typical amplitude :

$$|h_{\nu}| \sim 10^{-21} \left(\frac{\alpha}{0.01} \right) \left(\frac{L_{\nu}}{10^{52} \text{ erg/s}} \right) \left(\frac{\delta t}{1 \text{ sec}} \right) \left(\frac{R}{10 \text{ kpc}} \right)^{-1}$$

$$|h_{\nu}| \sim h_{\text{bounce}} \sim 10^{-21} \quad (10 \text{ kpc})$$

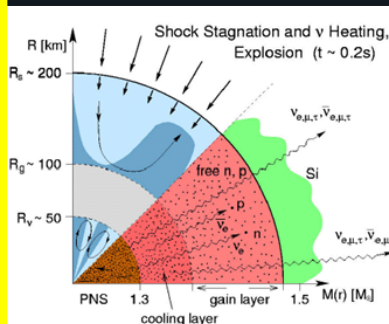
~ as large as the matter GW

Typical frequency :

$$t_{\nu} \sim \frac{1}{\sqrt{G\rho}} \geq 10 \text{ msec} \left(\frac{\rho_{\text{trap}}}{10^{11} \text{ gcm}^{-3}} \right)^{-1/2}$$

$$\nu_{\nu} \sim \frac{1}{t_{\nu}} \leq 100 \text{ Hz}$$

, longer than the matter GW signal because the dynamical time scale is determined at the position of neutrino sphere, where forms further out from the center.



Frequencies of GWs from neutrinos are typically lower than ~ 100 Hz.

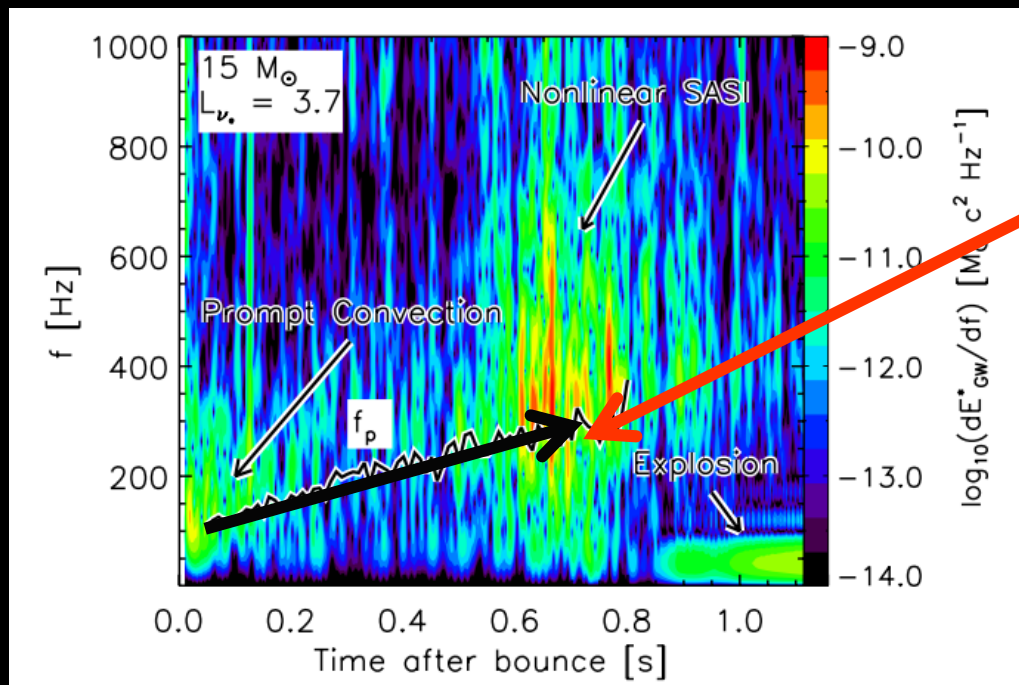
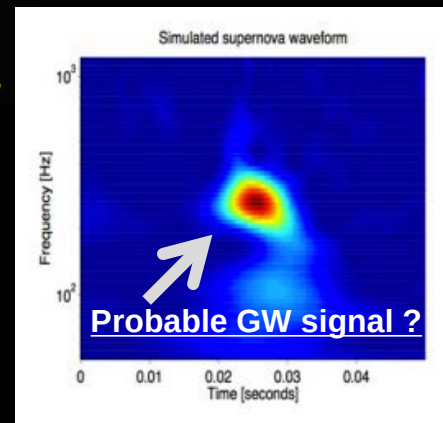
How to detect GWs with no-template features...

✓ **Excess power method:** Flanagan & Hugh (1998)

⇒ Decompose data-stream into time-frequency domains

⇒ Search for “hot” regions with excess power in the spectrogram !

✓ **GW spectrogram** from Murphy et al. ('09) ApJ.



★ Increase of typical frequency
⇒ the g-mode frequency of PNS

$$f_p = \frac{N}{2\pi} = \frac{1}{2\pi} \frac{GM}{R^2} \sqrt{\frac{(\Gamma - 1)m_n}{\Gamma k_b T}} \left(1 - \frac{GM}{Rc^2}\right)^{3/2}$$

M, R, T : mass, radius & temperature of PNS, Γ : stiffness of EOS

Due to mass accretion, M ↑, R ↓

⇒ f_p ↑

(see complete derivation in B. Mueller et al. ('13), ApJ)

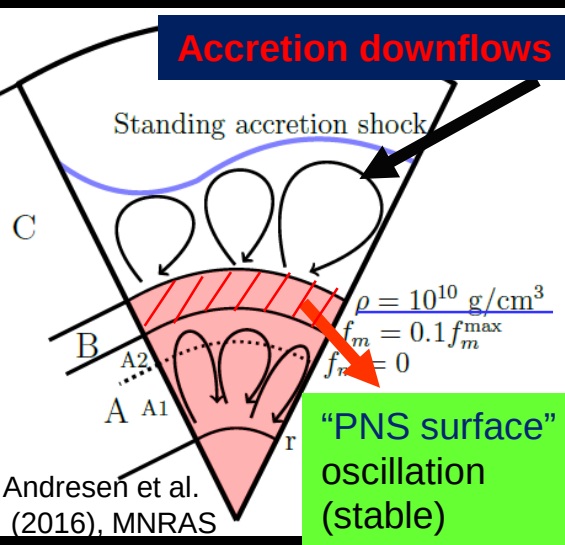
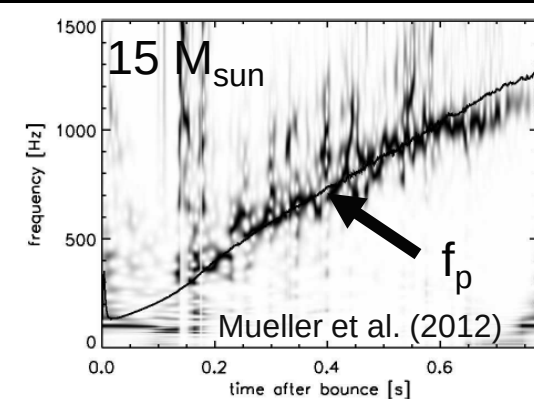
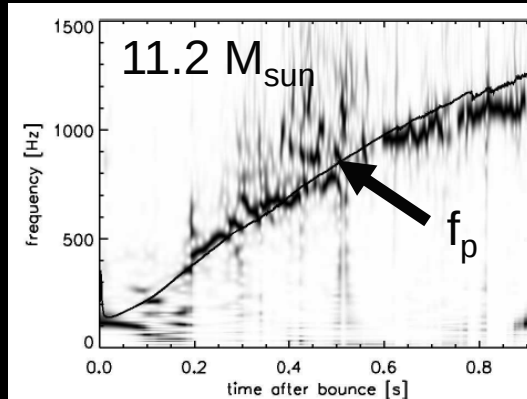
- ✓ (With no template character...) **Three generic phases are in the spectrogram !**
- ✓ Secular increase of typical GW frequency (f_p) **reflects the PNS evolution.**
- ✓ On top of f_p, the high frequency component comes from strong downflows to PNS.
- ✓ **These qualitative features : Common to more recent 2D and 3D models !**

"PNS" asteroseismology

: How to derive f_p ?



Knock a
"watermelon"
the best time
to eat



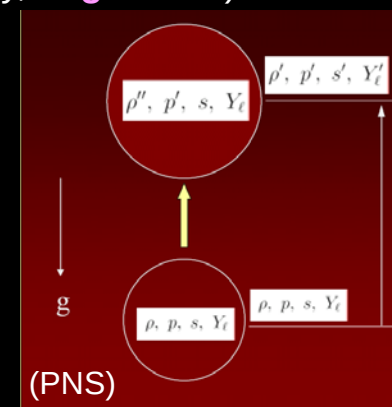
✓ Buoyancy frequency (Brunt-Vaisälä (BV) frequency, or **g-mode**)

$$N^2 = \frac{1}{\rho} \frac{\partial \Phi}{\partial r} \left[\frac{1}{c_s^2} \frac{\partial P}{\partial r} - \frac{\partial \rho}{\partial r} \right]$$

[check the unit !] $\Phi \sim \frac{\text{cm}^2}{\text{s}^2}$, $N^2 \sim \frac{\Phi}{r^2} \sim \frac{1}{\text{s}^2}$, $f_p = \frac{N}{2\pi}$ [Hz]
(characteristic frequency)

✓ Relativistic Extension of f_p

$$N^2 = \frac{\partial \alpha c^2}{\partial r} \frac{\alpha}{\rho h \phi^4} \left(\frac{1}{c_s^2} \frac{\partial P}{\partial r} - \frac{\partial \rho (1 + \epsilon/c^2)}{\partial r} \right)$$



PNS in hydrostatic equilibrium
: "1PN" ToV equation

$$\frac{dP}{dr} = - \frac{\rho GM}{R^2 \alpha}$$

$$\alpha = 1 - \frac{GM}{R}, \phi = \alpha^{-2}$$

$$h = 1 + \epsilon/c^2 + P/(\rho c^2)$$

$$P = \rho k_B T / m_n \text{ (non-relativistic baryon)}$$

$$\frac{d\rho}{dr} = \frac{m_n}{k_B T} \frac{dP}{dr}$$

$$c_s^2 = \Gamma k_B T / m_u$$

3+1 decomposition

$$g_{\mu\nu} = \begin{pmatrix} -\alpha^2 + \beta_i \beta^i & \beta_1 & \beta_2 & \beta_3 \\ \beta_1 & & & \\ \beta_2 & & \gamma_{ij} & \\ \beta_3 & & & \end{pmatrix}$$

$$\gamma_{ij} = \phi^4 \hat{\gamma}_{ij}$$

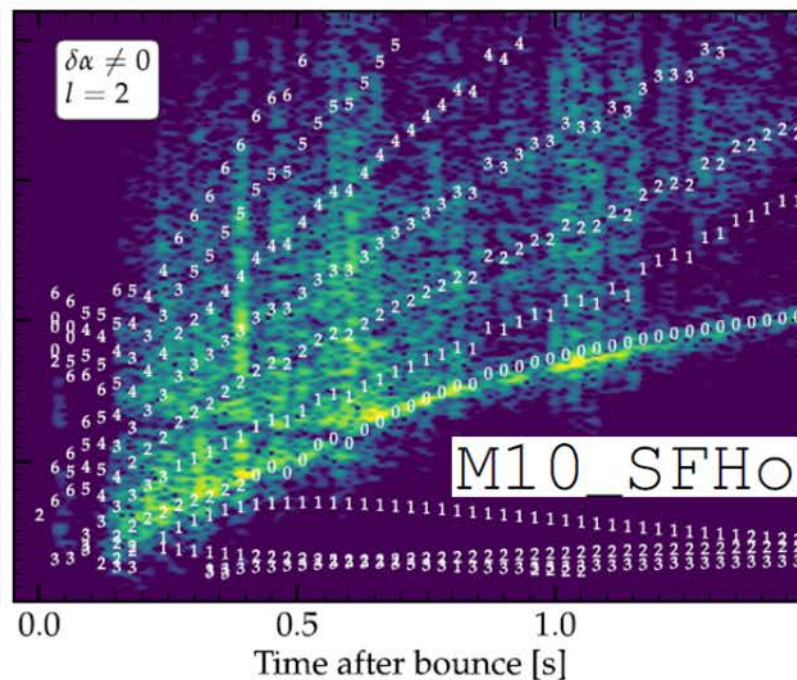
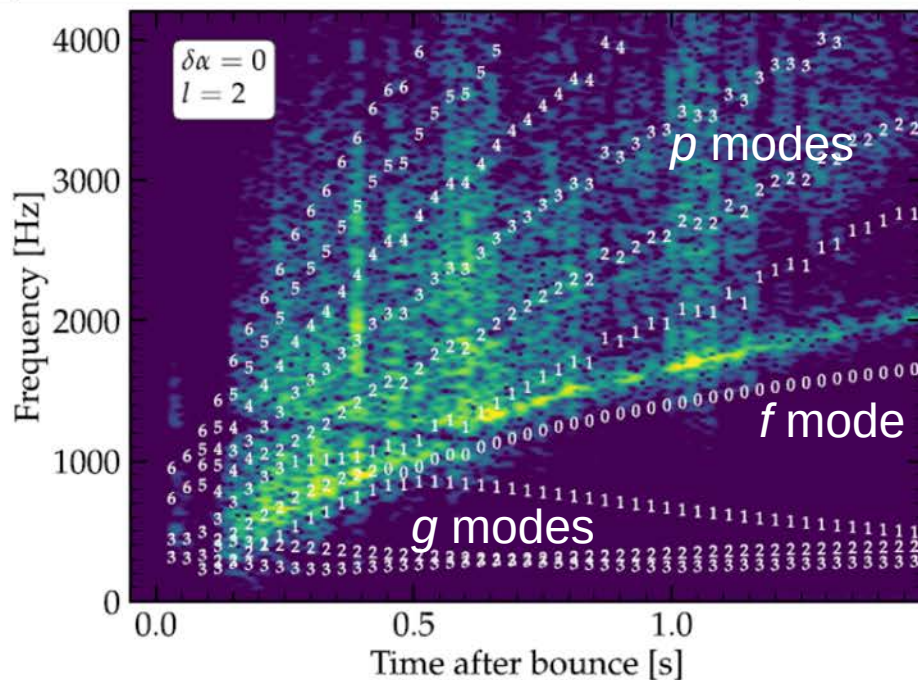
$$f_p = \frac{N}{2\pi} = \frac{1}{2\pi} \frac{GM}{R^2} \sqrt{\frac{(\Gamma - 1)m_n}{\Gamma k_B T}} \left(1 - \frac{GM}{Rc^2} \right)^{3/2}$$

Gravitational redshift

“PNS” asteroseismology (2/2)

✓ GR (1PN) correction important !

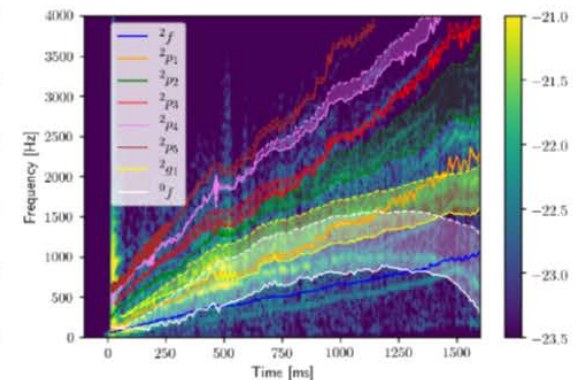
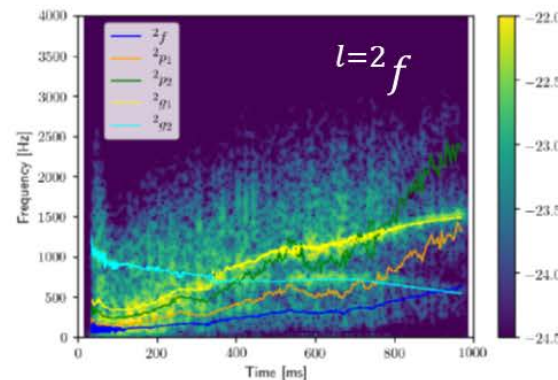
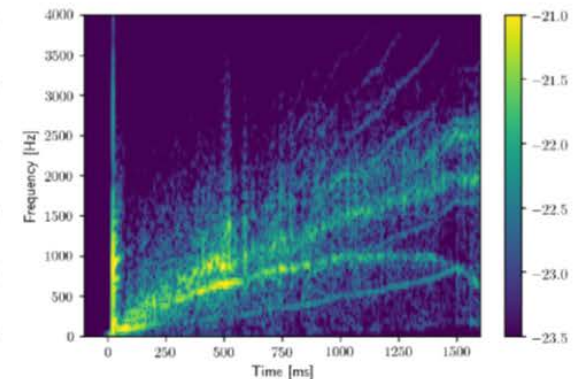
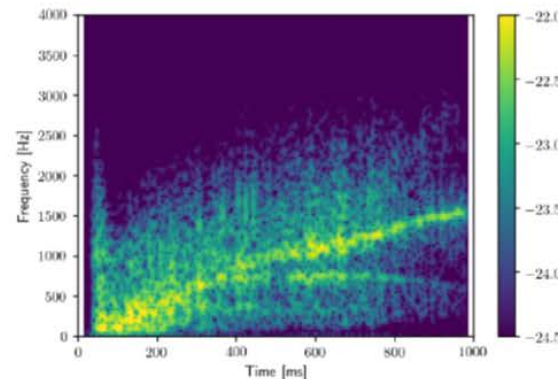
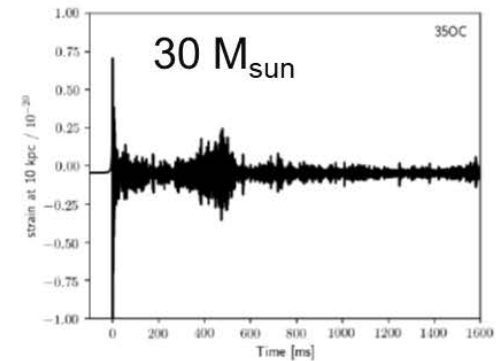
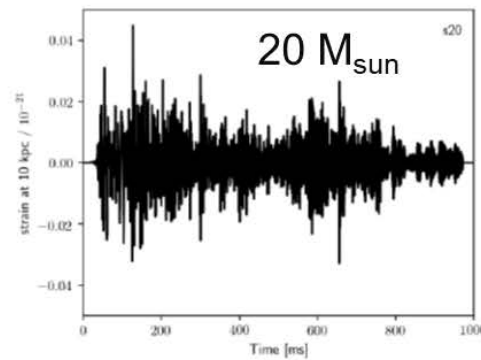
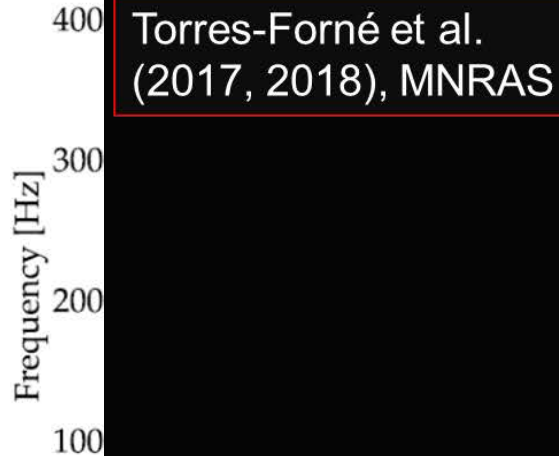
Morozova et al. (2018), MNRAS



“PNS” asteroseismology (2/2)

✓ GR (1PN) correction important !

Morozova et al. (2018), MNRAS



- ✓ Important:
Toward
“template-based”
GW search !
- ✓ Just started !
(e.g., Sotani et al. (2017). PRD)
- Detectability :
yet to be understood.

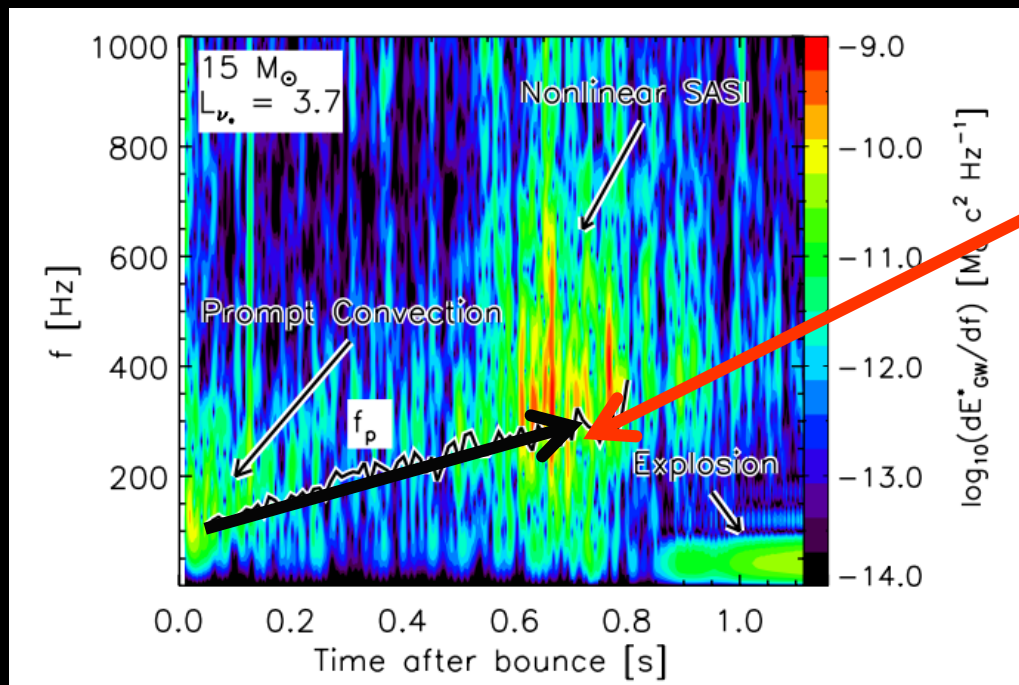
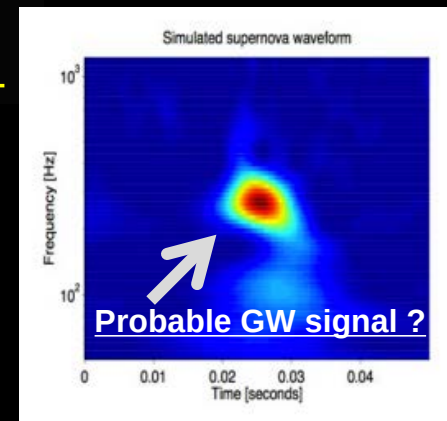
How to detect GWs with no-template features...

✓ **Excess power method:** Flanagan & Hugh (1998)

⇒ Decompose data-stream into time-frequency domains

⇒ Search for “hot” regions with excess power in the spectrogram !

✓ **GW spectrogram** from Murphy et al. ('09) ApJ.



★ Increase of typical frequency
⇒ the g-mode frequency of PNS

$$f_p = \frac{N}{2\pi} = \frac{1}{2\pi} \frac{GM}{R^2} \sqrt{\frac{(\Gamma - 1)m_n}{\Gamma k_b T}} \left(1 - \frac{GM}{Rc^2}\right)^{3/2}$$

M, R, T : mass, radius & temperature of PNS, Γ : stiffness of EOS

Due to mass accretion, M ↑, R ↓

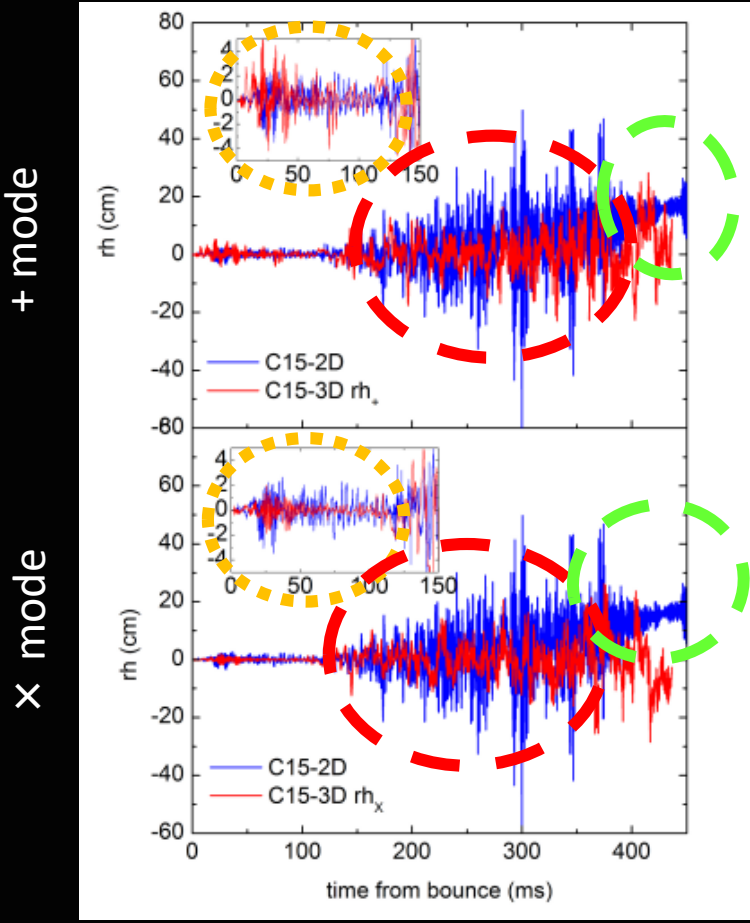
⇒ f_p ↑

(see complete derivation in B. Mueller et al. ('13), ApJ)

- ✓ (With no template character...) **Three generic phases are in the spectrogram !**
- ✓ Secular increase of typical GW frequency (f_p) **reflects the PNS evolution.**
- ✓ On top of f_p, the high frequency component comes from strong downflows to PNS.
- ✓ **These qualitative features : Common to more recent 2D and 3D models !**

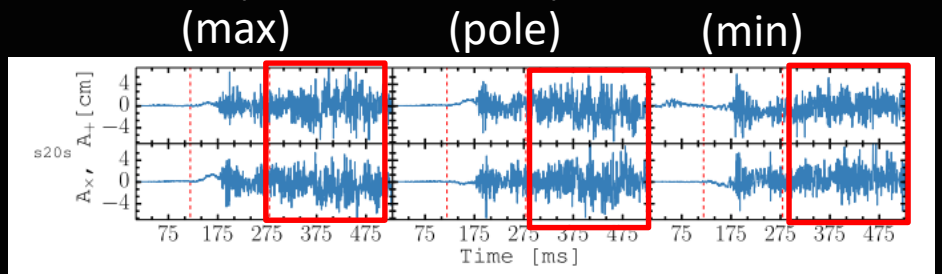
Recent GW predictions from 3D CCSN models with neutrino transport

- Yakunin, Mezzacappa et al. (2017)
- ✓ “Three generic phases” also seen in 3D
- ✓ 2D overestimates GW amp. relative to 3D

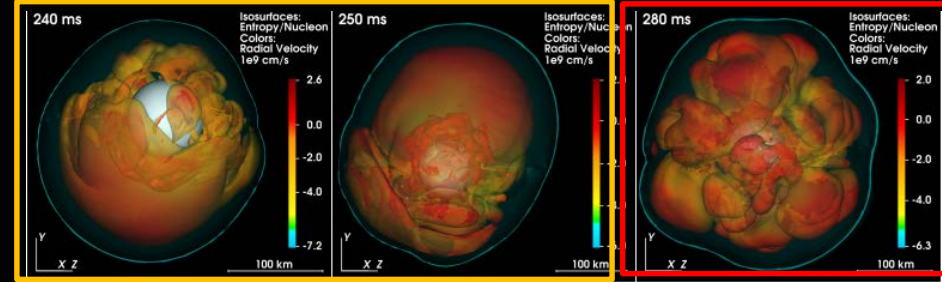


Based on 15 M_{sun} model
from Lentz et al. (2015), ApJL

- Andresen, B & E Müller and Janka (2017) MNRAS
- ✓ Wave amplitudes; rather insensitive to direction (to the observer).



especially when convection dominates over SASI.



- ✓ The horizon of LIGO is limited to nearby events.
Third generation detectors (ET) could detect any Galactic events !

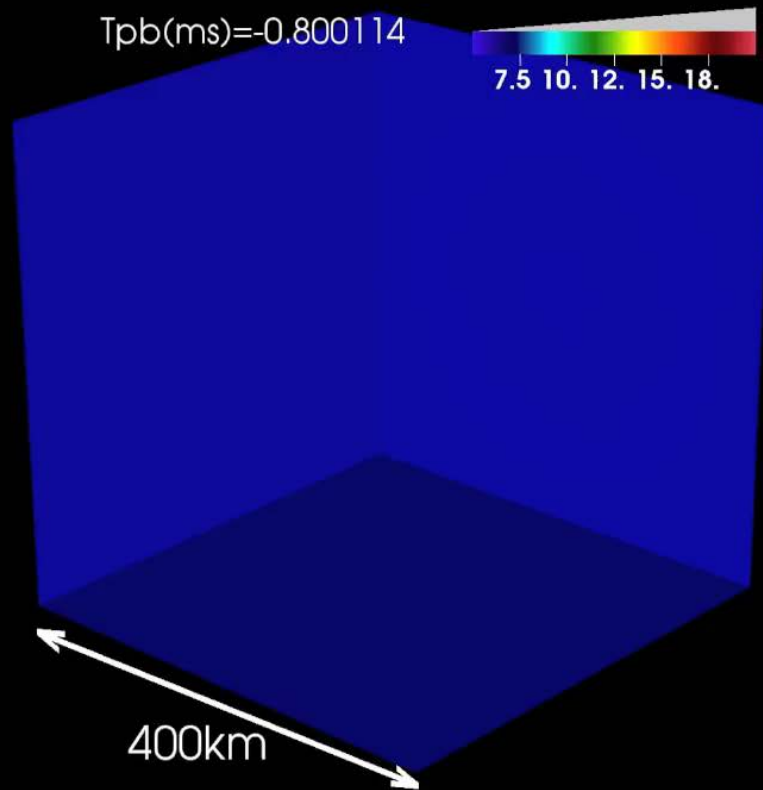
	s27				s20			
	Low	High	Total	Low/High	Low	High	Total	Low
AdvLIGO	3.7	4.5	8.8	0.82	5.3	7.7	9.4	0.82
ET-C	50.0	64.0	81.3	0.78	73.9	109.3	131.9	0.83
ET-B	78.5	73.7	107.7	1.07	113.9	127.0	170.6	0.74

GW signatures from 3D-GR models with **strong SASI** vs. **weak SASI** activity

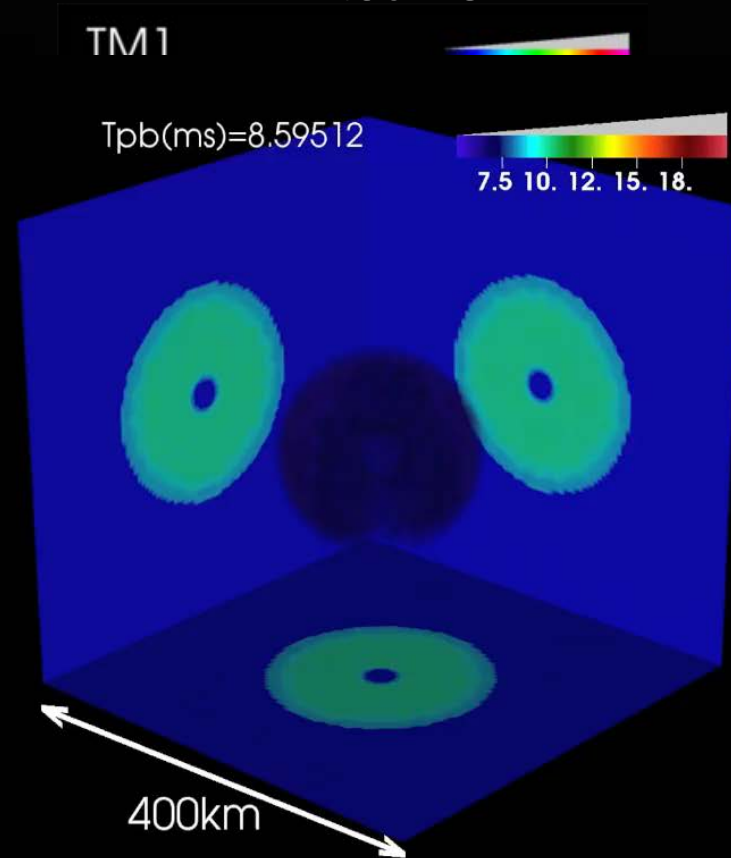
(from Kuroda, KK, & Takiwaki ApJL (2016), see also Andresen, B, E Müller and Janka (2017))

- ✓ Two EOSs → **SFHx** (Steiner et al. (2013), fits well with experiment/NS radius, Steiner+(2011)),
HS(TM1) (Shen et al. (1998)).
- ✓ $15 M_{\text{sun}}$ star (Woosley & Weaver (1995))

SFHx :softer



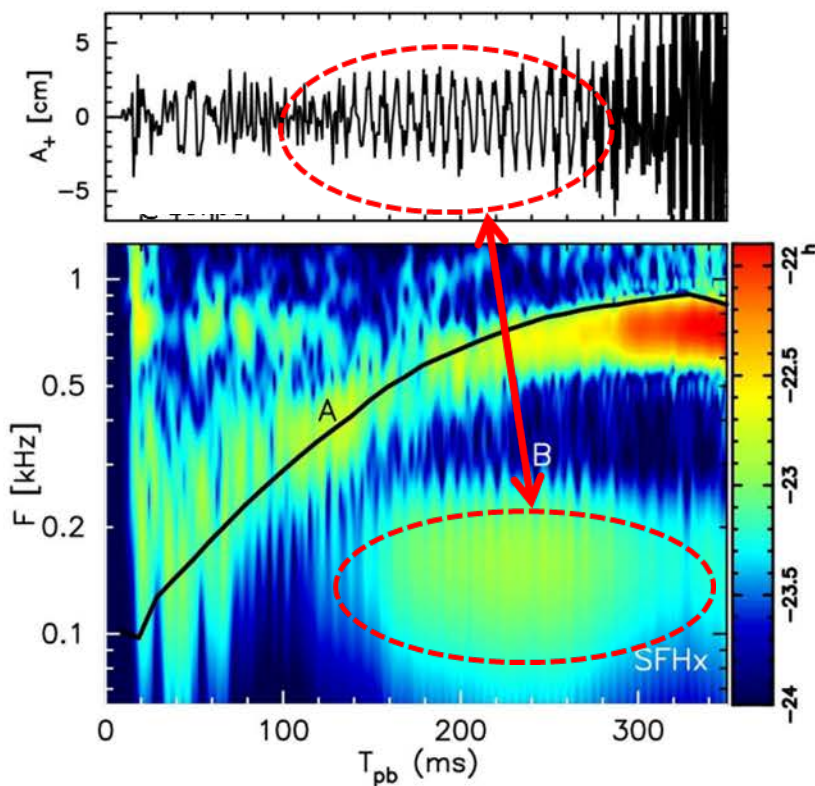
TM1 :stiffer



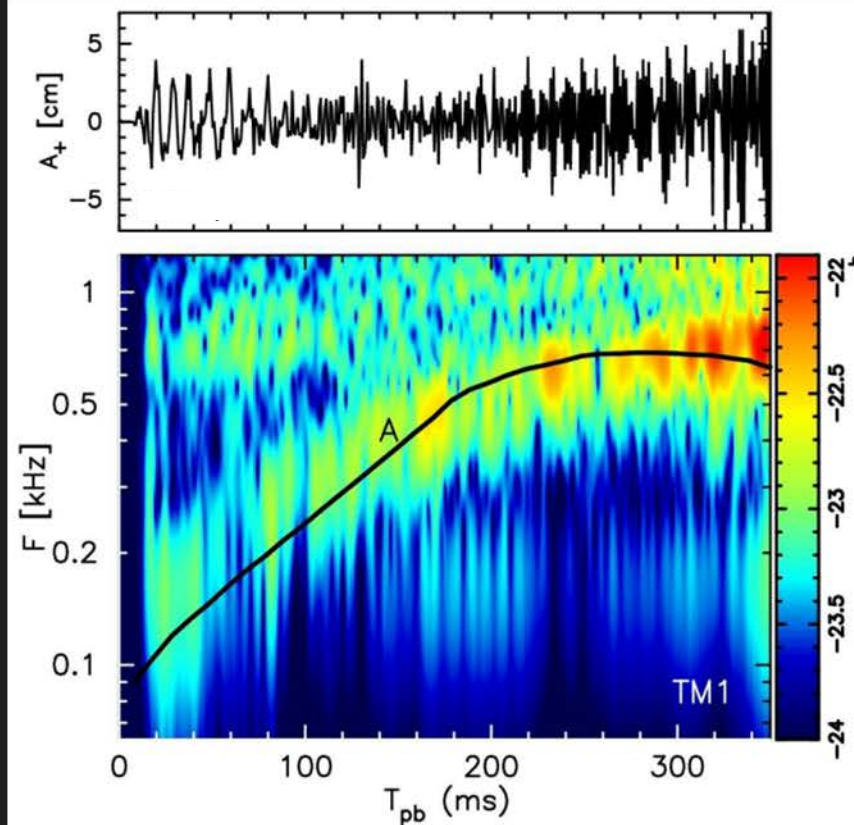
✓ **SASI activity higher for softer EOS** (due to high growth rate, e.g., Foglizzo et al. ('06)).

GW Spectrograms from 3D-GR models with strong SASI vs. weak SASI activity

SFHx :softer



TM1 :stiffer



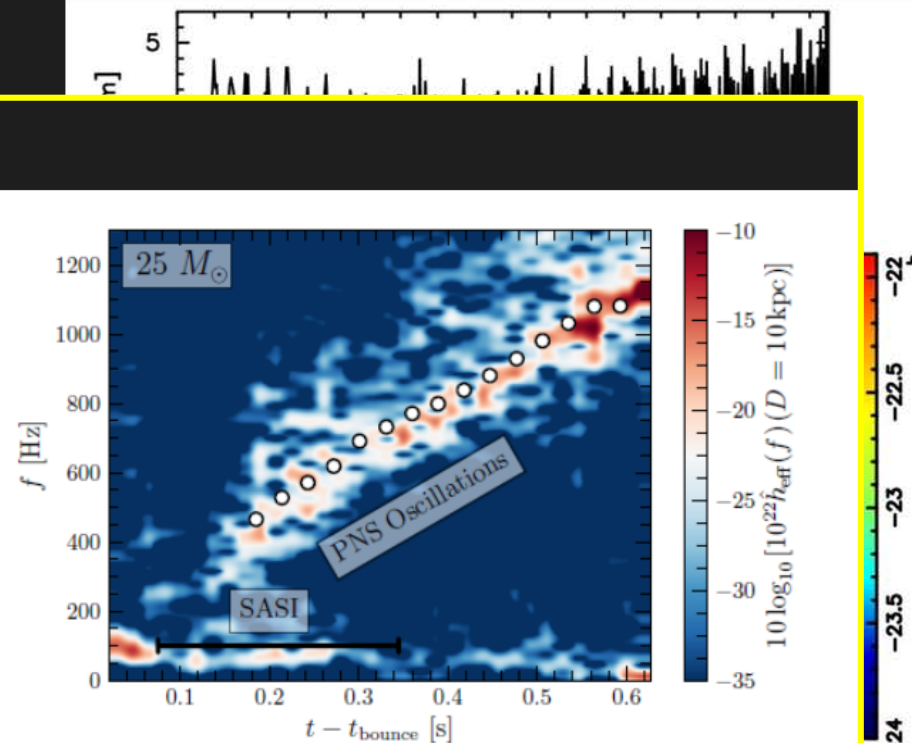
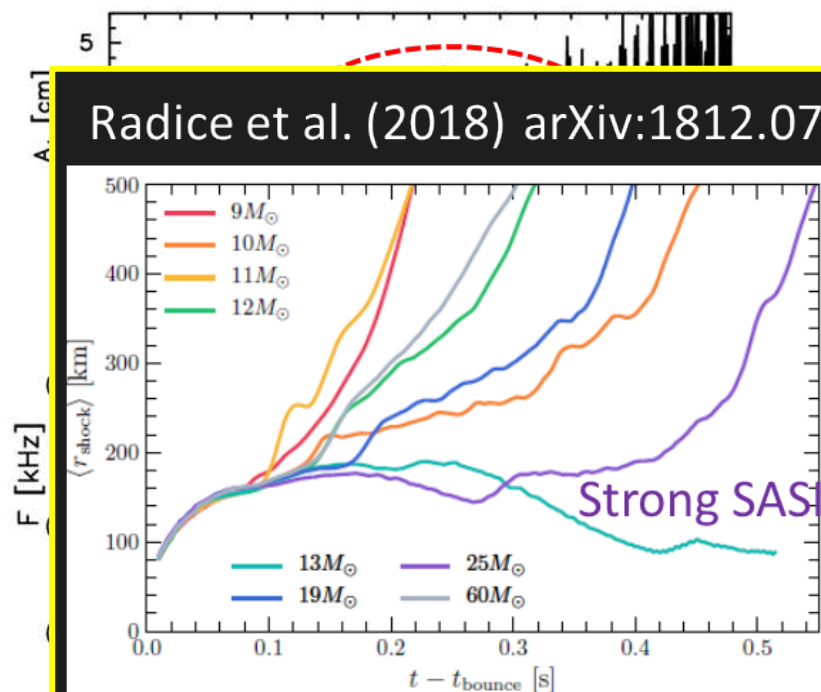
- ✓ The **quasi-periodic modulation** is associated with SASI, clearly visible **with realistic EOS**.
- ✓ By **coherent network analysis** of LIGO, VIRGO, and KAGRA, the detection horizon is only 2~3 kpc, but could extend out to 100 kpc when **ET and CE are on-line (>2035)**.
- ✓ **Detection of neutrinos** (Super-K, IceCube) important to get timestamp of GW detection.
- ✓ The SASI activity, if very high, results in characteristic signatures in both GWs and neutrino signals (e.g., Tamborra et al. (2012) for SASI-induced neutrino signals).

GW Spectrograms from 3D-GR models with strong SASI vs. weak SASI activity

SFHx :softer

TM1 :stiffer

Radice et al. (2018) arXiv:1812.07703



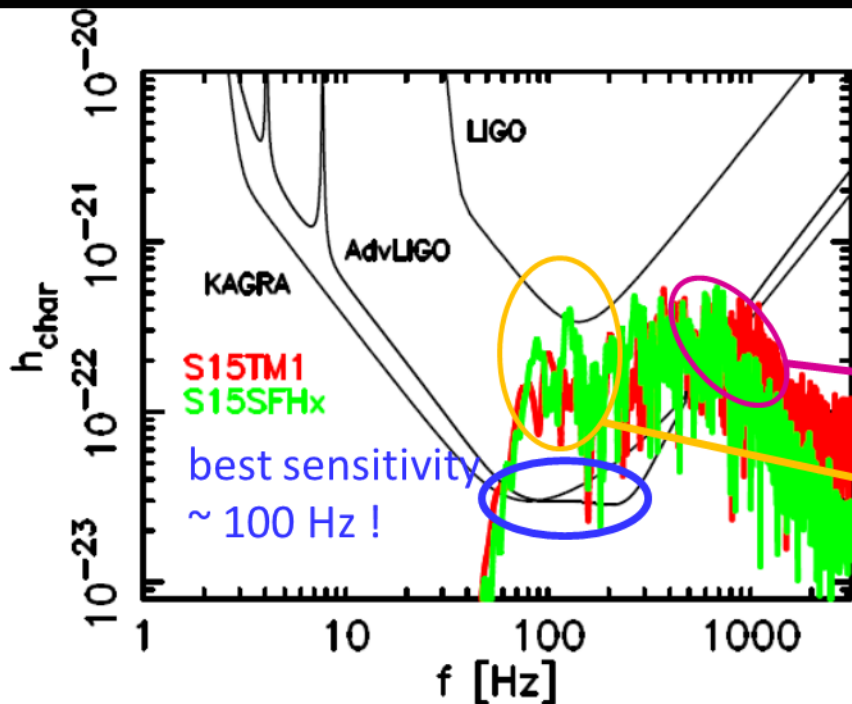
- ✓ The **quasi-periodic modulation** is associated with SASI, clearly visible **with realistic EOS**.
- ✓ By **coherent network analysis** of LIGO, VIRGO, and KAGRA, the detection horizon is only 2~3 kpc, but could extend out to 100 kpc when **ET and CE are on-line (>2035)**.
- ✓ **Detection of neutrinos** (Super-K, IceCube) important to get timestamp of GW detection.
- ✓ The SASI activity, if very high, results in characteristic signatures in both GWs and neutrino signals (e.g., Tamborra et al. (2012) for SASI-induced neutrino signals).

GW Spectrograms from 3D-GR models with strong SASI vs. weak SASI activity

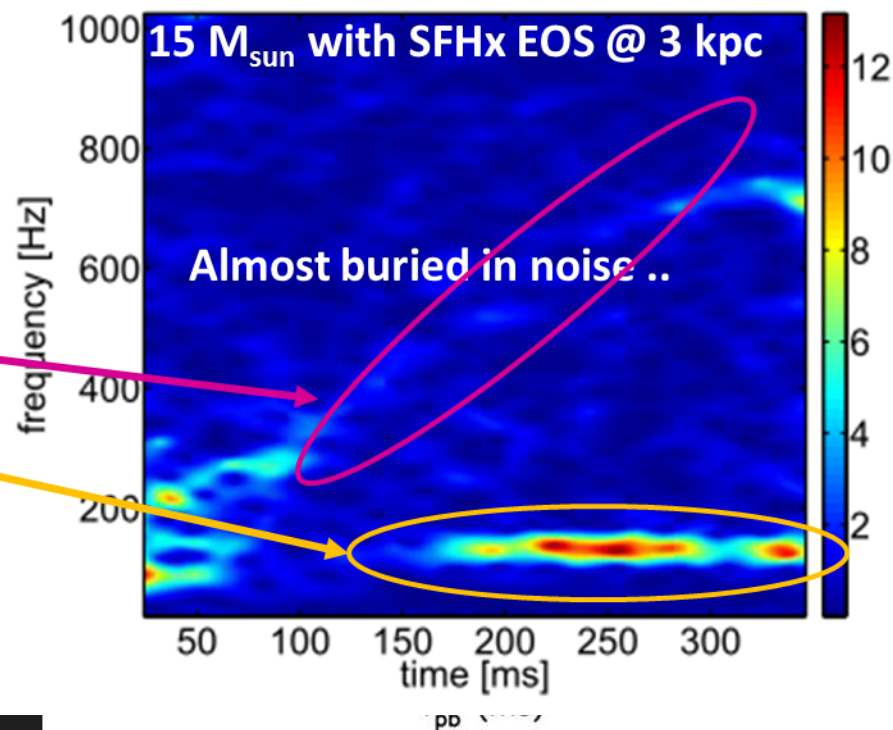
SFHx :softer

TM1 :stiffer

Sensitivity curves and model predictions



The reconstructed GW spectrogram

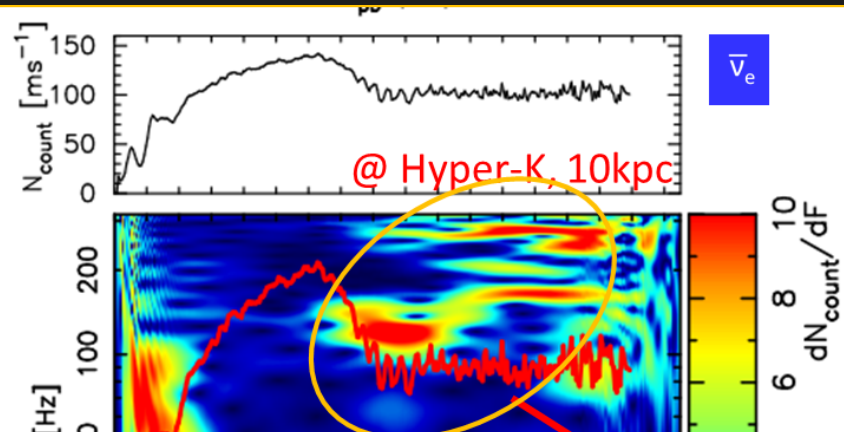


- ✓ The **quasi-periodic modulation** is associated with SASI, clearly visible **with realistic EOS**.
- ✓ By **coherent network analysis** of LIGO, VIRGO, and KAGRA, the detection horizon is only 2~3 kpc, but could extend out to 100 kpc when **ET and CE are on-line (>2035)**.
- ✓ **Detection of neutrinos** (Super-K, IceCube) important to get timestamp of GW detection.
- ✓ The SASI activity, if very high, results in characteristic signatures in both GWs and neutrino signals (e.g., Tamborra et al. (2012) for SASI-induced neutrino signals).

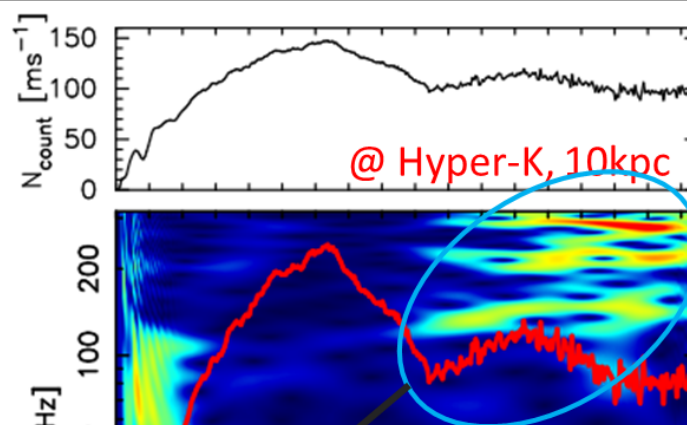
GW Spectrograms from 3D-GR models with strong SASI vs. weak SASI activity

SFHx :softer

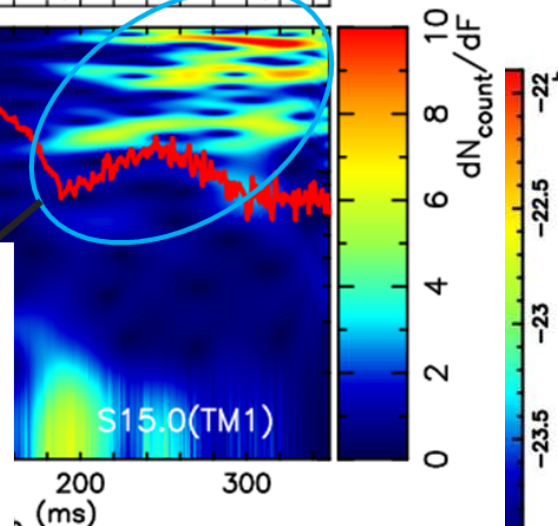
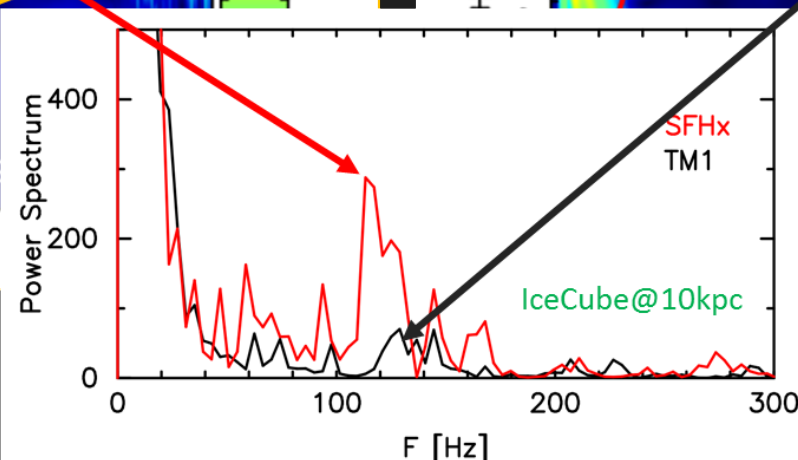
TM1 :stiffer



@ Hyper-K, 10kpc



@ Hyper-K, 10kpc



$$T_{\text{SASI}} = \tau_{\text{adv}} + \tau_{\text{ac}} = \int_{r_V}^{r_{\text{sh}}} \frac{dr}{|v_r|} + \int_{r_V}^{r_{\text{sh}}} \frac{dr}{c_s - |v_r|}$$

- ✓ The **quasi-periodic modulation** is associated with SASI, clearly visible **with realistic EOS**.
- ✓ By **coherent network analysis** of LIGO, VIRGO, and KAGRA, the detection horizon is only 2~3 kpc, but could extend out to 100 kpc when **ET and CE are on-line (>2035)**.
- ✓ **Detection of neutrinos** (Super-K, IceCube) important to get timestamp of GW detection.
- ✓ The SASI activity, if very high, results in characteristic signatures in both GWs and neutrino signals (e.g., Tamborra et al. (2012) for SASI-induced neutrino signals).

“New” GW messenger is Circular Polarization of GW :Non-axisymmetric instabilities

(incl. low T/|W|, spiral SASI)

Stokes Parameters:

Hayama et al. (2016), PRL (see also Klimenko et al. (2015) PRD)

$$\begin{pmatrix} \langle h_R(f, \hat{n}) h_R(f', \hat{n}')^* \rangle & \langle h_L(f, \hat{n}) h_R(f', \hat{n}')^* \rangle \\ \langle h_R(f, \hat{n}) h_L(f', \hat{n}')^* \rangle & \langle h_L(f, \hat{n}) h_L(f', \hat{n}')^* \rangle \end{pmatrix} \\ = \frac{1}{4\pi} \delta_D^2(\hat{n} - \hat{n}') \delta_D(f - f') \\ \times \begin{pmatrix} I(f, \hat{n}) + \underline{V(f, \hat{n})} & Q(f, \hat{n}) - iU(f, \hat{n}) \\ Q(f, \hat{n}) + iU(f, \hat{n}) & I(f, \hat{n}) - \underline{V(f, \hat{n})} \end{pmatrix}$$

V parameter =
Asymmetry of right and left modes

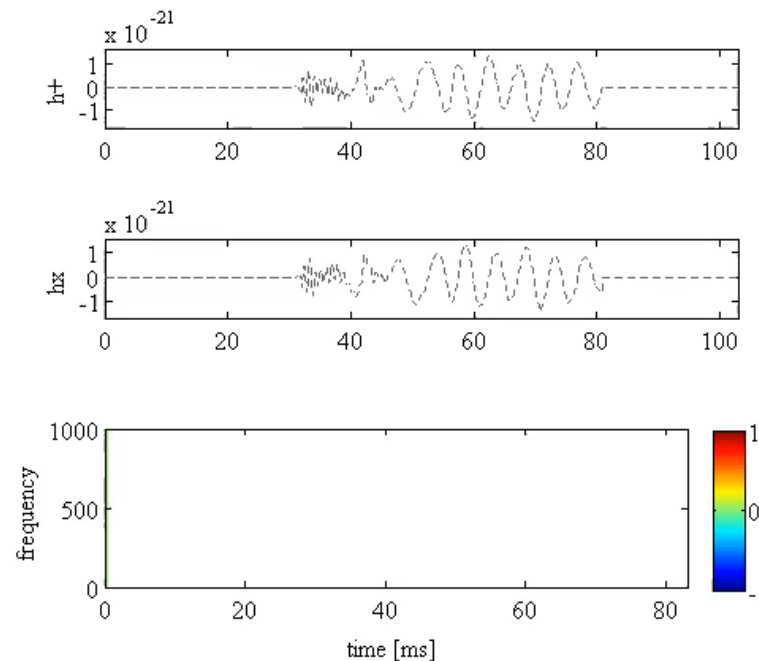
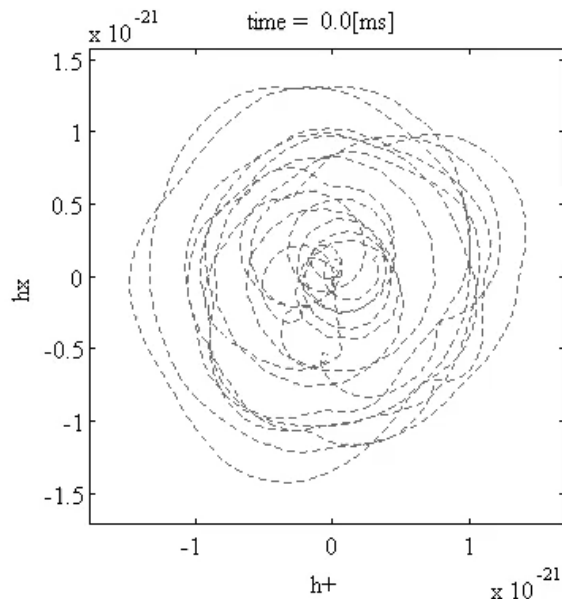
$$h_R := (h_+ - i h_\times) / \sqrt{2}$$

$$h_L := (h_+ + i h_\times) / \sqrt{2}$$

(See definitions in Seto and Taruya (2007), PRL)

Rapidly rotating 15 M_{sun} (early postbounce phase)
from Kuroda, Takiwaki, KK (2014) PRD

@10kpc



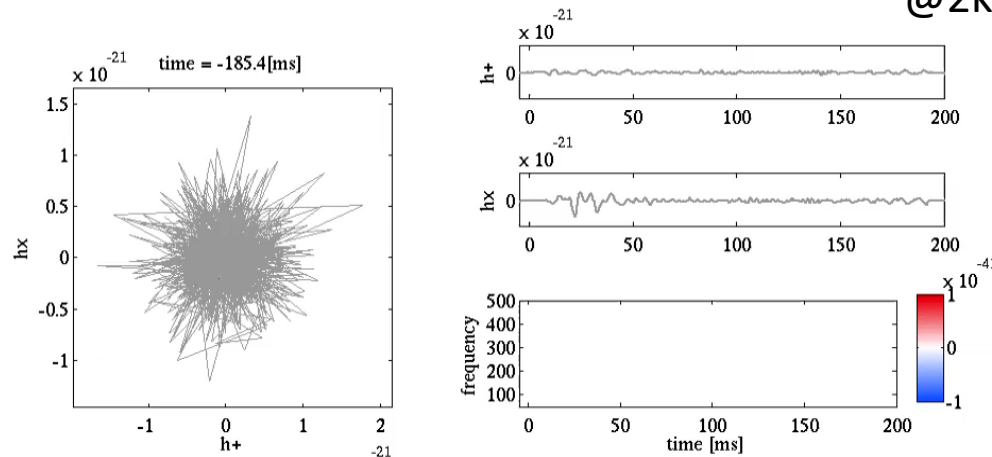
CP (if seen from the spin axis)
:evidence of
“rapid rotation”.

What about Circular GW polarization in “Non-rotating” progenitors ?

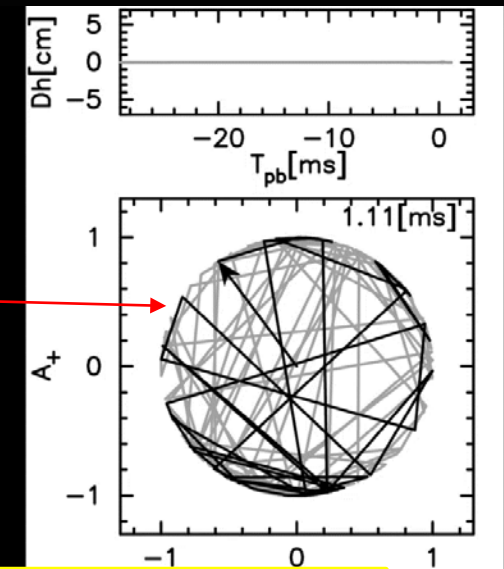
Hayama, KK et al. (2018)

Non-rotating 11.2 M_{sun} star ; Convection dominant

@2kpc



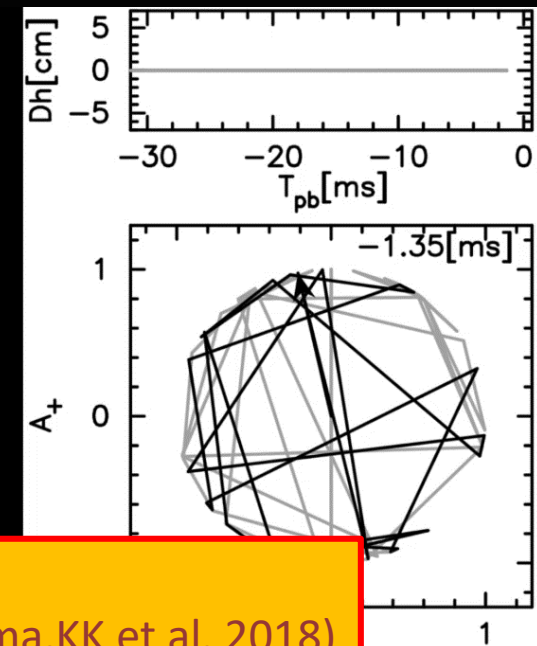
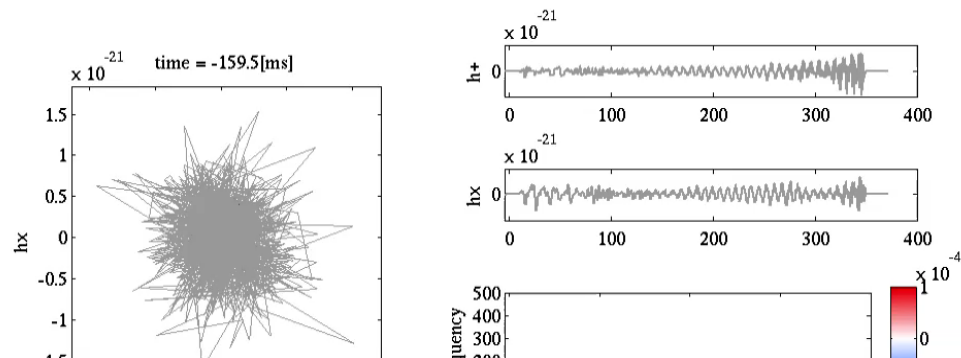
Normalized polarization
(e.g.,
 $\frac{x}{\sqrt{x^2 + y^2}}$,
 $x = h_+, h_x$)



If the core is convection-dominant (likely for low ξ stars), no clear signature of CP !

Non-rotating 15 M_{sun} star (SFHx EOS); SASI dominant

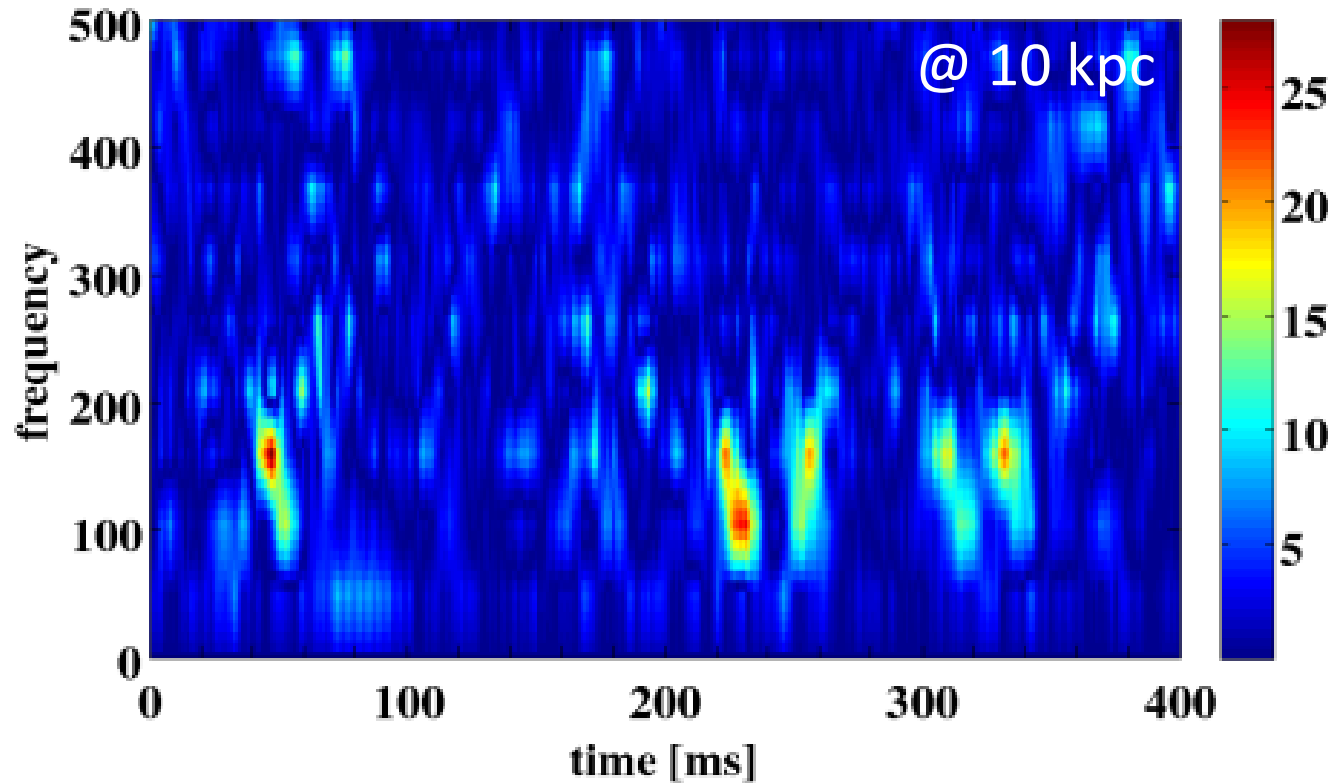
@5kpc



If the SASI dominant (likely for high ξ stars), clear signature of CP !

⇒ indication of SASI motions non-spherical mass accretion (Hayama, KK et al. 2018)

SNR of Circular Polarization of GW relative to background



- ✓ The detection of GW amplitude is within several kpc using LIGO (e.g., Andresen et al. (2017))
 - ✓ The detection of CP could extend (far) beyond the detection horizon of GW waveform !
 - ✓ The CP would provide new window to detect GW signals !
- (Hayama, Takiwaki, KK, Kuroda, MNRAS Letters, (2018))

The Origin of the Nobel-Prize-awarded BHs ($7 \sim 40 M_{\text{sun}}$) ?



The Nobel Prize in Physics 2017

Rainer Weiss, Barry C. Barish, Kip S. Thorne

Share this: [f](#) [G+](#) [t](#) [+](#) [e](#) 1.3K

The Nobel Prize in Physics 2017



© Nobel Media. Ill. N. Elmehed
Rainer Weiss
Prize share: 1/2



© Nobel Media. Ill. N. Elmehed
Barry C. Barish
Prize share: 1/4

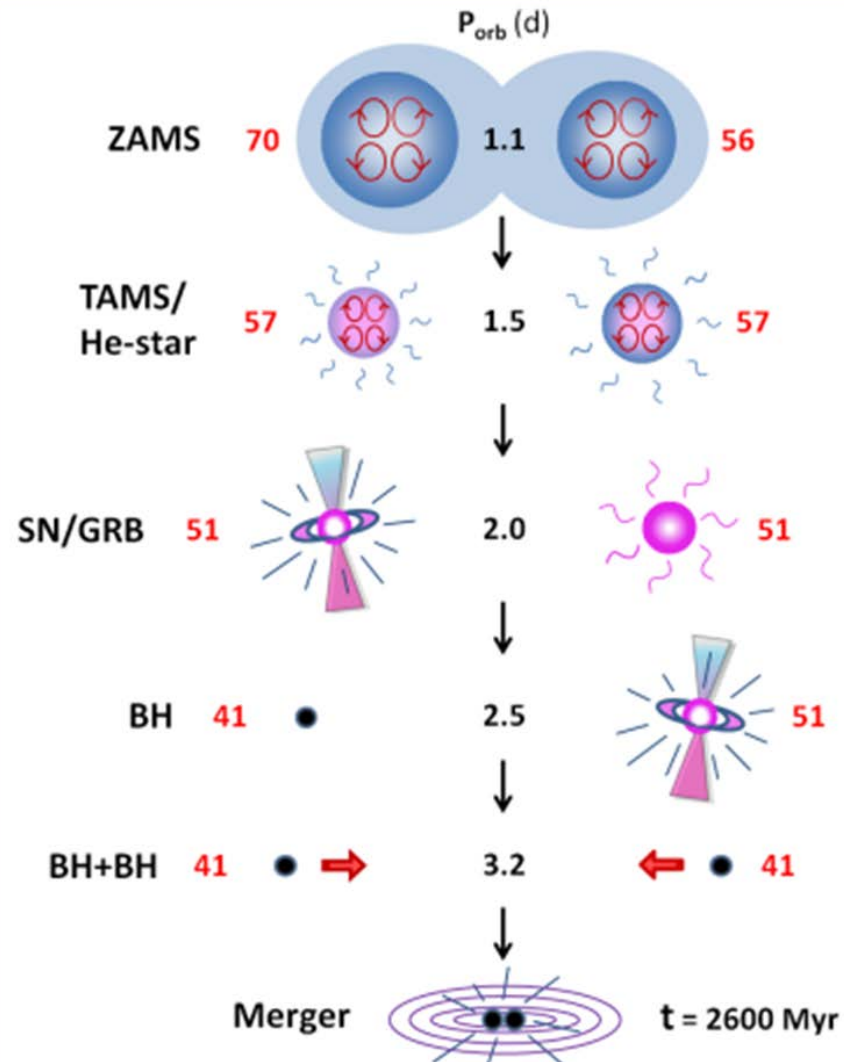


© Nobel Media. Ill. N. Elmehed
Kip S. Thorne
Prize share: 1/4

The Nobel Prize in Physics 2017 was divided, one half awarded to Rainer Weiss, the other half jointly to Barry C. Barish and Kip S. Thorne *"for decisive contributions to the LIGO detector and the observation of gravitational waves"*.

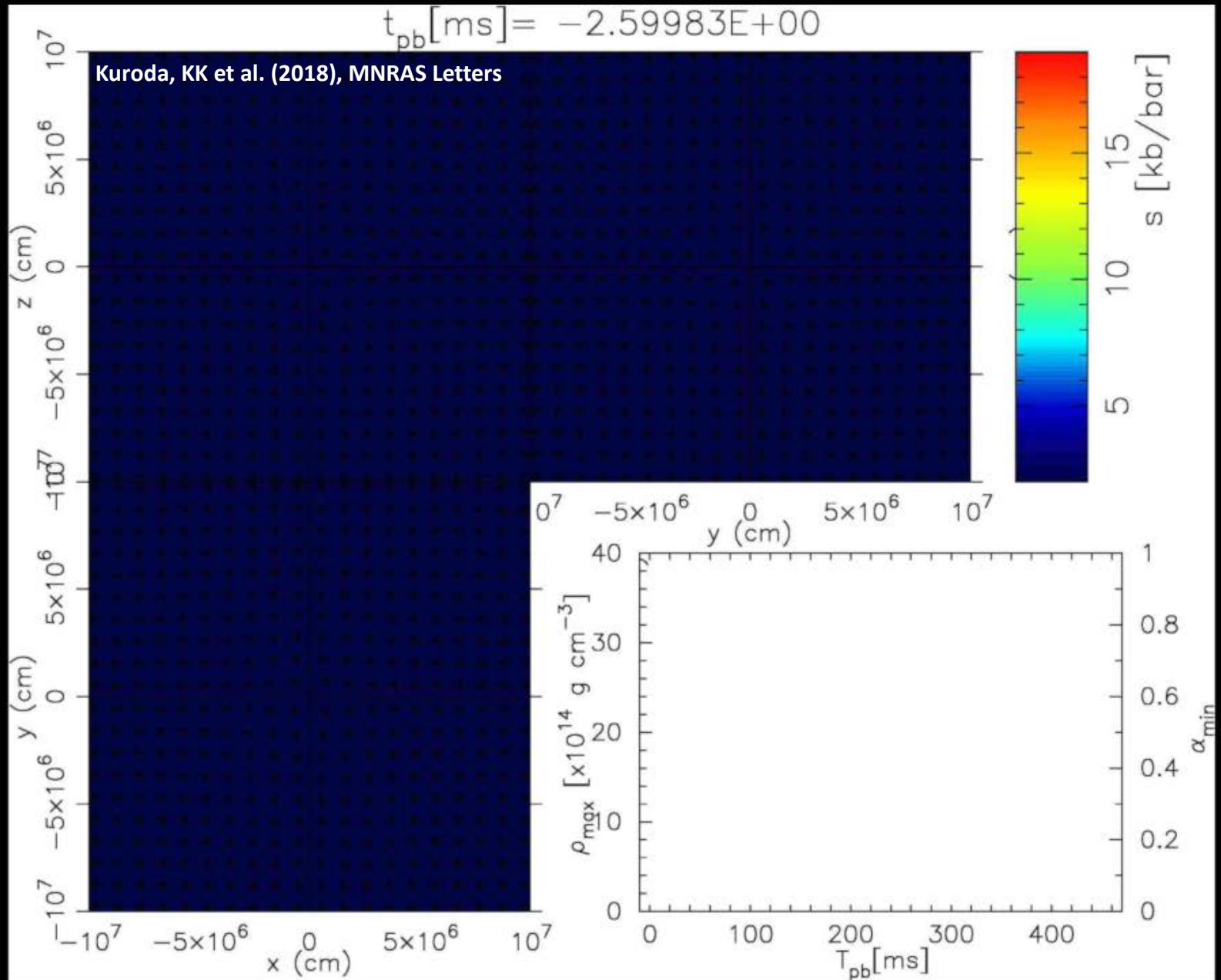
✓ **Low metallicity** environment needed for large stellar mass BH formation.
(e.g., Kinugawa et al. (2014,2016))

✓ One of them.. "Isolated binaries"



Marchant, Langer, Podsiadlowski et al. (2006)

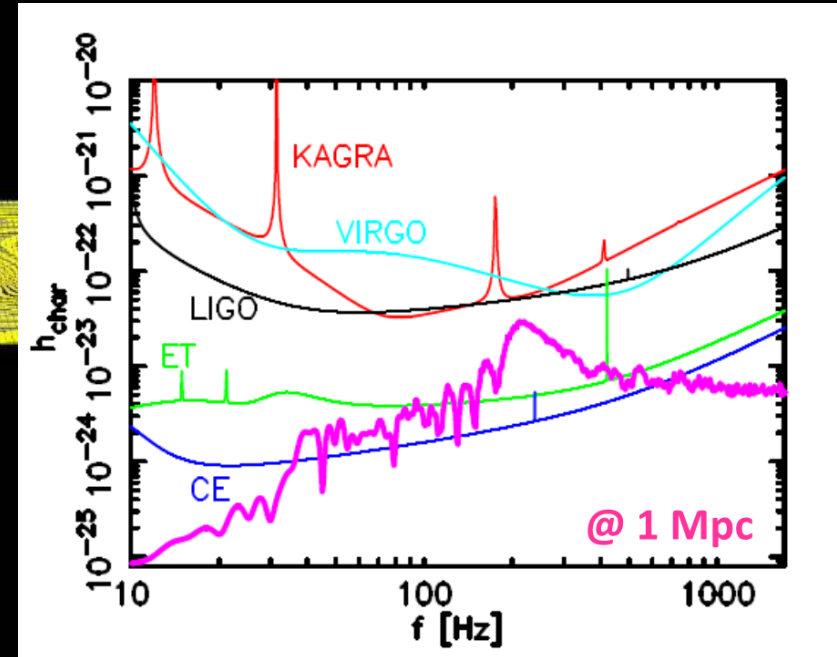
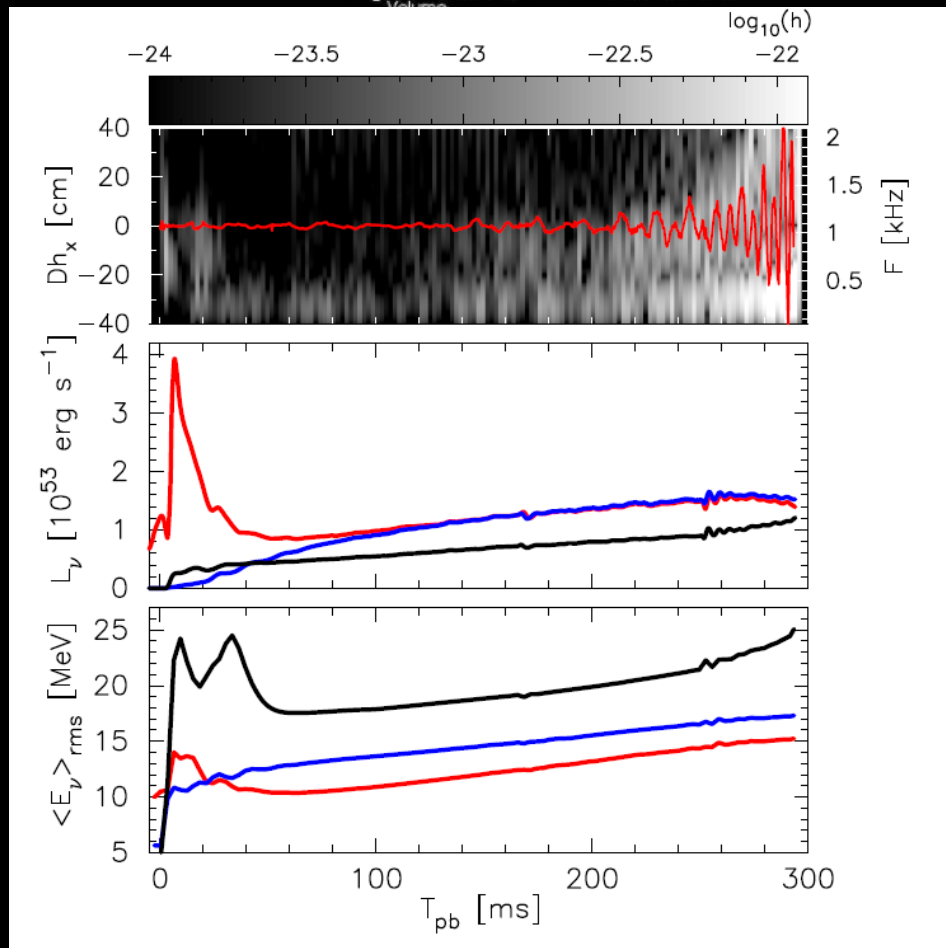
✓ 3D-GR results of $70 M_{\text{sun}}$ ($M_{\text{CO}} \sim 28.5 M_{\text{sun}}$) (progenitor from Takahashi et al. (2014))



✓ 3D-GR results of 70 M_{sun} ($M_{\text{CO}} \sim 28.5 M_{\text{sun}}$)

Z70.0(LS220)
 $T_{\text{pb}}(\text{ms}) = -2.59983$

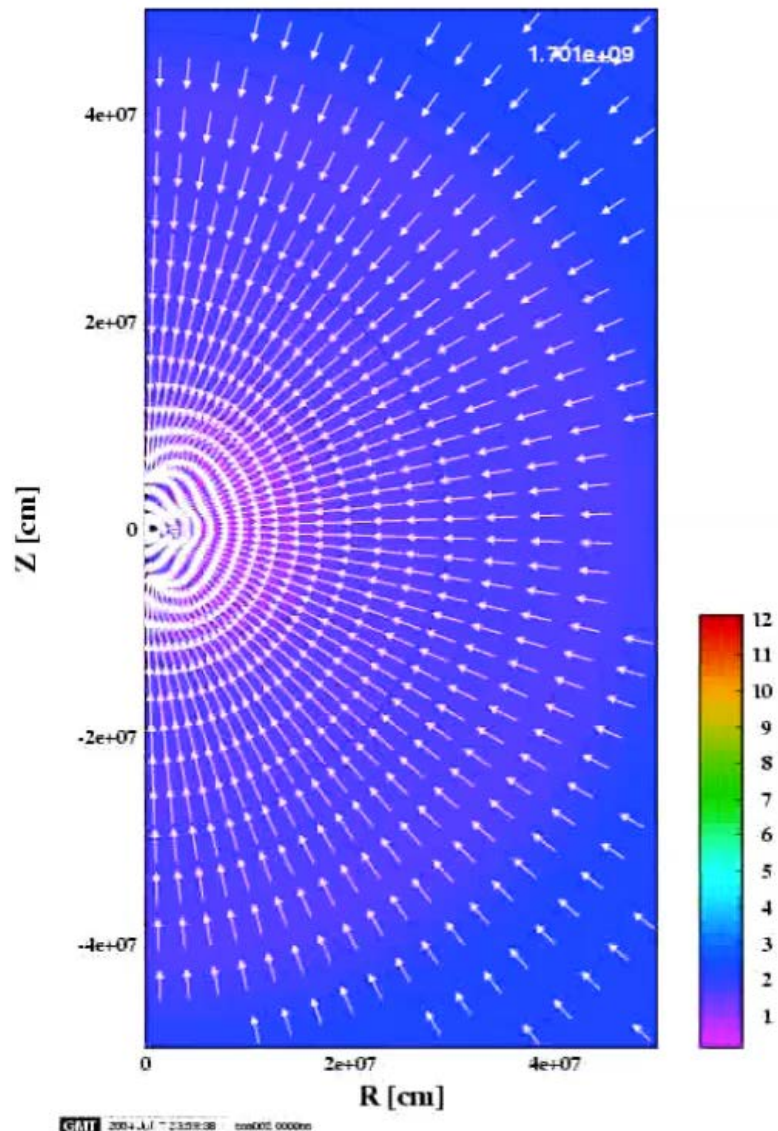
Kuroda, KK et al. (2018), MNRAS Letters



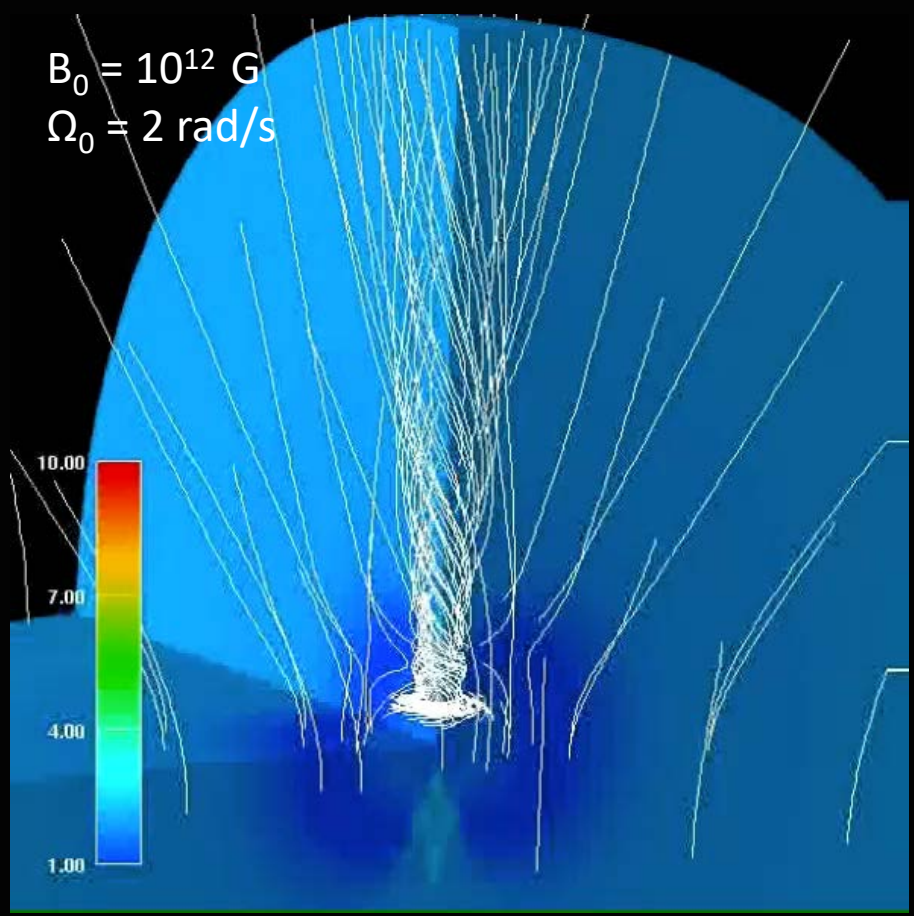
- ✓ The first BH forming simulation in 3D !
- ✓ Before the BH formation, monotonic increase of neutrino luminosity and rms energy. (consistent with 1D, e.g., Sumiyoshi+ (2006), Fischer+ (2009), Huedepohl+(2016))
- ✓ Strong GW emission is visible to 1 Mpc, but not O(100) Mpc...

Switching gears to **MHD mechanism (rapid rotation required !!)** My research life....

Kotake, Yamada, Sato (2003) ApJ



Magnetohydrodynamics Simulations
Kotake et al. (2004), (2006)

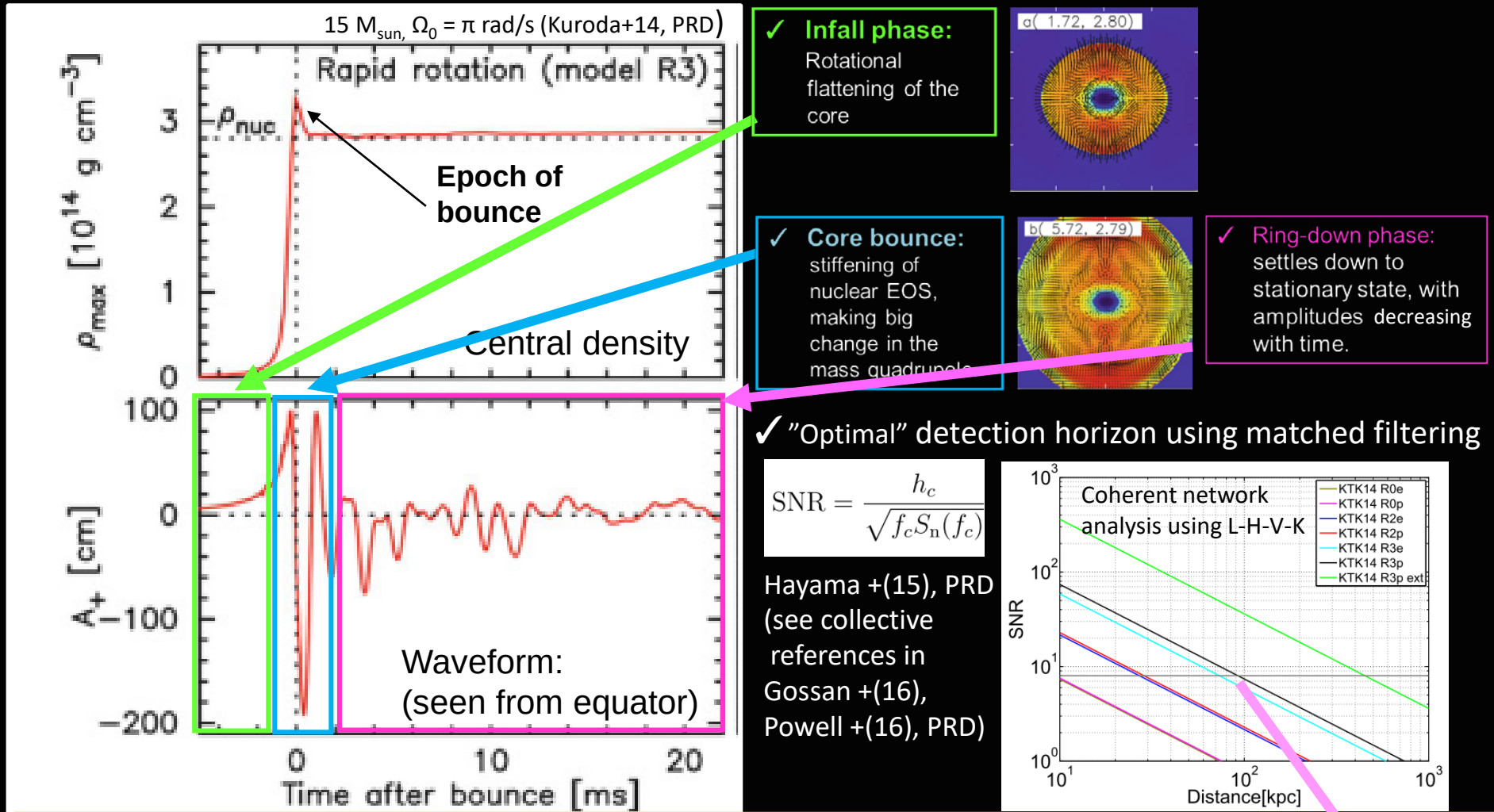


See recent developments in Moesta et al.(2015)
Masada et al. (2015), Ramirez et al. (2016), Sawai et al. (2014)

Switching gears to **MHD mechanism (rapid rotation required !!)**

GW from Rapidly Rotating Core-Collapse and Bounce

(Dimmelmeier et al. (07, PRL), Scheidegger et al. (10, A&A) Ott et al. (12, ApJ), Abdikamalov+(14, PRD), Kuroda+(14,PRD))



Bounce GW signal (in the context of rapidly rotating collapse and bounce):

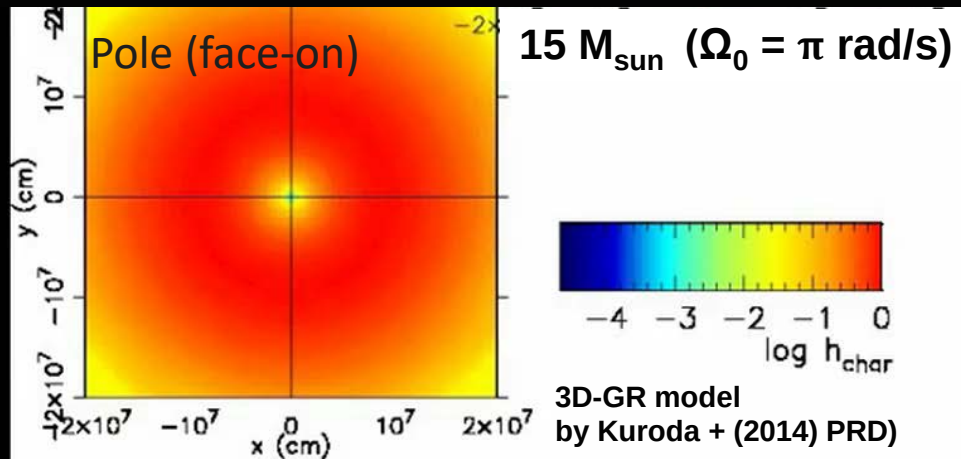
- ✓ Characterized by “one” big spike at bounce followed by smaller peaks: “type I” signal
- ✓ Matched filtering (or PCA) likely applicable.
- ✓ Horizon distance can reach beyond LMC (50kpc)

GWs from (Rotation-induced) Non-Axisymmetric Instabilities

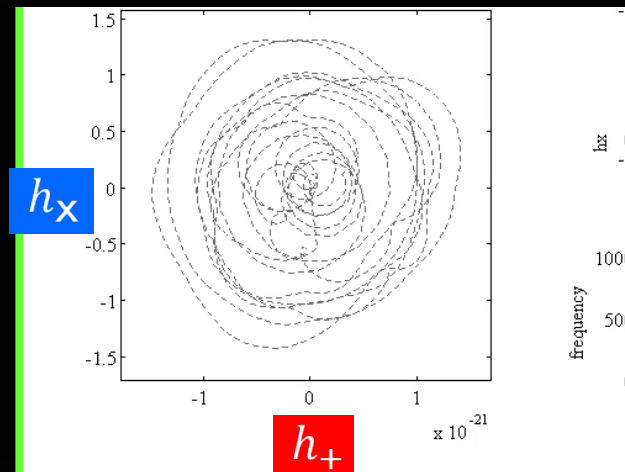
✓ **Low $T/|W|$ instability** is most likely to develop (Ott + (05, ApJL), Scheidegger + (10, A&A))

GW emissivity: $\psi \equiv (A_+ I)^2 + (A_\times I)^2 = (\ddot{I}_{xx} - \ddot{I}_{yy})^2 + (2\ddot{I}_{xy})^2$

Circular polarization of h_+ and $h_\times$



Strong emission from one-armed spiral wave

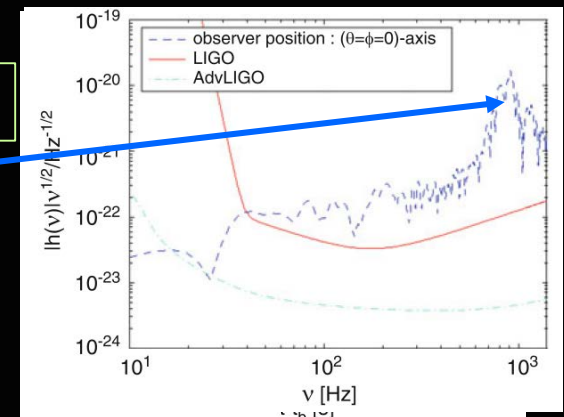


Hayama et al. (2016), PRL
(see also Klimentenko et al. (2015) PRD)

✓ GW from non-axisym. instabilities (incl. low $T/|W|$, spiral SASI) : **Quasi-Periodicity**
(Ott + (07, PRL), Scheidegger + (10, A&A), Kuroda + (14, PRD))

⇒ The effective amplitude scales as the # of GW cycles as

$$h_{\text{eff}} \propto h \sqrt{N}$$



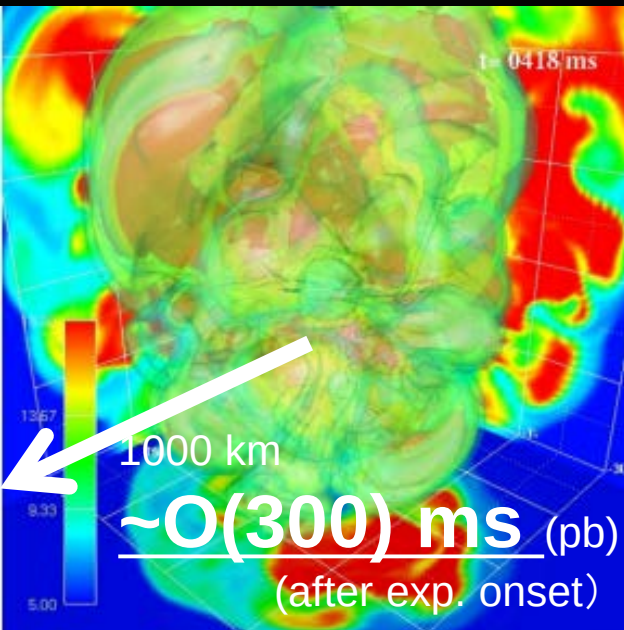
- ✓ Circular polarization can be **evidence of “rapid rotation”**.
- ✓ **“Quasi-periodicity”** enhances the chance of detection.

Summary

	Neutrino mechanism	MHD mechanism
Progenitor	Non- or slowing- rotating star ($\Omega_0 < \sim 0.1$ rad/s)	Rapidly rotating star with strong B fields ($\Omega_0 > \sim \pi$ rad/s, $B_0 > \sim 10^{11}$ G)
Main GW signatures	Three generic phases: Prompt convection, neutrino-driven convection & SASI, and explosion	Rotating bounce (< 20 ms p.b.) and non-axisymmetric instabilities (< ? ms)
Detection Prospect	✓ Requires 3rd generation detector to see every Galactic event (with high SNR). ✓ Closeby events (2~3kpc) detectable, LIGO/Virgo/KAGRA ✓ If detected, critical information about SN engine (convection-dominant vs. SASI dominant) can be obtained. ✓ Detection of circular polarization: important probe of SASI.	✓ Bounce GW signal: detection horizon of LIGO, depending on Ω_0, can cover our Milky way and beyond. ✓ GWs from non-axisymmetric instabilities: “quasi-periodicity” enhances chance of detection. ✓ Detection of circular polarization: important probe of core rotation.

Next 10 years: Where are we and are we going?

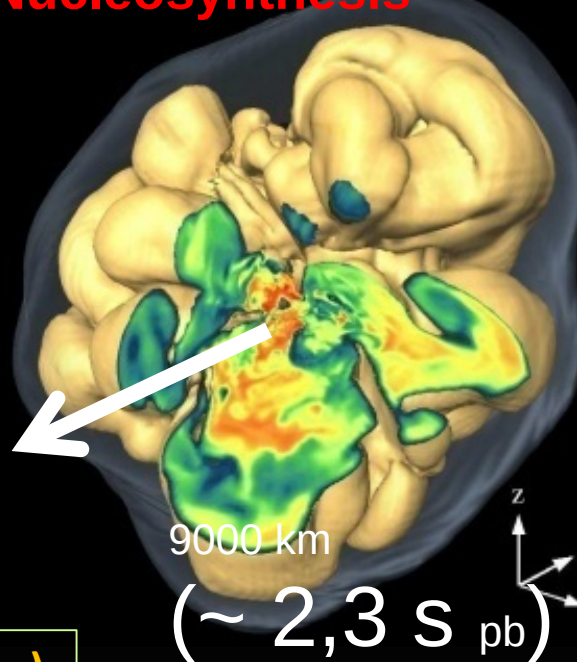
“A” self-consistent 3D model



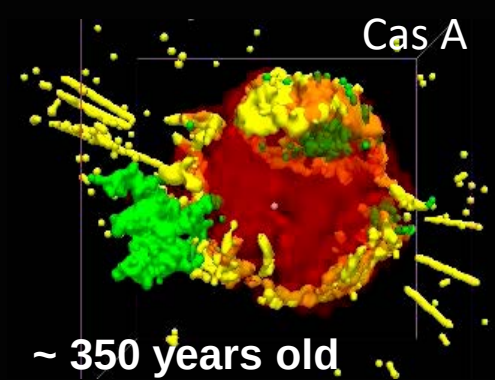
Melson+15, Takiwaki+16, Ott+18

**For progenitors ($11.2, 15, 20, 27 M_{\text{sun}}$),
the stalled shock revived !
(5D/4D with approximate transport)**

Gray-transport simulation **Nucleosynthesis**

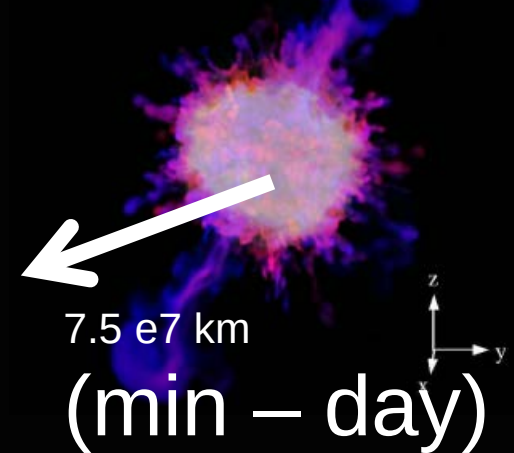


Wongwathanarat et al. (2015)



DeLaney et al. (2010)

Hydrodynamic model:
Mixing, RT, RM instabilities



Wongwathanarat et al. (2016)

**To-do-1: Long-term evolution in self-consistent 3D (GR) models
 $\Rightarrow$ confront CCSN theory with observation $\Rightarrow$ Pragmatism**

To-do-2 : Full GR and Boltzmann project :

$\Rightarrow$ ultimately test whether the stalled shock would revive. $\Rightarrow$ Perfectionism

SN 20xx ! in the Galactic center: End-to-End Bridging Simulations

sec min hours day years
 - 4 - 3 - 2 - 1 0 2 >3

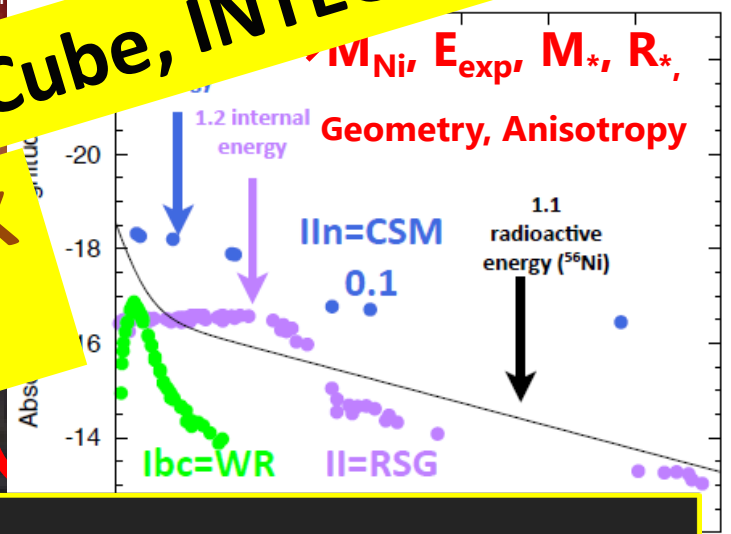
SK detects ~ 10,000 neutrinos

< 15min SURGE meeting (Supernova Urgent Response Group)

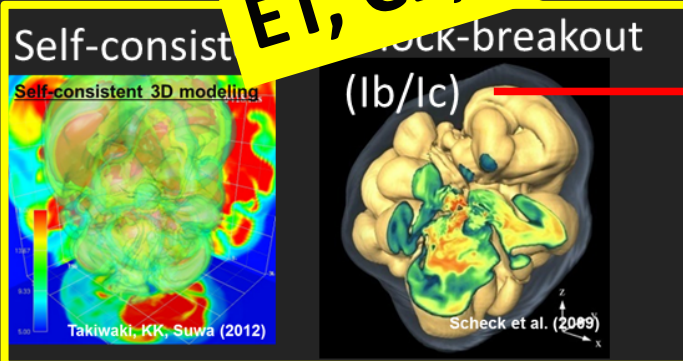
< 1 hour SK provide alert: Astronomers begin observations (onset of neutrino emission)



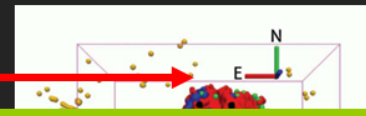
LIGO, Virgo, SK, KamLAND, K, IceCube, INTEGRAL
 KAGRA, Hyper-K, and Post-K
 ET, CE, DUNE.....



Self-Consistent Core-collapse



(IIp)



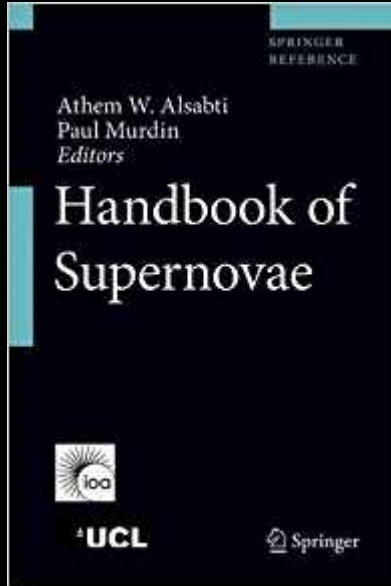
Multi-messenger research
 in steady progress !

Hammer et al. (2011)

~ 350 years old

Useful references

1. Review on GW signatures from CCSNe



~1000 \$!

Chapter 7

“Gravitational waves
From Core-Collapse
Supernovae”

By Kotake and Kuroda



Contents lists available at SciVerse ScienceDirect

Comptes Rendus Physique

www.sciencedirect.com



Gravitational waves / Ondes gravitationnelles

Multiple physical elements to determine the gravitational-wave signatures of core-collapse supernovae

Éléments physiques multiples déterminant les signatures des ondes gravitationnelles de supernovas à effondrement de cœur

Kei Kotake ^{a,b,*}

TOPICAL REVIEW

The gravitational-wave signature of core-collapse supernovae

Christian D Ott

Published 23 February 2009 • 2009 IOP Publishing Ltd

Classical and Quantum Gravity, Volume 26, Number 6

2. Recent publications on CCSN GWs

Summary of publication lists (by Ewald Mueller):

https://wwwmpa.mpa-garching.mpg.de/rel_hydro/GWlit_catalog.shtml

Asteroseismology:

Morozova et al. (2018), ApJ

Torres-Forne et al. (2019), MNRAS

Sotani et al. (2017), PRD