

Conserved charge fluctuations in the chiral limit

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(HotQCD Collaboration)

Criticality in QCD and the Hadron Resonance Gas



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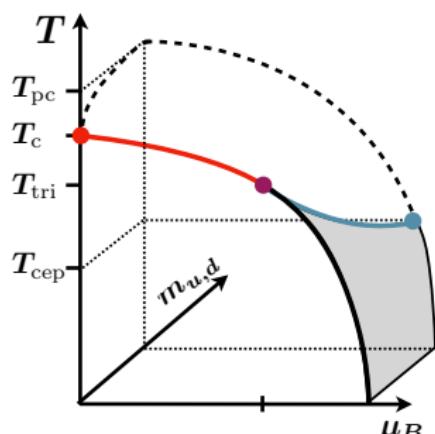
Fakultät für Physik



July 30, 2020

Chiral phase transition in QCD phase diagram

Chiral phase transition in the limit of vanishing up and down quark masses



[F. Karsch, arxiv:1905.03936]

Expected to belong to the universality class of $3d\ O(4)$ spin model

[R.D. Pisarski, F. Wilczek, PRD 29 338 (1984)]

- ⇒ Determination of the order and nature of the chiral phase transition
- ⇒ Imprint of the criticality on the thermodynamics at physical quark masses

$$T_c = 132^{+3}_{-6} \text{ MeV}$$

Anirban's &
Christian's talk

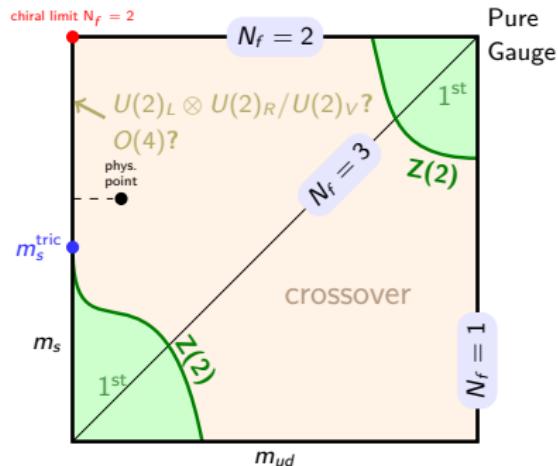
$$T_{pc} = 156.5(1.5) \text{ MeV}$$

[HotQCD, PRL 123, 062002 (2019)]

[HotQCD, PLB 795 (2019) 15–21]

Exploring the chiral phase transition

- measurements towards the chiral limit with strange mass fixed at physical value
- investigate behavior of conserved charge fluctuations, which are also measurable in HIC, in the chiral limit
- alternative scenario : 3d Ising universality transition at non-zero mass, with a first-order chiral limit transition



[O. Philipsen and C. Pinke, PRD93, 114507, 2016]

Lattice Setup

- ⇒ Gauge ensembles generated with HISQ fermion discretization and Symanzik-improved gauge action, used in chiral T_c determination [HotQCD, arxiv:1905.11610].
- ⇒ Ensembles for smaller-than-physical quark (up, down) masses $m_l = ms/27, ms/40, ms/80, ms/160$, keeping strange quark mass m_s fixed at physical value.
- ⇒ Corresp. pion masses : 140 MeV, 110 MeV, 80 MeV, 55 MeV
- ⇒ Thermodynamic and continuum limit not yet performed.
Measurements done at the largest simulated volumes for each mass at fixed time extent $N_\tau = 8$.
- ⇒ Computing resources : Jülich, Piz Daint, JLAB, Bielefeld and Wuhan supercomputing facilities.

Critical behavior of thermodynamic quantities I

Analysis of QCD thermodynamics in the vicinity of the chiral phase transition, acc. to Wilson's RG approach

infinite volume

$$\frac{p}{T^4} = \frac{1}{VT^3} \ln Z(T, \vec{\mu}, m_l) = h^{(2-\alpha)/\beta\delta} f_f(z) + f_r(T, \vec{\mu}, m_l), \quad z \equiv t/h^{1/\beta\delta}$$

singular

regular \rightarrow no linear m_l dependence

$$t = \frac{1}{t_0} \left(\frac{T-T_c}{T_c} + \kappa_2^X \left(\frac{\mu_X}{T} \right)^2 \right)$$

“energy-like” coupling

$$h = \frac{1}{h_0} \frac{m_l}{m_s}$$

X=B,Q,S

“magnetic-like” coupling

critical exponents

Chiral phase transition at
 $m_l \equiv m_u = m_d = 0$ ($h = 0$)
 $T = T_c$ ($t = 0$) at $\mu = 0$

	α	β	δ
$O(4)$	-0.21	0.38	4.82
$O(2)$	-0.017	0.349	4.78
$Z(2)$	+0.109	0.325	4.8

Critical behavior of thermodynamic quantities II

(Anirban's talk)

Conserved charge fluctuations are energy-like w.r.t to chiral phase transition (also Polyakov loop, see David's talk)

magnetic-like

$$\frac{\partial^2 \ln Z}{\partial h^2}$$

$$\sim h^{1/\delta - 1}$$

$$\sim h^{-0.79}$$

divergence : **strong**

mixed

$$\frac{\partial^2 \ln Z}{\partial h \partial t}$$

$$\sim h^{(\beta-1)/\beta\delta}$$

$$\sim h^{-0.34}$$

moderate

energy-like

$$\frac{\partial^2 \ln Z}{\partial t^2}$$

$$\sim h^{-\alpha/\beta\delta}$$

$$\sim h^{+0.11}$$

vanishes

Conserved charge fluctuations at $\mu = 0$ (Singular part) :

$$\chi_{2n}^X = -\left. \frac{\partial^{2n} f/T^4}{\partial (\mu_X/T)^{2n}} \right|_{\mu_X=0} \sim - (2\kappa_2^X)^n h^{(2-\alpha-n)/\beta\delta} f_f^{(n)}(z)$$

Scaling Expectation : $\frac{\partial}{\partial T} \sim \kappa_2^X \frac{\partial^2}{\partial \mu^2}$

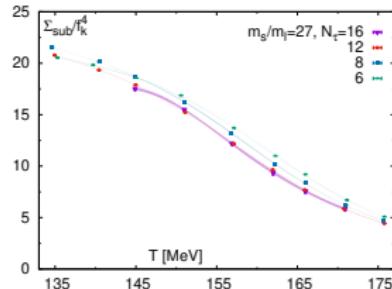
$\Rightarrow \chi_2 \sim \text{Energy density}, \chi_4 \sim \text{Specific heat}$

Scaling at physical quark masses

⇒ Scaling expectation for singular part : $\frac{\partial}{\partial T} \sim \kappa \frac{\partial^2}{\partial \mu^2}$.

Scaling at physical quark masses

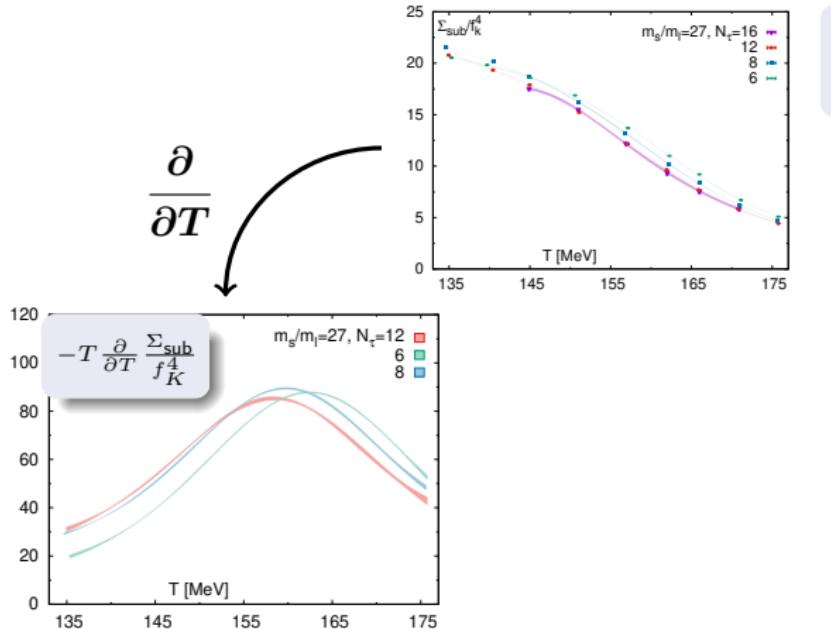
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[HotQCD, PLB 795 (2019) 15]

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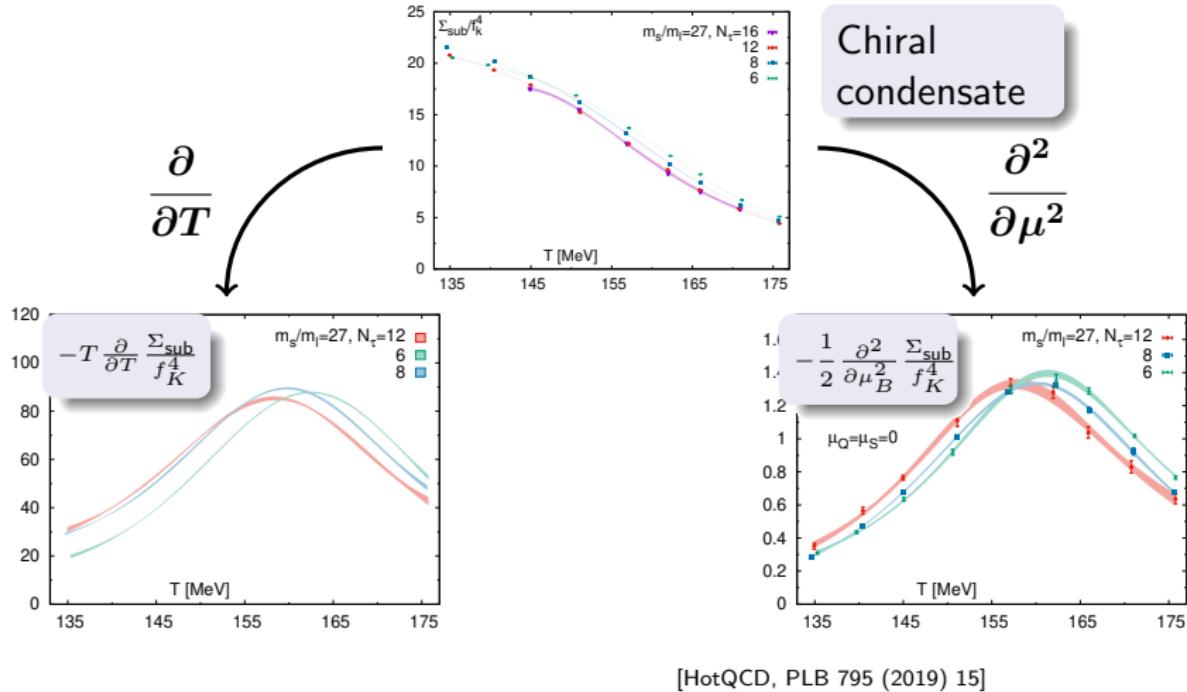


Chiral
condensate

[HotQCD, PLB 795 (2019) 15]

Scaling at physical quark masses

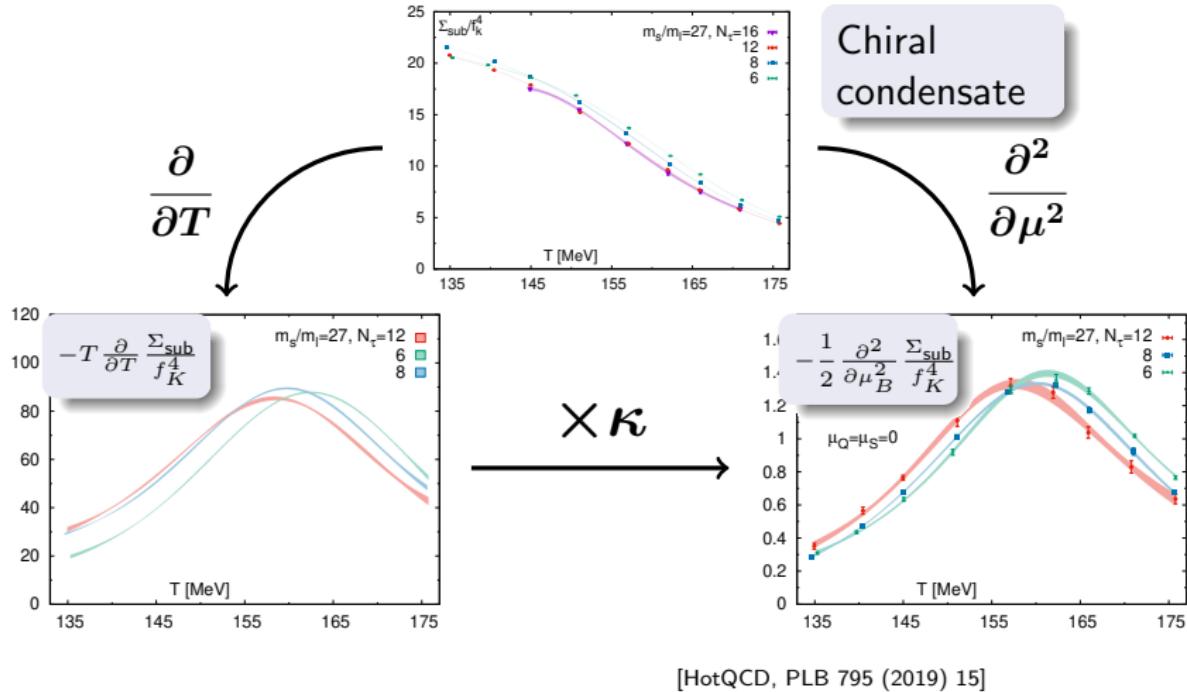
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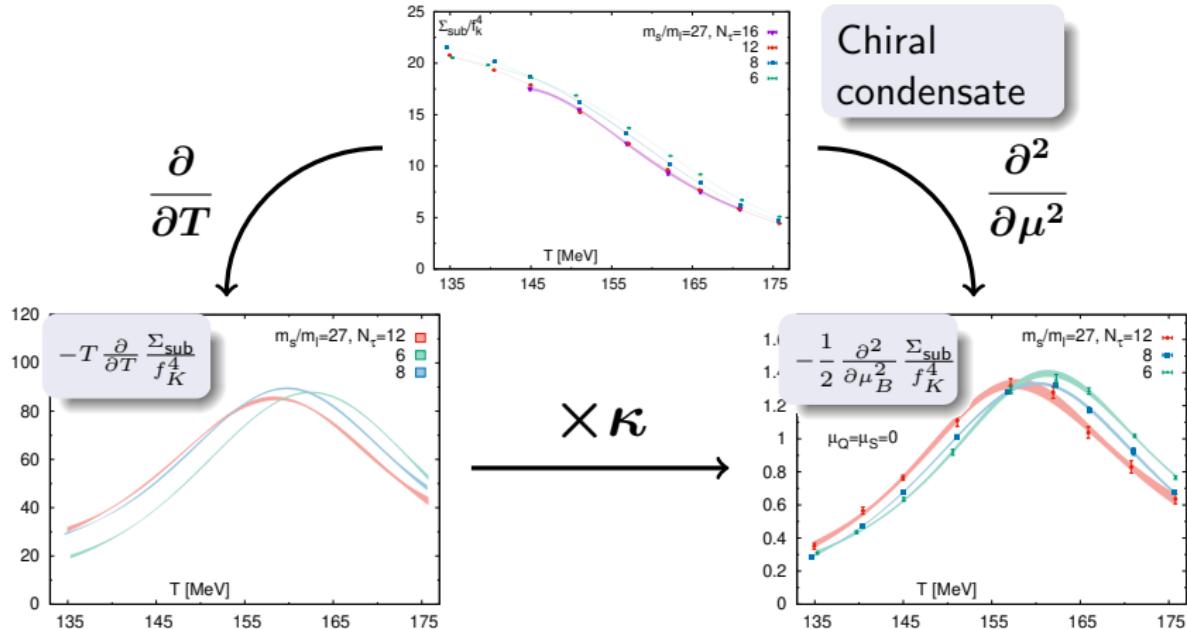
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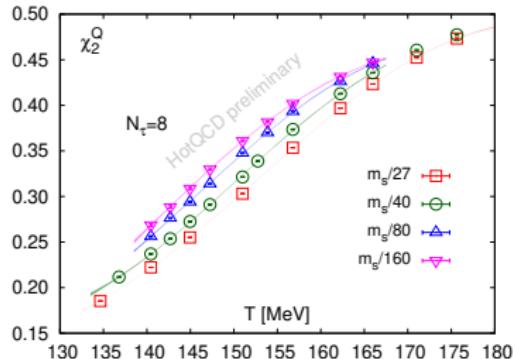
Scaling expectation already holds at physical quark masses

Scaling expectation for χ_2^Q

⇒ Scaling expectation if singular part dominates : $\frac{\partial}{\partial T} \sim \kappa \frac{\partial^2}{\partial \mu^2}$.

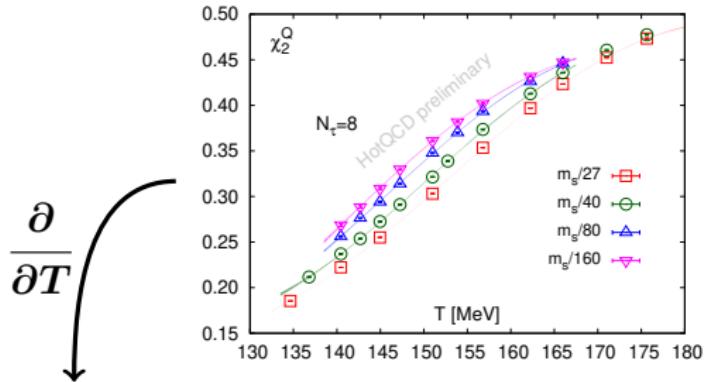
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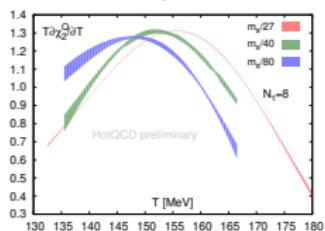


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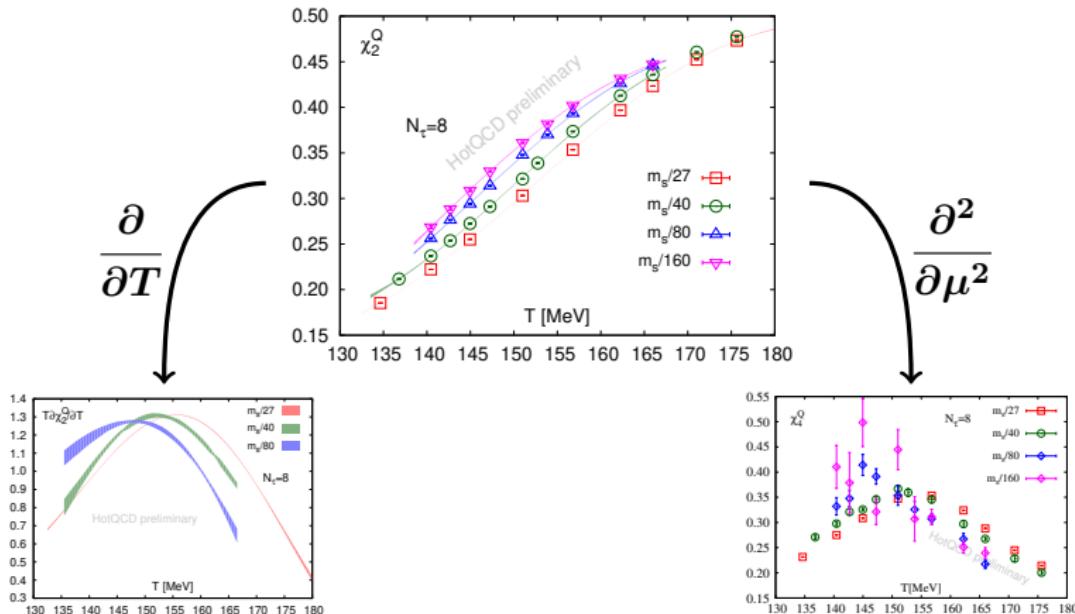


$$\frac{\partial}{\partial T}$$



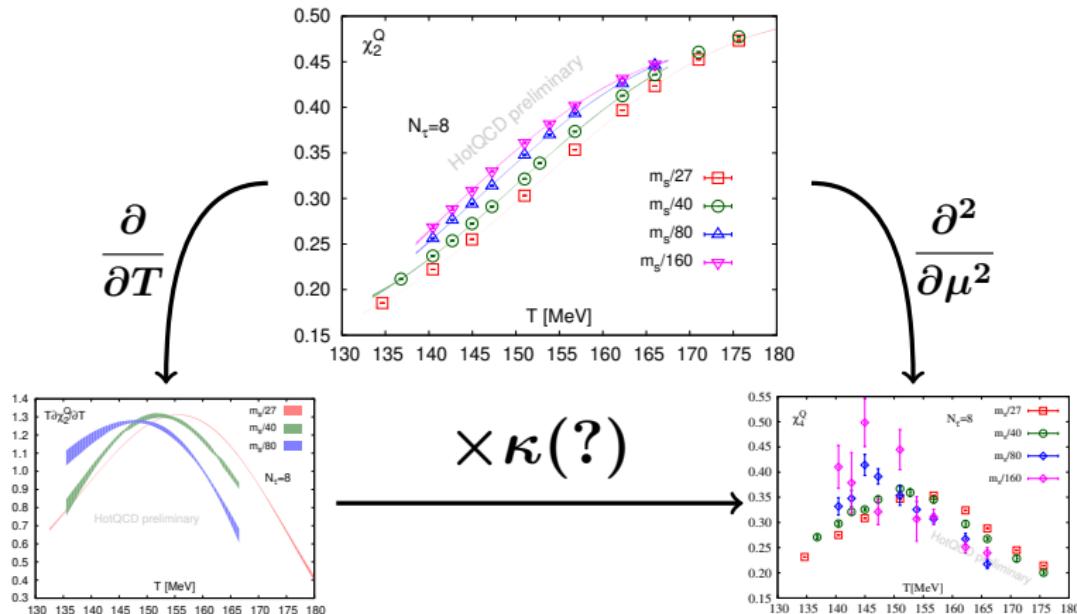
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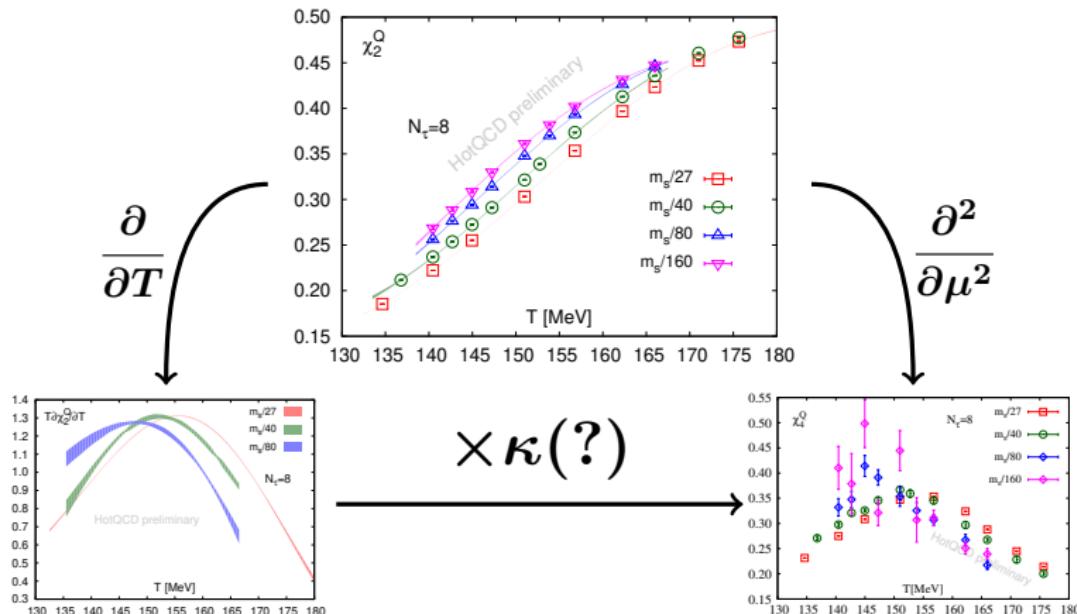
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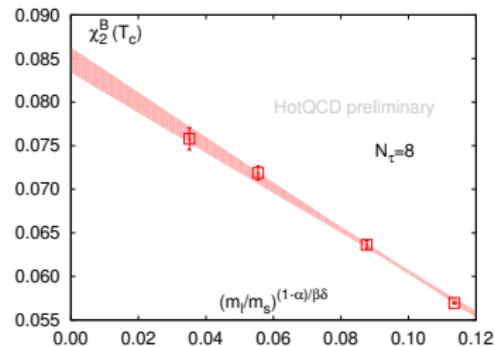
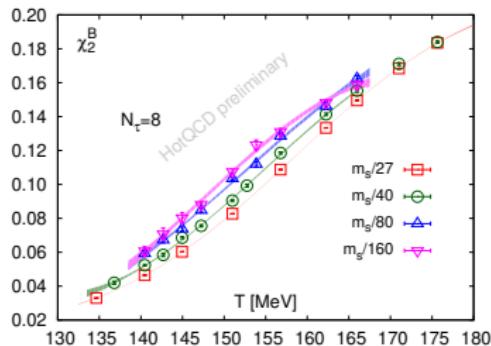
Singular term in χ_2 (energy-like) is not dominant
unlike in order parameter

Estimation of the singular contribution to χ_2^B

$$\chi_2^X(T_c, m_l) \sim -\kappa_2^X h^{(1-\alpha)/\beta\delta} f_f^{(1)}(0) + \text{const. reg. term} + \mathcal{O}(m_l^2)$$

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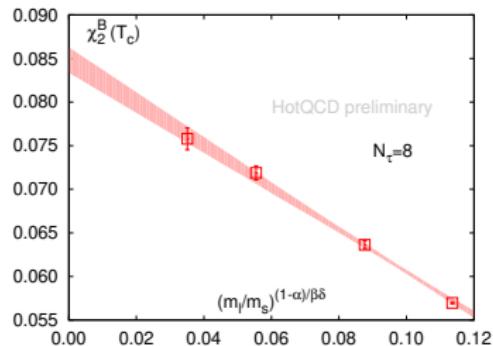
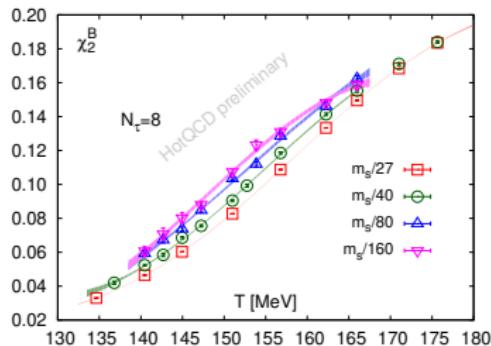


$\chi_2^X(T_c, m_l = 0) - \chi_2^X(T_c, m_l = m_s/27) = \text{Singular part of } \chi_2^X$

► Linear in H plot

Estimation of the singular contribution to χ_2^B

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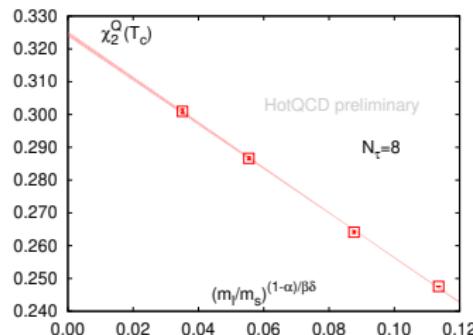
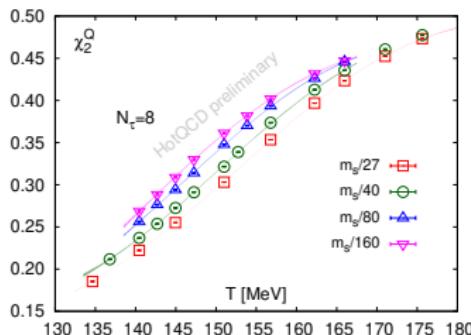


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Singular contribution to χ_2^B at physical masses $\sim 50\%$

Scaling behavior of χ_2^Q

$$\chi_2^X(T_c, m_l) \sim -\kappa_2^X h^{(1-\alpha)/\beta\delta} f_f^{(1)}(0) + \text{const. reg. term} + \mathcal{O}(m_l^2)$$



$$\frac{\text{Singular part of } \chi_2^Q}{\text{Singular part of } \chi_2^B} = \frac{\kappa_2^Q}{\kappa_2^B} \sim 2.6$$

Close to result for physical mass :
 1.81 ± 0.40

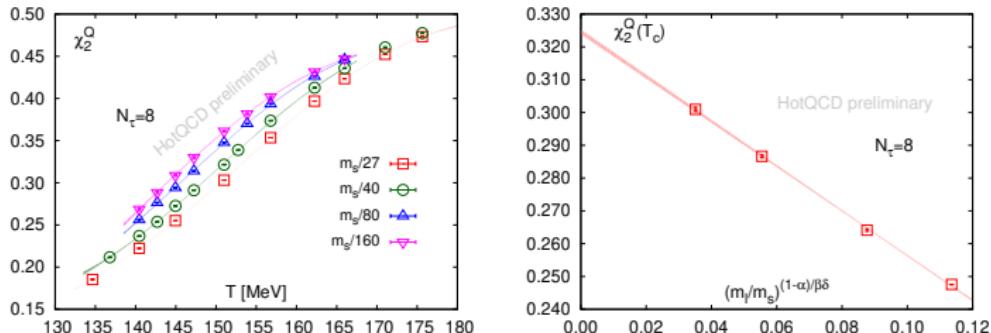
[HotQCD, Phys. Lett. B 795 (2019) 15]

Linear in H plot

Singular contribution to
 χ_2^Q at physical masses
 $\sim 30\%$

Scaling behavior of χ_2^Q

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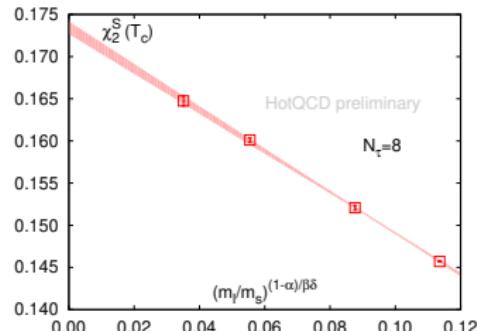
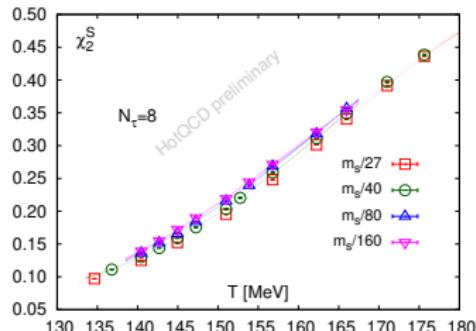
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Linear in H plot

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Scaling behavior of χ_2^S

$$\chi_2^X(T_c, m_l) \sim -\kappa_2^X h^{(1-\alpha)/\beta\delta} f_f^{(1)}(0) + \text{const. reg. term} + \mathcal{O}(m_l^2)$$



Singular contributions in χ_2^B and χ_2^S
are similar.

Consistent with previous result for
physical mass

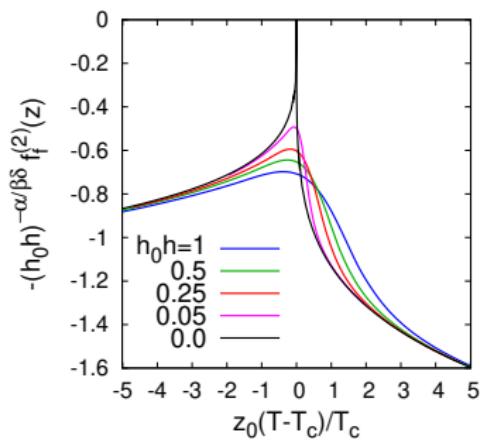
[HotQCD, Phys. Lett. B 795 (2019) 15]

Singular contribution to
 χ_2^S at physical masses
 $\sim 20\%$

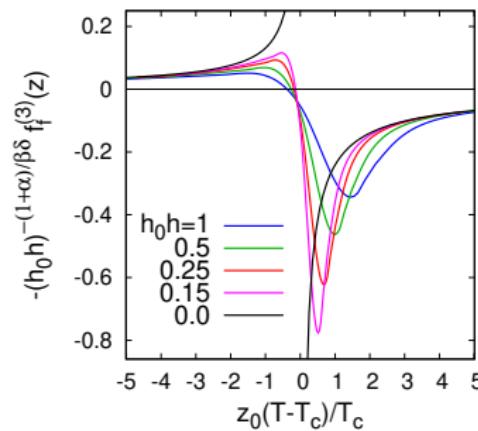
Scaling expectation of higher fluctuations

Derivatives of the singular part for 3d $O(4)$ universality class

$$h^{0.116} f_f^{(2)}$$



$$h^{-0.429} f_f^{(3)}$$



[B. Friman, F. Karsch, K. Redlich, V. Skokov, Eur. Phys. J. C (2011) 71:1694]

$\sim \chi_4$

$\sim \chi_6$

Scaling behavior of χ_4^Q

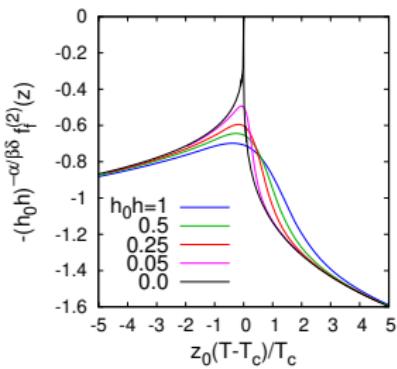
Singular part :

$$\chi_4^Q \sim h^{-\alpha/\beta\delta} f_f^{(2)}(z)$$

$$O(4) : -\alpha/\beta\delta = 0.116$$

Not divergent

– but pronounced spike



Scaling behavior of χ_4^Q

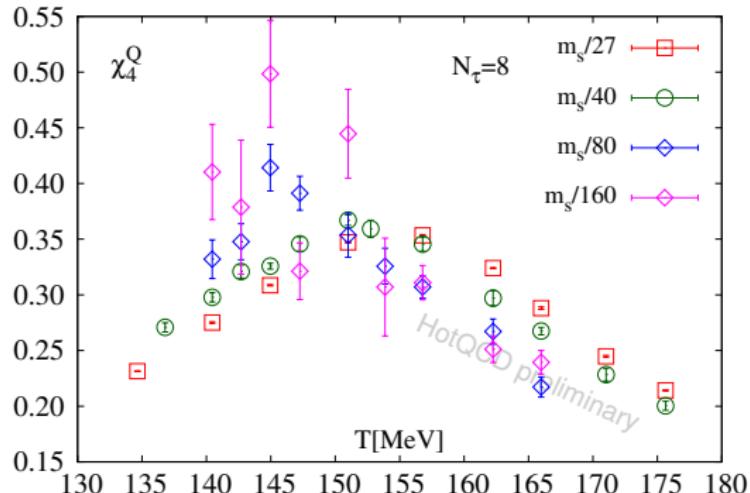
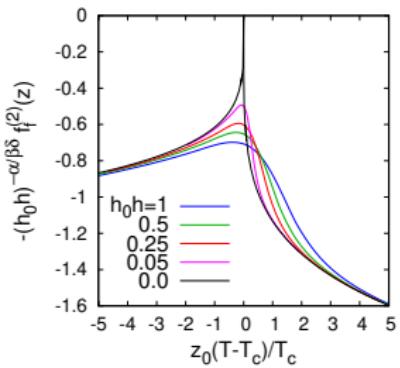
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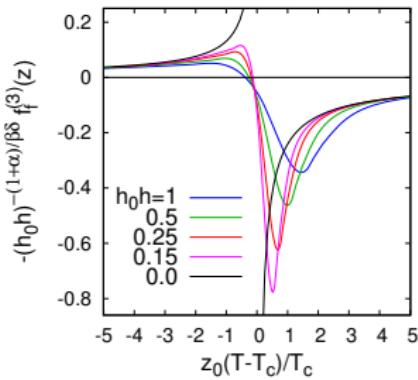
⇒ Expected features are apparent for
 $m_l = m_s/27, m_s/40$

Scaling behavior of χ_6^Q

Singular part :

$$\begin{aligned}\chi_6^Q &\sim h^{-(1+\alpha)/\beta\delta} f_f^{(3)}(z) \\ &\sim h^{-0.429}\end{aligned}$$

Moderate divergence

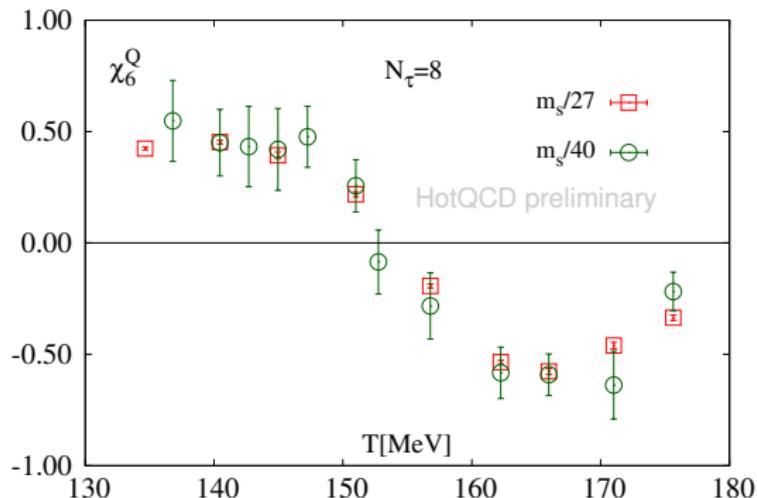
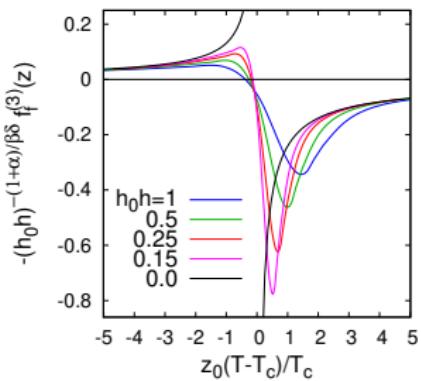


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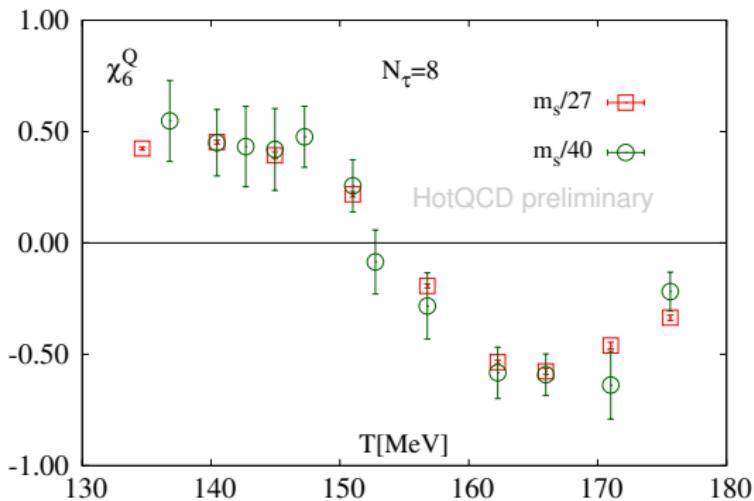
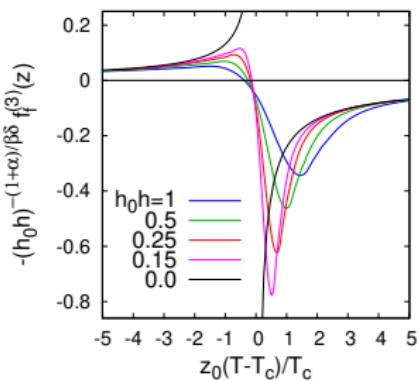


Scaling behavior of χ_6^Q

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Moderate divergence



Ratio of peak heights expected from scaling : $(\chi_6^Q)_{1/40}^{max}/(\chi_6^Q)_{1/27}^{max} \sim 1.18$

Mixed observables I

H -derivatives of energy-like observables

$$\chi_2^S \sim -\kappa_2^S H^{(1-\alpha)/\beta\delta} f'_f(z) + f_{\text{reg}}$$

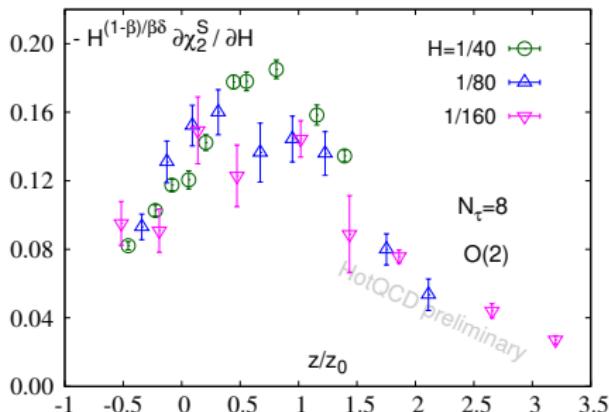
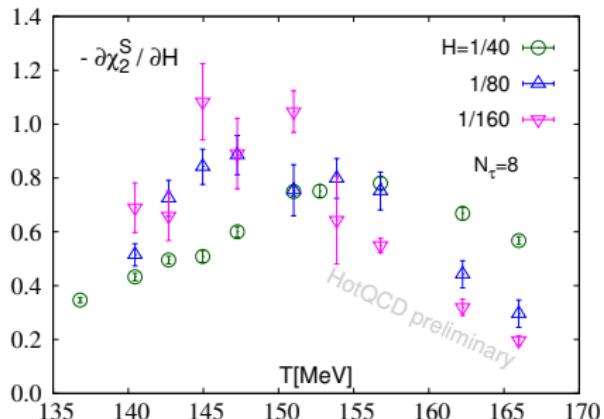
$$H = \frac{m_L}{m_s}$$

$$z \equiv z_0 \frac{t}{H^{1/\beta\delta}}$$

$$\frac{\partial \chi_2^S}{\partial H} \sim \kappa_2^S H^{(\beta-1)/\beta\delta} f'_G(z) + \frac{\partial f_{\text{reg}}}{\partial H}$$

$$\frac{\beta-1}{\beta\delta} = -0.39, O(2)$$

Divergent already for 2nd order



Dominant singular part

Mixed observables II

Strange chiral condensate \Rightarrow energy-like observable

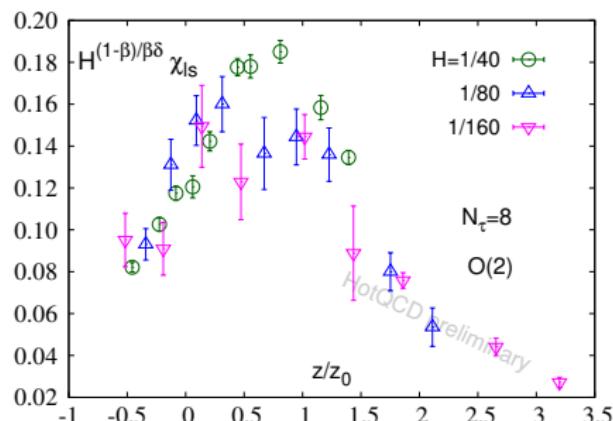
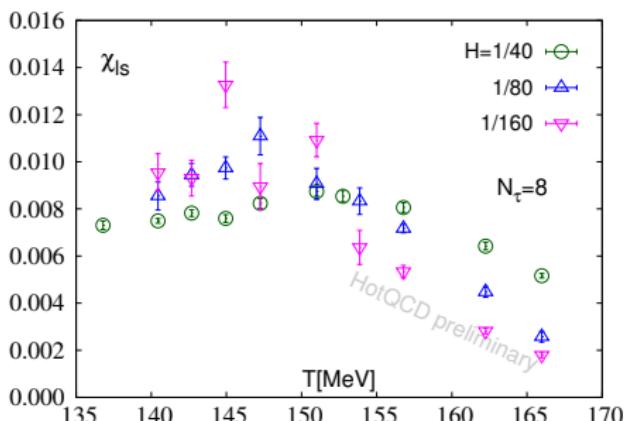
$$\langle \bar{\psi} \psi \rangle_s = AH^{(1-\alpha)/\beta\delta} f'_f(z) + f_{\text{reg}}$$

$$H = \frac{m_L}{m_S}$$

$$\chi_{ls} \equiv \frac{\partial \langle \bar{\psi} \psi \rangle_s}{\partial H} = -AH^{(\beta-1)/\beta\delta} f'_G(z) + \frac{\partial f_{\text{reg}}}{\partial H}$$

$$\frac{\beta-1}{\beta\delta} = -0.39, O(2)$$

Divergent at 2nd order



Similar divergence, different regular contribution

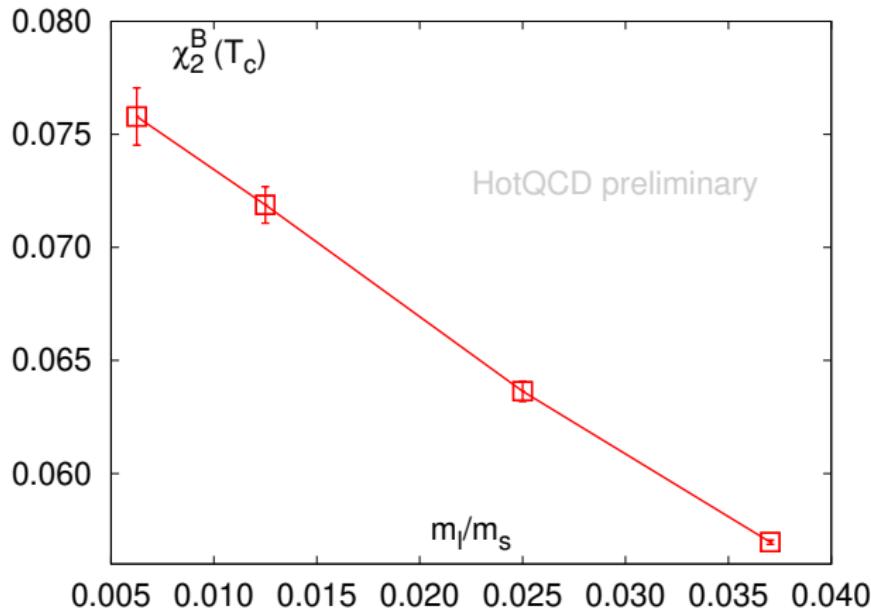
Conclusions and Outlook

- ⇒ Fluctuations of conserved charges seem consistent with chiral phase transition belonging to $O(4)$ or $O(2)$ universality class. Expected energy-like behavior of fluctuations w.r.t. chiral phase transition.
- ⇒ Singular part can be extracted from χ_2^X and may be used to determine the curvature coefficients of the chiral critical line. Singular contribution to 2nd order fluctuations roughly $\leq 50\%$
- ⇒ Strange fluctuations and strange quark condensate behaves as energy-like quantities in the 2-flavour chiral limit.
- ⇒ Future comparison with HRG with smaller masses obtained from ChPT.
- ⇒ Requires more statistics at lower masses and proper continuum and thermodynamic limits.

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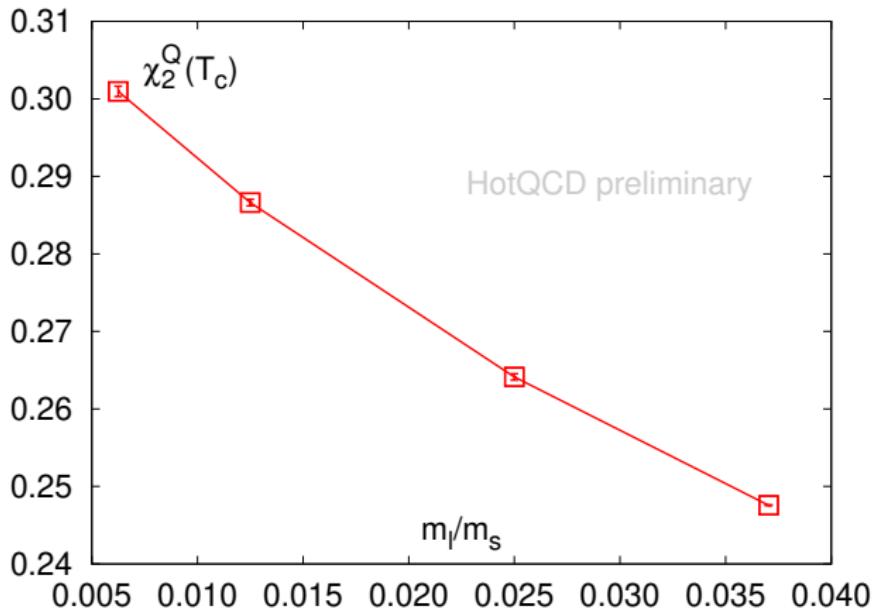
Thank you for your attention



Points are connected to guide the eye. Lines are NOT fits.

[◀ Go Back](#)

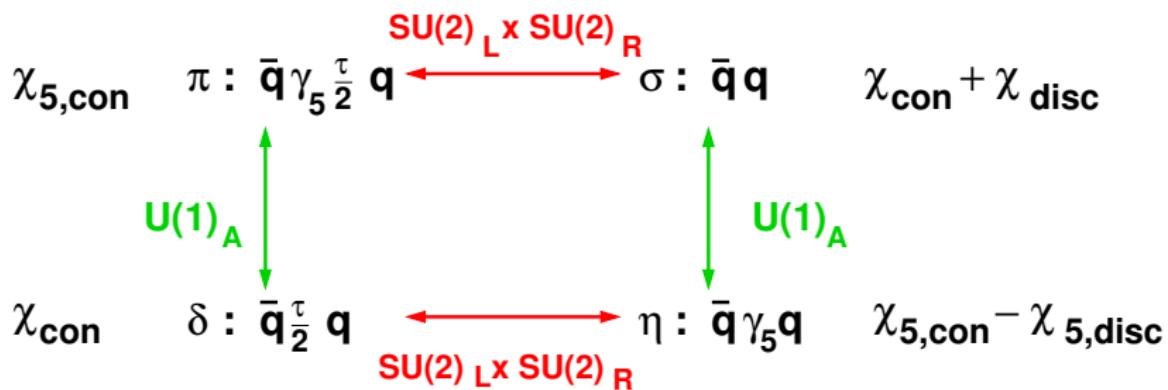
Backup



Points are connected to guide the eye. Lines are NOT fits.

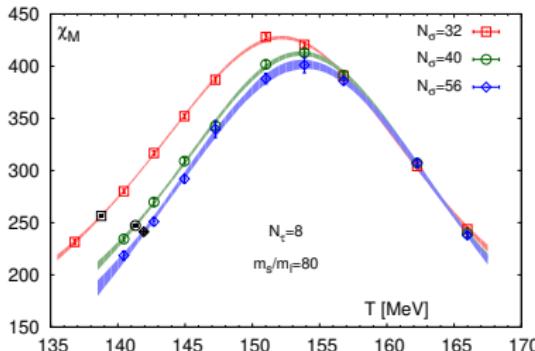
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Symmetry transformations



$O(4)$ criticality : Bounds on first order transition

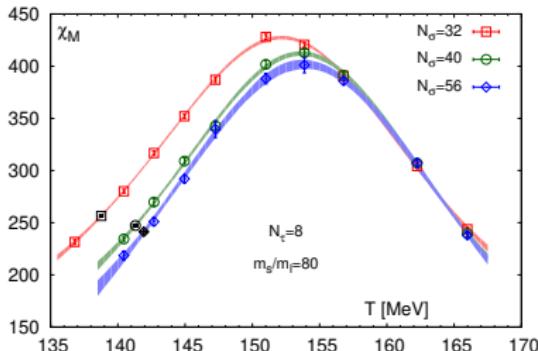
[HotQCD, arxiv:1905.11610]



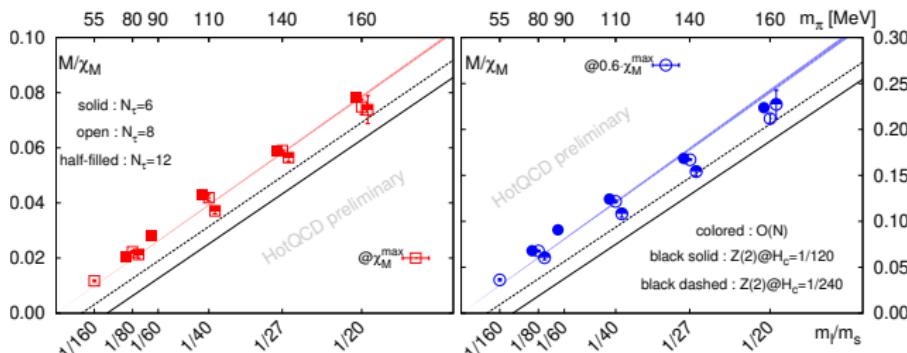
- ⇒ No indications of first order volume scaling at $m_\pi = 80$ MeV
- ⇒ Consistent with $O(4)$ volume scaling at non-zero H

$O(4)$ criticality : Bounds on first order transition

[HotQCD, arxiv:1905.11610]



- ⇒ No indications of first order volume scaling at $m_\pi = 80$ MeV
- ⇒ Consistent with $O(4)$ volume scaling at non-zero H



$$\frac{M(T_X)}{\chi_M(T_X)} = (H - H_c) \frac{f_G(z_X)}{f_X(z_X)}$$

Seems no indication of a 1st order transition till $m_\pi \sim 46$ MeV

[H.T. Ding et al, Nucl. Phys. A 982 (2019) 211]

$O(4)$ criticality : Chiral T_c estimation

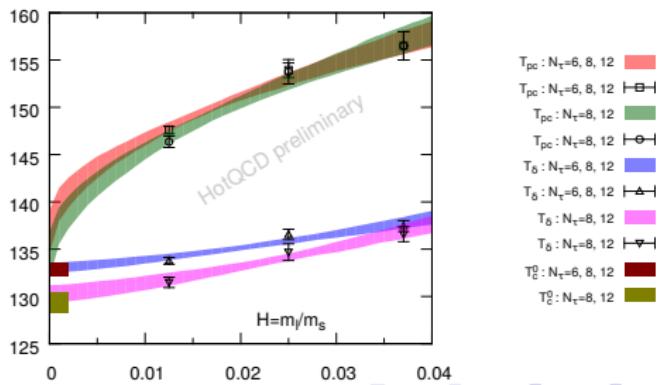
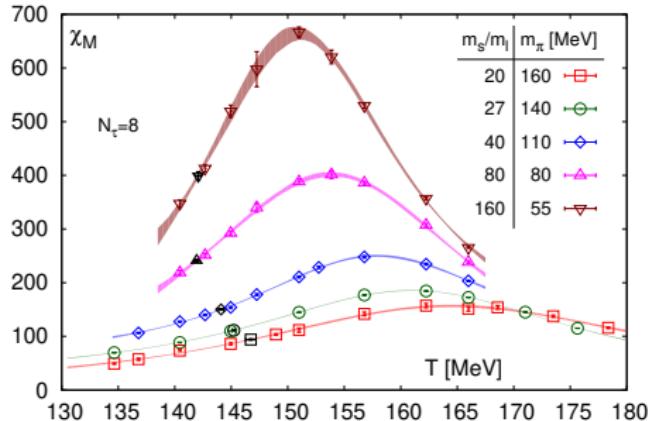
Chiral phase transition at the chiral limit extracted by $O(4)$ scaling arguments.

$$T_X(H) = T_c \left(1 + \frac{z_X}{z_0} H^{1/\beta\delta} \right)$$

$$X = \delta, 60, pc$$

$$\frac{H\chi_M(T_\delta, H)}{M(T_\delta, H)} = \frac{1}{\delta}$$

$$\chi_M(T_{60}, H) = 0.6 \chi_M^{\max}$$



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$$T_c = 132^{+3}_{-6} \text{ MeV}$$

[HotQCD, arxiv:1905.11610]

