

*Constraining the critical temperature for color
superconductivity in heavy-ion collisions from
neutron star phenomenology*

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(in collaboration with
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OUTLINE

- Introduction
- Equation of State (EoS)
- Constructing a hybrid EoS I
- Preliminary results I
- Nambu-Jona-Lasinio (NJL) model
- Constructing a hybrid EoS II
- Preliminary results II
- Conclusion

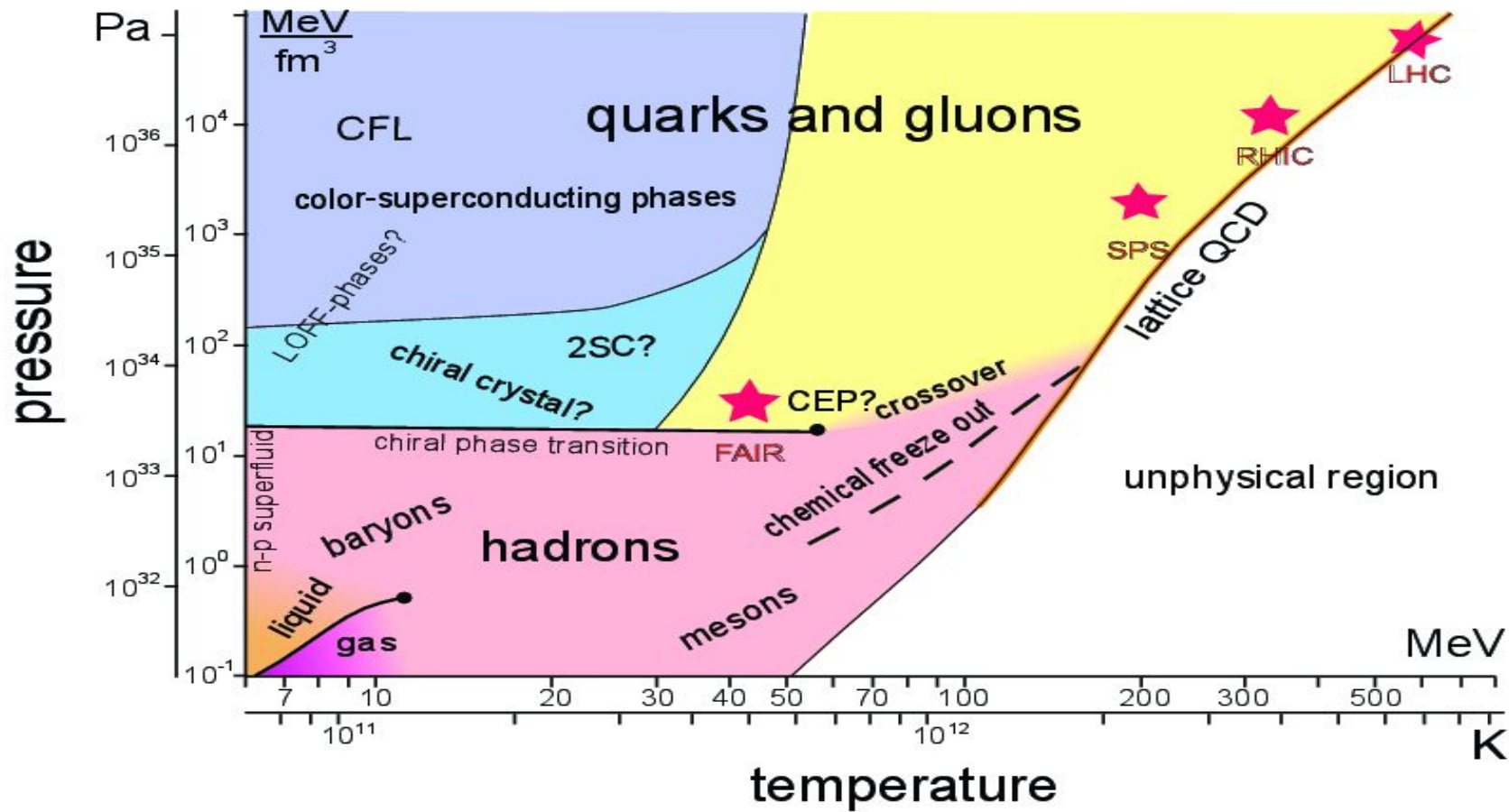
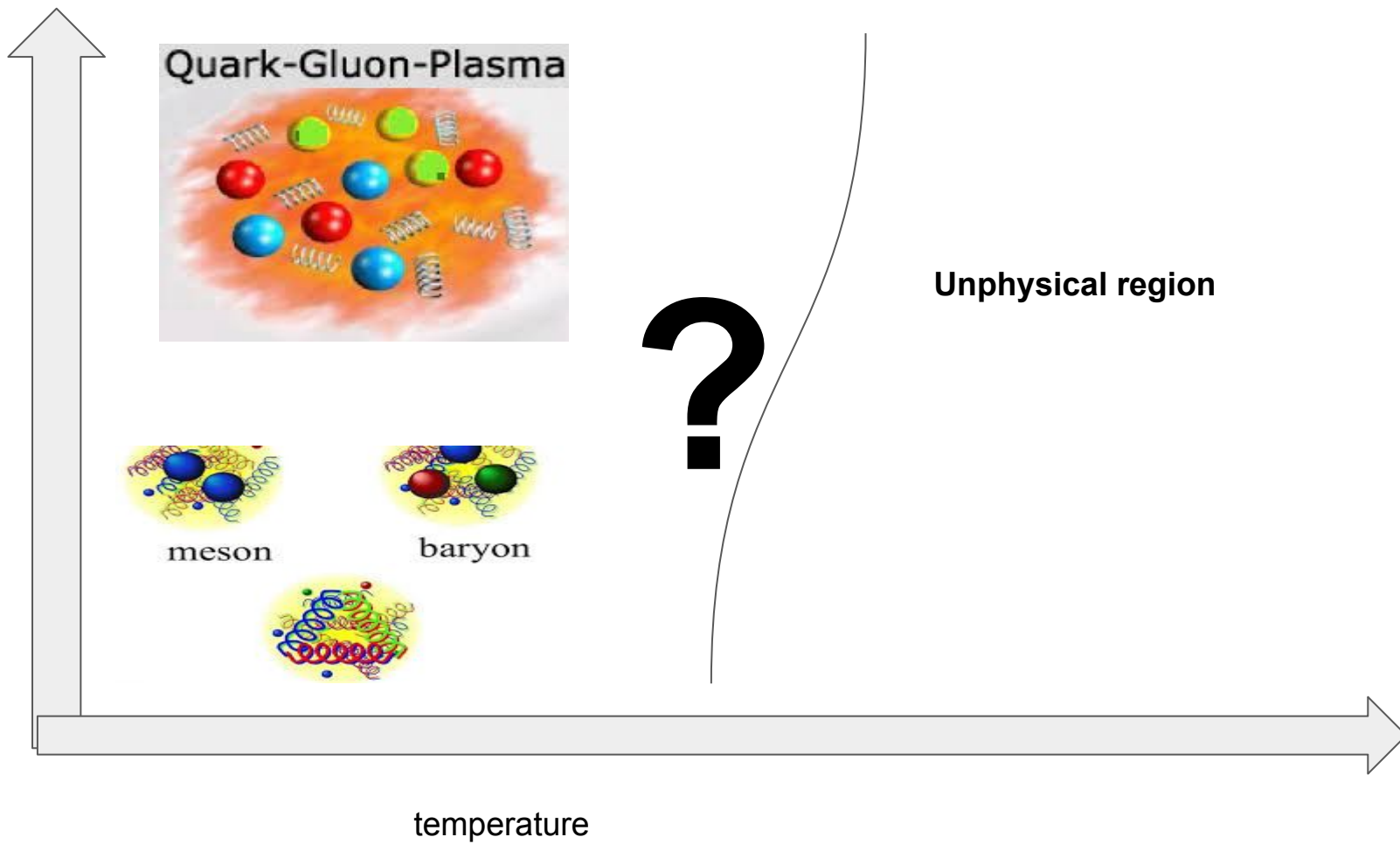


Figure courtesy: K. Heckmann (TU Darmstadt, 2012)

OPEN QUESTION




ADOPTED
PRESCRIPTION



- M.Albright, J. Kapusta, C. Young, Phys. Rev. C 90, 024915 (2014)
- J. Kapusta, T. Welle, Phs. Rev. C 104, L012801 (2021)

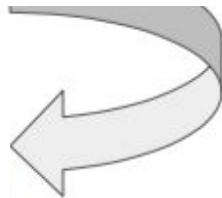
SWITCHING FUNCTION
B/W
HADRONS & QGP


$$S(\mu) = \exp\left\{-\left[\left(\frac{T}{T_0}\right)^r + \left(\frac{\mu}{\mu_0}\right)^r\right]^{-1}\right\}$$

Unified EoS...?

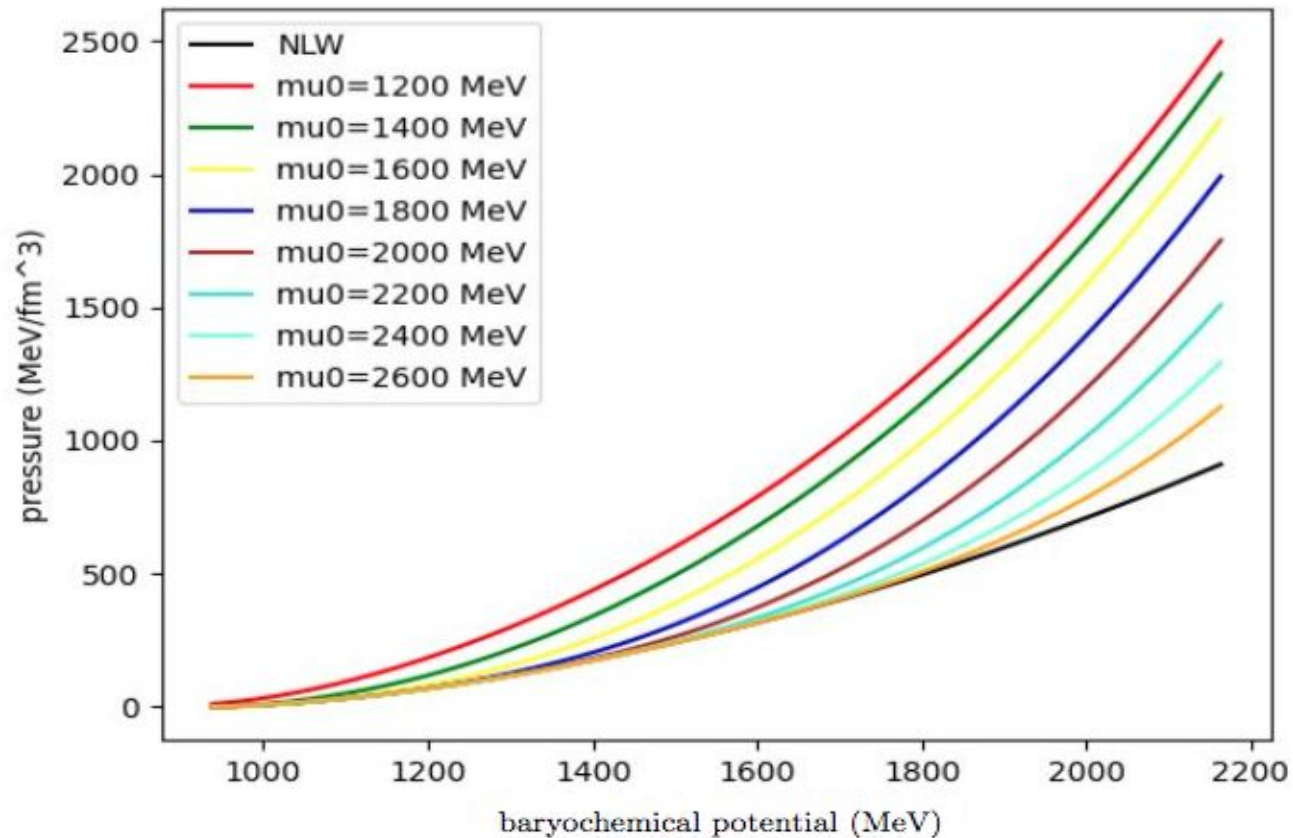


$$P(\mu, T) = S(\mu, T)P_q(\mu, T) + [1 - S(\mu, T)]P_h(\mu, T)$$



PRELIMINARY RESULTS - I

Direct NLW-CFL crossover with diquark gap = 200 MeV



NLW = Non-linear Walecka

Phenomenological EoS

pQCD

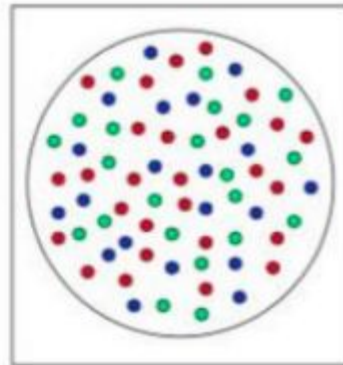
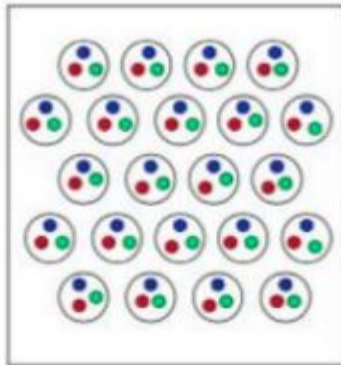
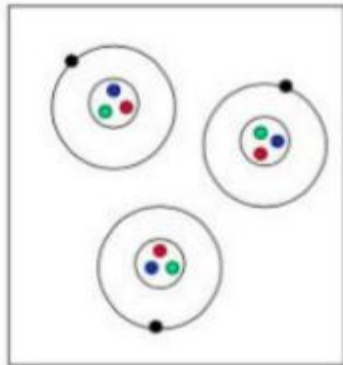
Cooper
pairing
<qq>

Quark
confinement

$$P_q(\mu) = \frac{3}{4\pi^2} a_4 \left(\frac{\mu}{3}\right)^4 + \frac{3}{\pi^2} \Delta^2 \left(\frac{\mu}{3}\right)^2 - B_{\text{eff}}$$

C. Zhang, R. Mann, Phys.Rev.D
103 (2021), 6

M. Alford, M. Braby, M. Paris,
S.Reddy, *Astrophys.J.* 629 (2005)
969-978



$$V_e(r) \sim \left\{ \frac{e_0^2}{r} \right\}$$

$$V_q(r) \sim r$$

Confinement?
Asymptotic
freedom

Lagrangian...
Symmetries?

Degrees of
freedom---> quarks,
diquarks, multi-quark
clusters...?

Book: The Physics of the Quark-Gluon Plasma, Edited by Sourav Sarkar, Helmut Satz, Bikash Sinha

Nambu-Jona-Lasinio (SU(2))

(with vector meson & diquark interactions)

Given (full) Lagrangian -

$$\mathcal{L} = \mathcal{L}_0 + \mathcal{L}_S + \mathcal{L}_V + \mathcal{L}_D$$

$$\mathcal{L}_0 = \bar{q}(i\not{\partial} - m_0 + \mu\gamma_0)q$$

$$\mathcal{L}_S = G_S[(\bar{q}q)^2 + (\bar{q}i\gamma_5\vec{\tau}q)^2]$$

$$\mathcal{L}_V = -G_V(\bar{q}\gamma_\mu q)^2$$

$$\mathcal{L}_D = G_D \sum_{A=2,5,7} (\bar{q}i\gamma_5\tau_2\lambda_A q^C)(\bar{q}^C i\gamma_5\tau_2\lambda_A q)$$

D. Blaschke, M. Buballa,
et. al,
Annals Phys. 348 (2014)
228-255

Partition function -

$$\mathcal{Z} = \int \mathcal{D}q \mathcal{D}\bar{q} \exp \left(\int d^4x_E \mathcal{L} \right)$$

Introducing Nambu-Gorkov bispinors $\Rightarrow \Psi \equiv \frac{1}{\sqrt{2}} \begin{pmatrix} q \\ q^C \end{pmatrix}$, and, $\bar{\Psi} \equiv \frac{1}{\sqrt{2}} (\bar{q} \quad \bar{q}^C)$

And the charge-conjugation of a Dirac spinor (in position space) as -

$$q_c(x) \equiv C \bar{q}^T(x)$$

where, $C \equiv i\gamma^2\gamma^0$

THERMODYNAMIC POTENTIAL



$$\Omega(T, \mu) = -\frac{T}{V} \ln \mathcal{Z}(T, \mu),$$

$$\Omega^{MFA} = -\frac{T}{V} \ln \mathcal{Z}^{MFA} = \frac{\bar{\sigma}^2}{2G_S} + \frac{\bar{\Delta}^2}{2G_D} - \frac{\bar{\omega}^2}{2G_V} - 2 \int \frac{d^3 \vec{p}}{(2\pi)^3} \xi(\vec{p}),$$

where

$$\xi(\vec{p}) = \sum_{\kappa, s=\pm} 2 \left\{ \epsilon_r^\kappa / 2 + T \ln \left[1 + e^{-\frac{\epsilon_r^\kappa + s \delta \tilde{\mu}_r}{T}} \right] \right\} + \sum_{\kappa, s=\pm} \left\{ \bar{E}_b^\kappa / 2 + T \ln \left[1 + e^{-\frac{\bar{E}_b^\kappa + s \delta \tilde{\mu}_b}{T}} \right] \right\}$$

Gap equations

$$\begin{aligned}\bar{\sigma} &= 2G_s \int \frac{d^3\vec{p}}{(2\pi)^3} \frac{g(\vec{p})}{E} M(\vec{p}) \\ &\times \sum_{\kappa} \left\{ 1 - n_F\left(\frac{\bar{E}_b^{\kappa} + \delta\tilde{\mu}_b}{T}\right) - n_F\left(\frac{\bar{E}_b^{\kappa} - \delta\tilde{\mu}_b}{T}\right) \right. \\ &\left. + \frac{2\bar{E}_r^{\kappa}}{\epsilon_r^{\kappa}} \left[1 - n_F\left(\frac{\epsilon_r^{\kappa} + \delta\tilde{\mu}_r}{T}\right) - n_F\left(\frac{\epsilon_r^{\kappa} - \delta\tilde{\mu}_r}{T}\right) \right] \right\},\end{aligned}$$

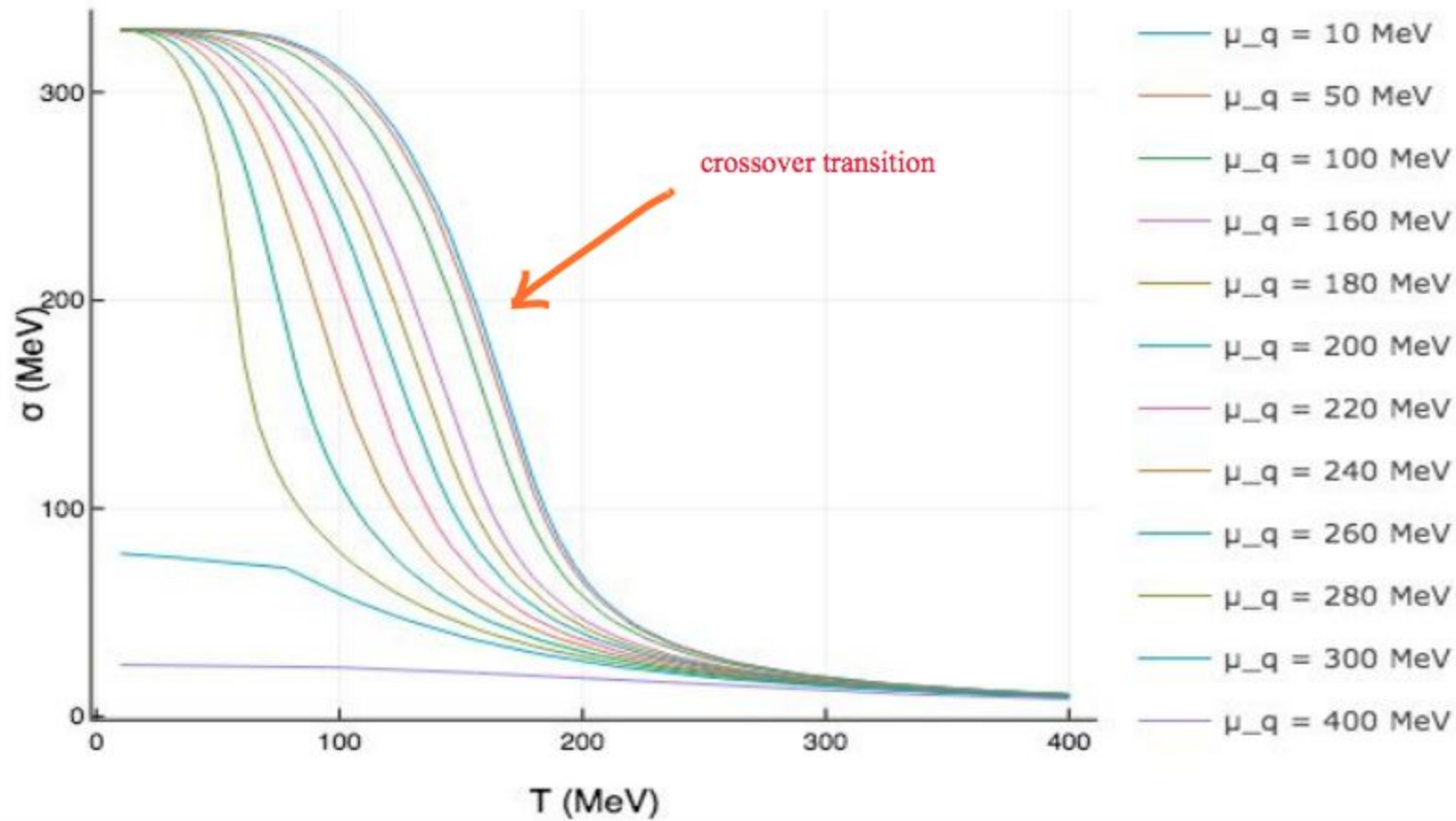
$$\begin{aligned}\bar{\Delta} &= 2G_D \bar{\Delta} \int \frac{d^3\vec{p}}{(2\pi)^3} g^2(\vec{p}) \\ &\times \sum_{\kappa=\pm} \left\{ \frac{2}{\epsilon_r^{\kappa}} \left[1 - n_F\left(\frac{\epsilon_r^{\kappa} + \delta\tilde{\mu}_r}{T}\right) - n_F\left(\frac{\epsilon_r^{\kappa} - \delta\tilde{\mu}_r}{T}\right) \right] \right\},\end{aligned}$$

and

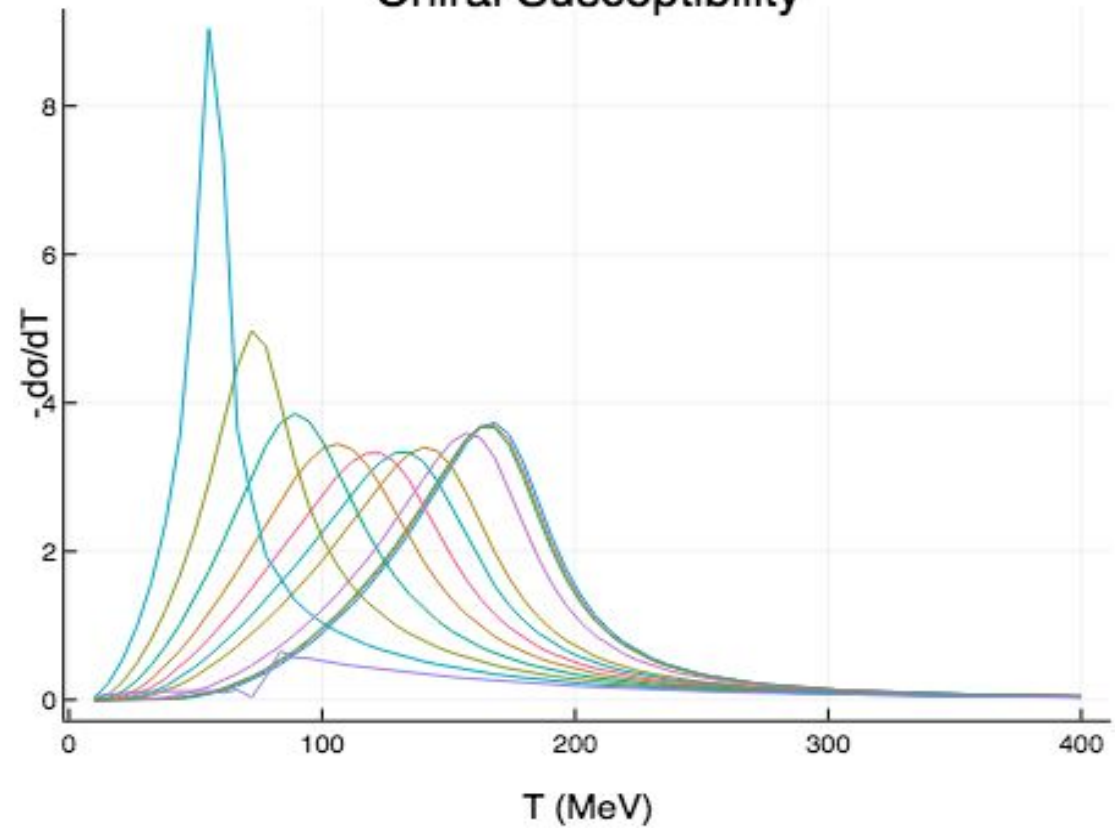
$$\begin{aligned}\bar{\omega} &= 2G_V \int \frac{d^3\vec{p}}{(2\pi)^3} \\ &\times \sum_{\kappa=\pm} \kappa \left\{ 1 - n_F\left(\frac{\bar{E}_b^{\kappa} + \delta\tilde{\mu}_b}{T}\right) - n_F\left(\frac{\bar{E}_b^{\kappa} - \delta\tilde{\mu}_b}{T}\right) \right. \\ &\left. + \frac{2\bar{E}_r^{\kappa}}{\epsilon_r^{\kappa}} \left[1 - n_F\left(\frac{\epsilon_r^{\kappa} + \delta\tilde{\mu}_r}{T}\right) - n_F\left(\frac{\epsilon_r^{\kappa} - \delta\tilde{\mu}_r}{T}\right) \right] \right\},\end{aligned}$$

PRELIMINARY RESULTS - II

quark mass gap

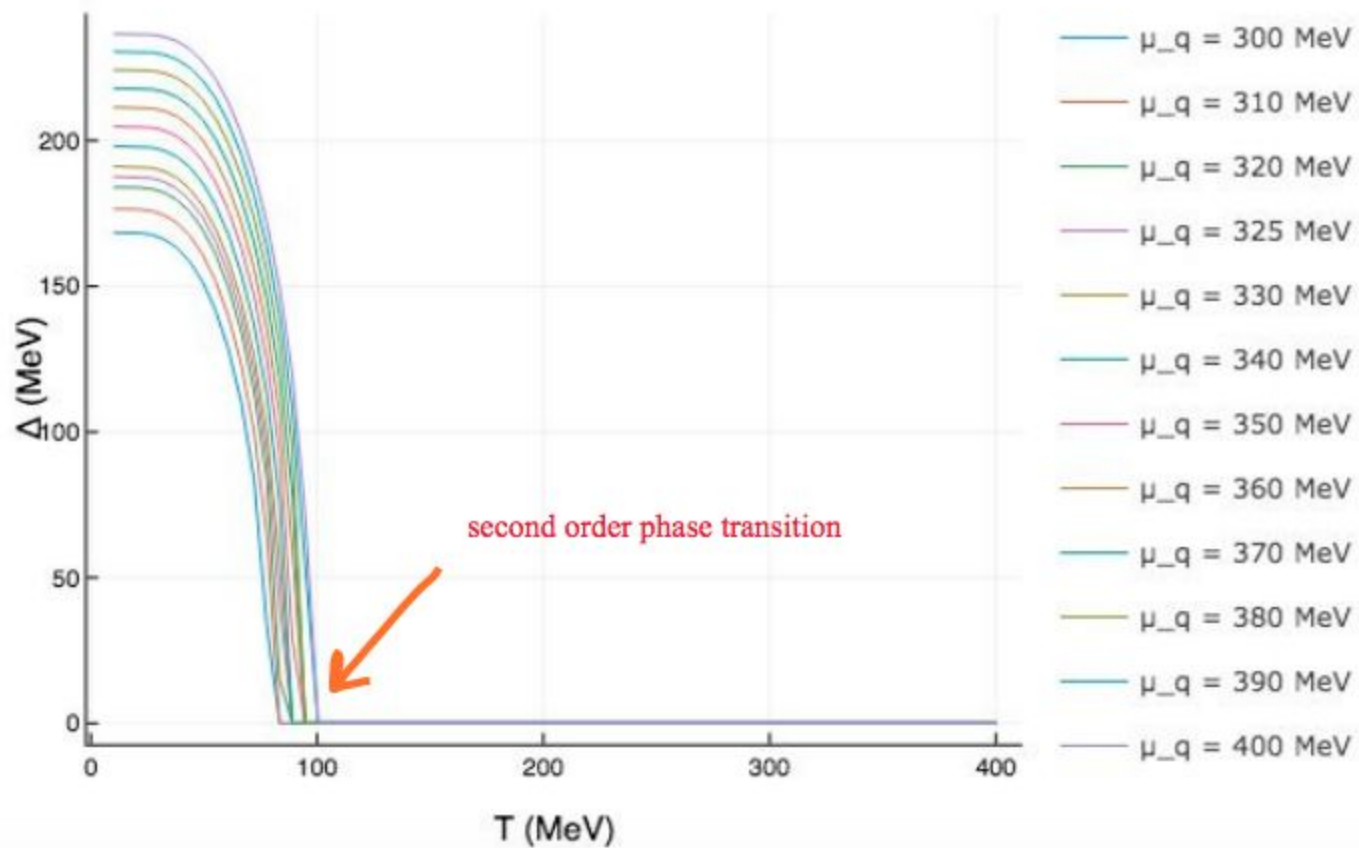


Chiral Susceptibility

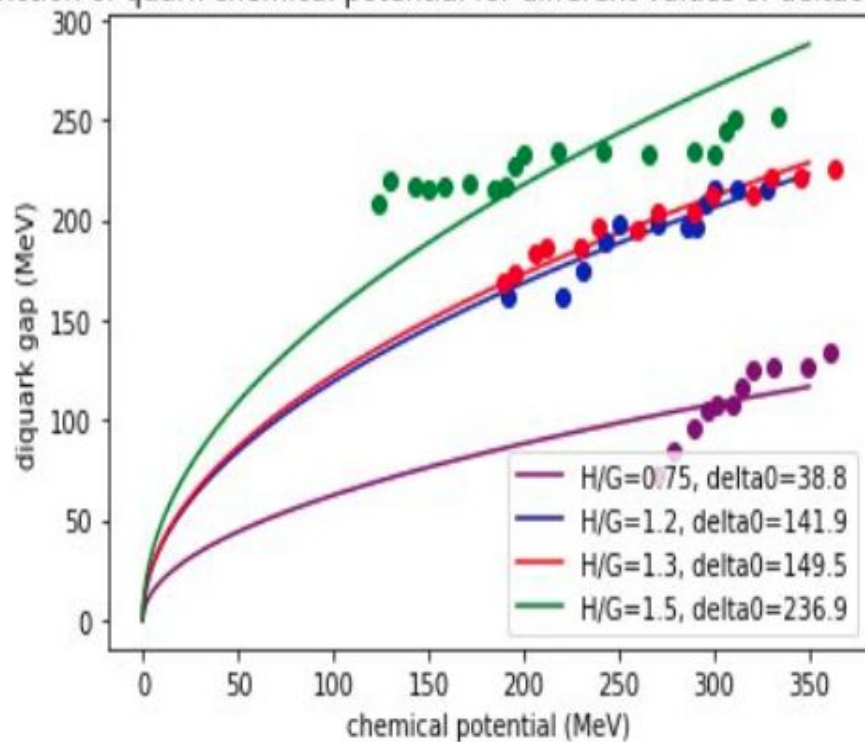


- $\mu_q = 10$ MeV
- $\mu_q = 40$ MeV
- $\mu_q = 50$ MeV
- $\mu_q = 100$ MeV
- $\mu_q = 160$ MeV
- $\mu_q = 180$ MeV
- $\mu_q = 200$ MeV
- $\mu_q = 220$ MeV
- $\mu_q = 240$ MeV
- $\mu_q = 260$ MeV
- $\mu_q = 280$ MeV
- $\mu_q = 300$ MeV

diquark pairing gap



Diquark gap as a function of quark chemical potential for different values of Δ_0 (fit to local NJL model data)



DIQUARK GAP ANSATZ



$$\Delta(\mu) = \sqrt{\Delta_0 \mu / 3}$$

COMPRESSIBILITY



SPEED OF SOUND



STIFFNESS OF EoS



$$P_q(\mu) = \frac{3}{4\pi^2} a_4 \left(\frac{\mu}{3}\right)^4 + \frac{3}{\pi^2} \Delta^2 \left(\frac{\mu}{3}\right)^2 - B_{\text{eff}}$$



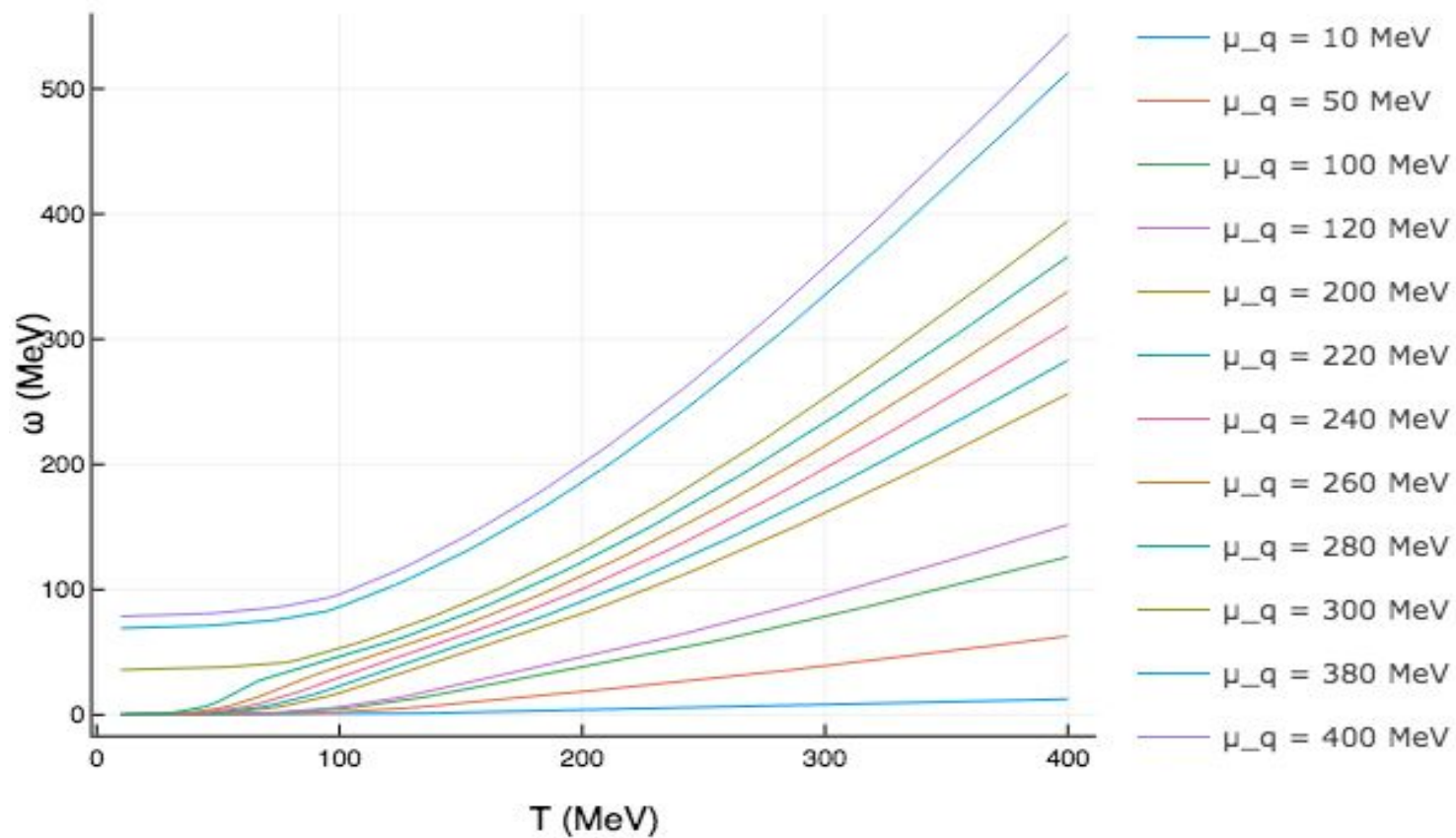
ANSATZ

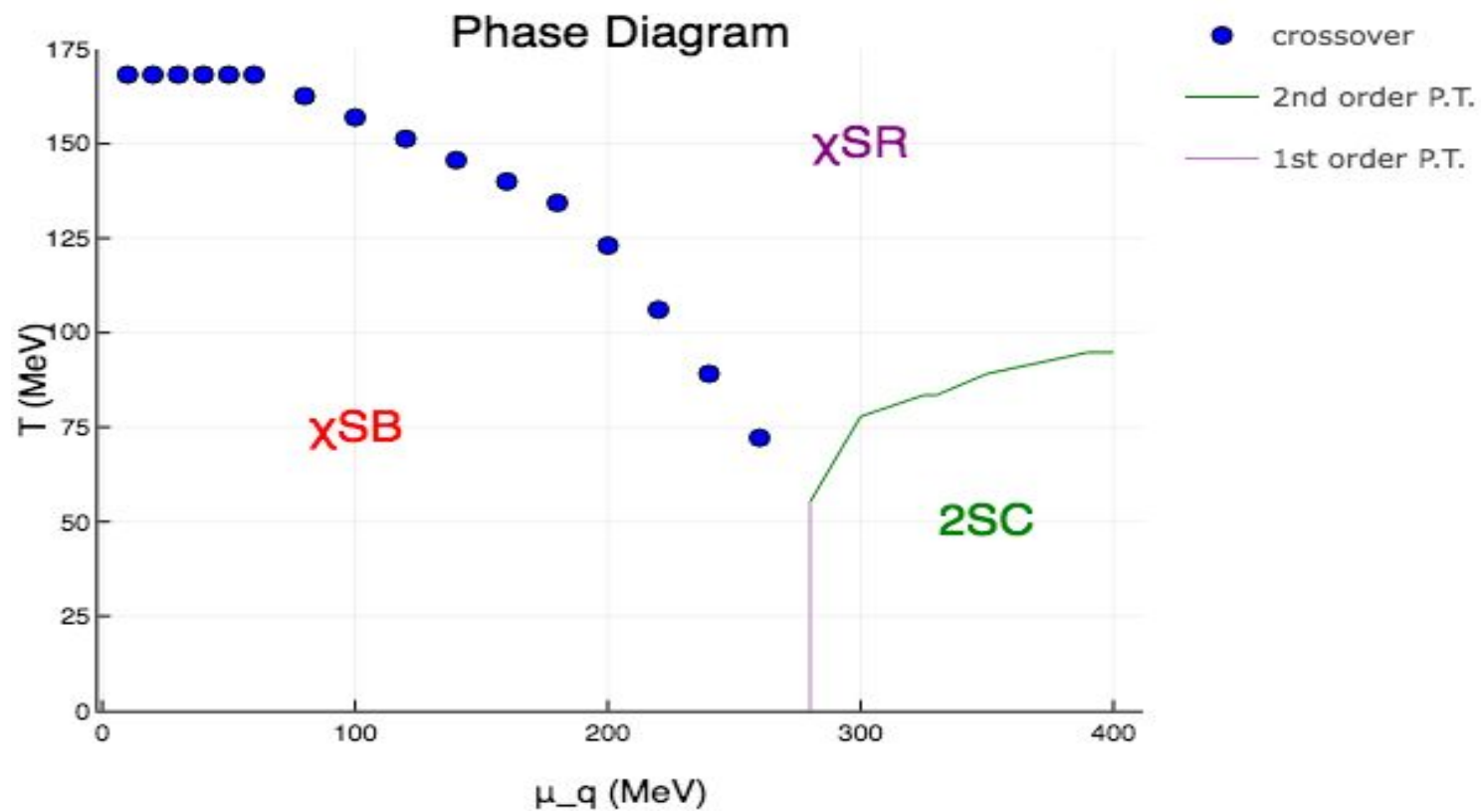


$$\Delta(\mu) = \sqrt{\Delta_0 \mu / 3}$$



$$P_q(\mu) = \frac{3}{4\pi^2} a_4 \left(\frac{\mu}{3}\right)^4 + \frac{3}{\pi^2} \Delta_0 \left(\frac{\mu}{3}\right)^3 - B_{\text{eff}}$$



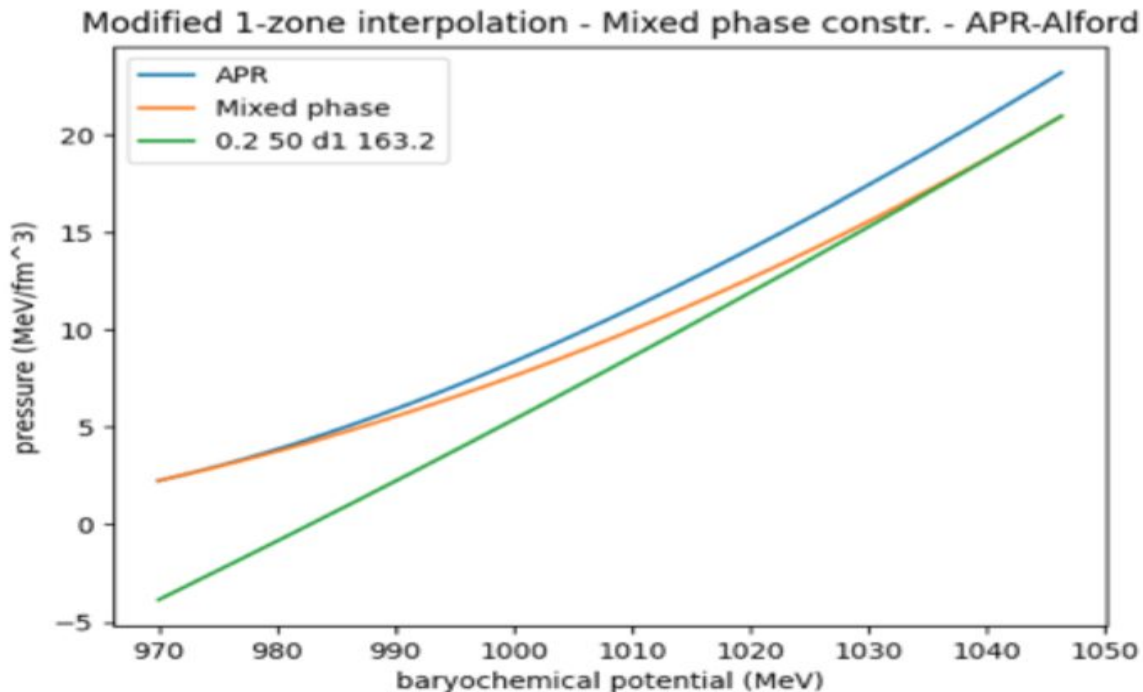


Mixed phase ansatz - $A \cdot \mu^2 + B \cdot \mu + C$

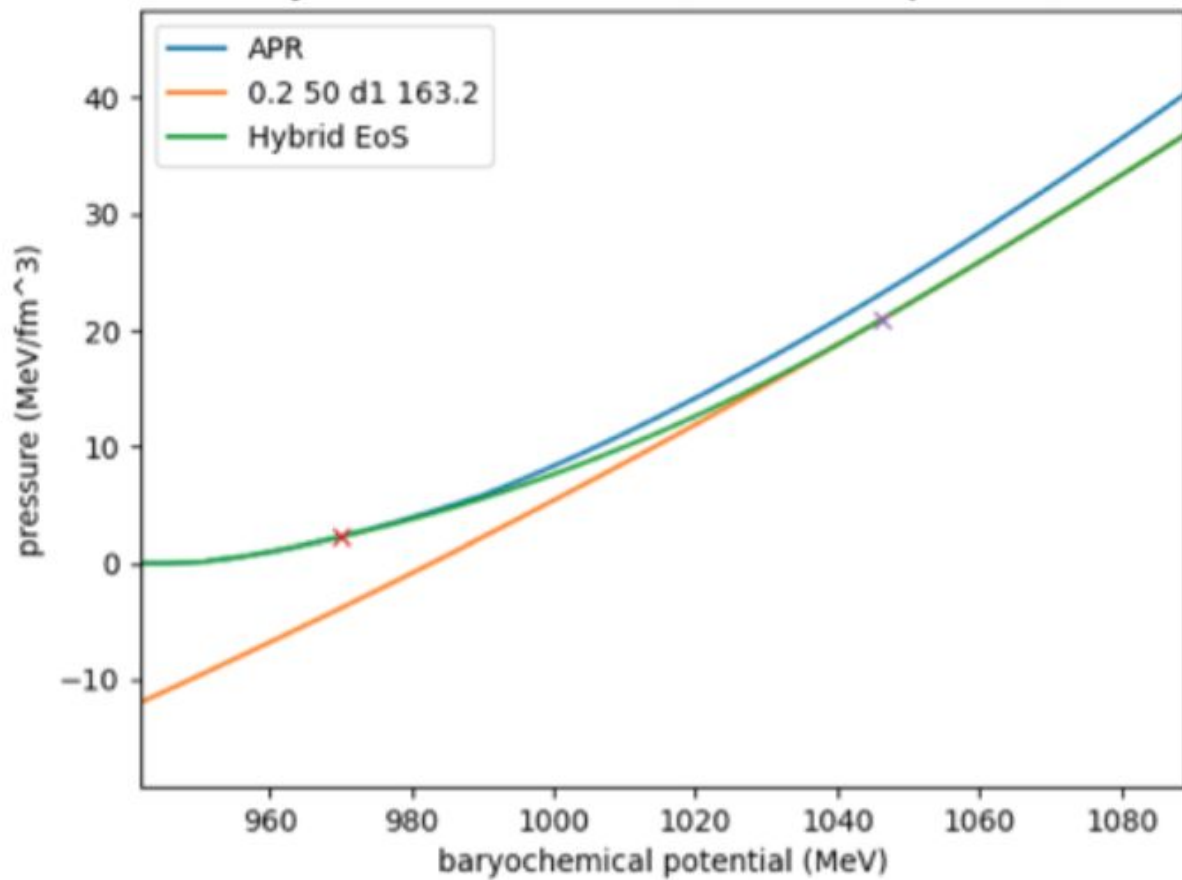
μ_H \rightarrow **fixed at saturation density** = 970 MeV

Unknowns - μ_Q , A, B, C (solved via four continuity equations)

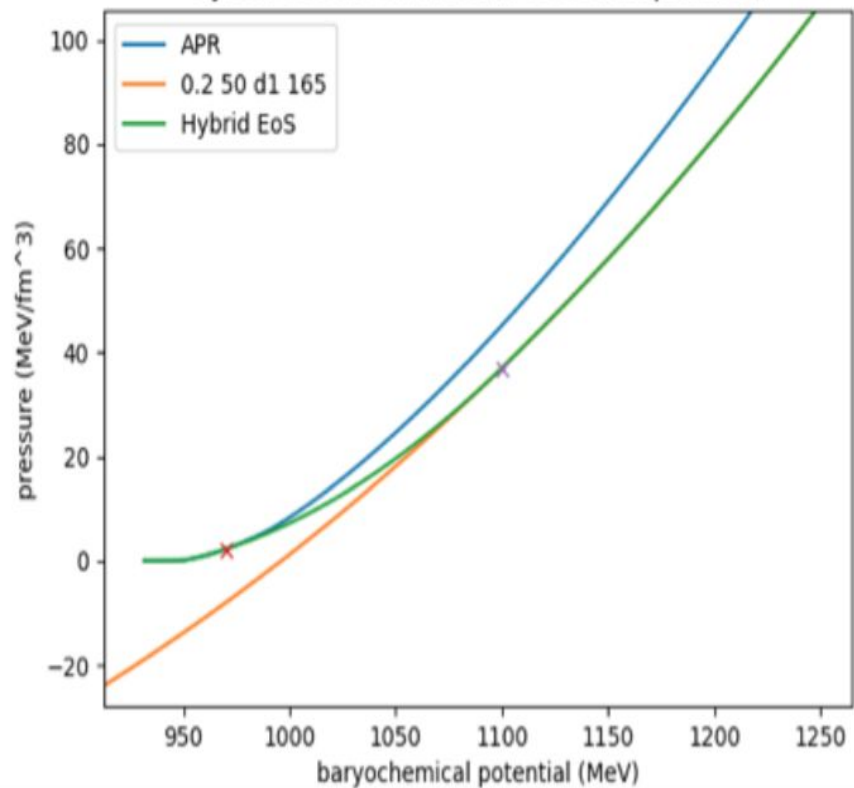
1) Quark EoS (Alford) - 0.2 50 d1 163.2 - (no crossing in the acceptable range of μ)



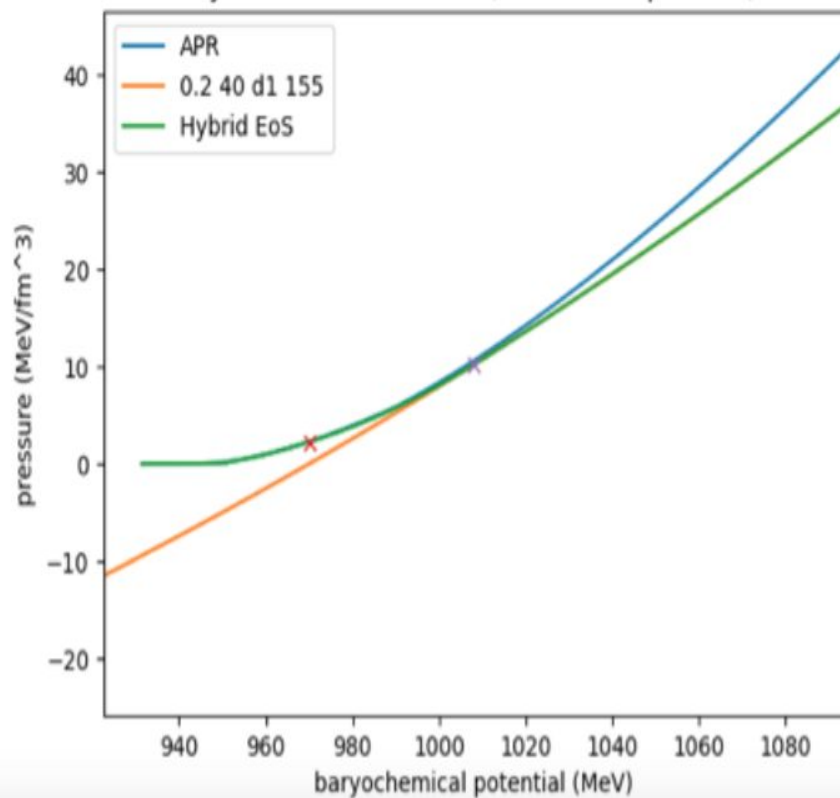
Hybrid EoS : APR-Alford (1-zone interpolation)



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Hybrid EoS : APR-Alford (1-zone interpolation)



Mass-Radius (observational constraint)

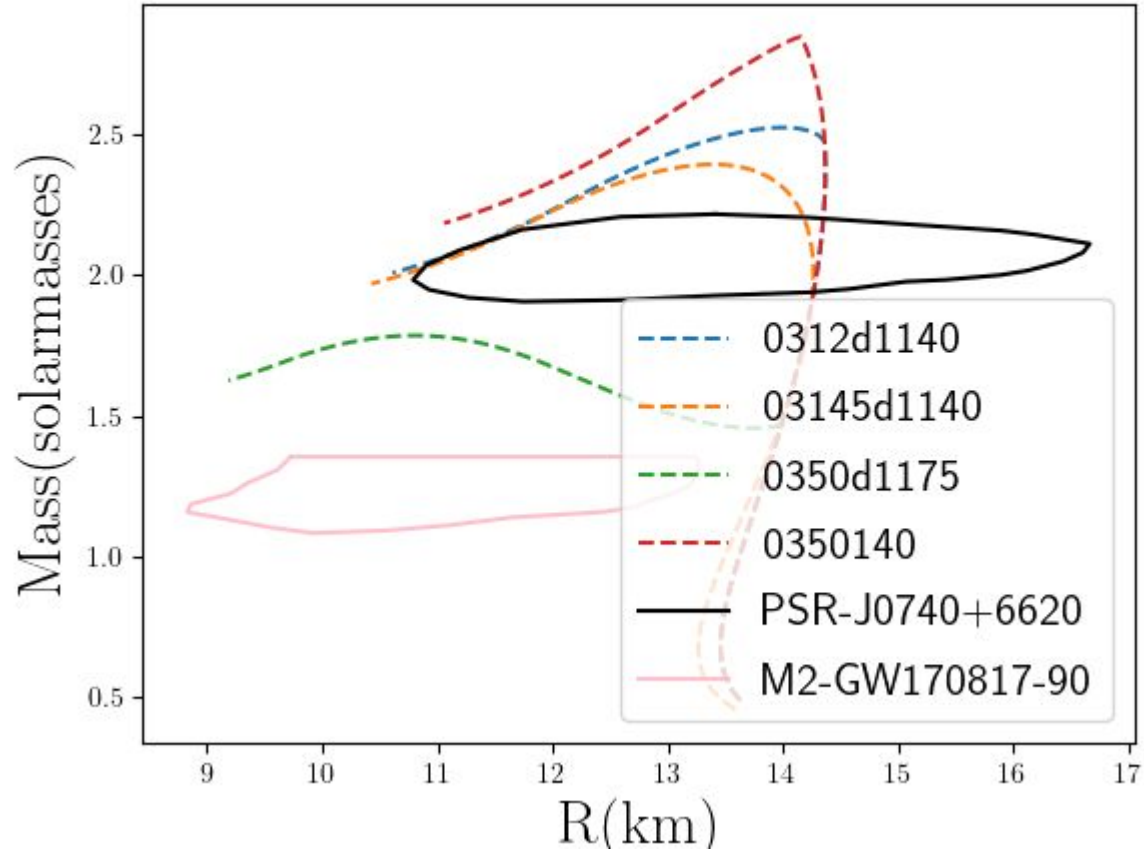


Figure courtesy : S. Liebing

Tidal Deformability (observational constraint)

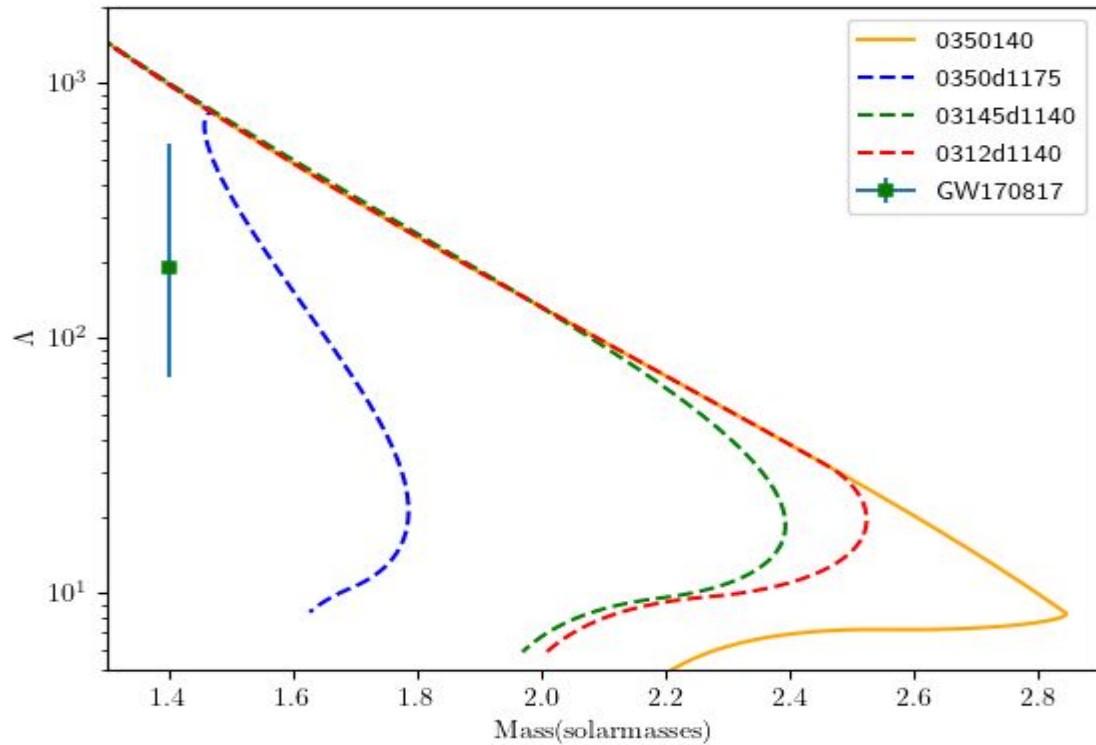


Figure courtesy : S. Liebing

OUTLOOK

- D. Blaschke, M. Buballa, A. Dubinin, G. Ropke, D. Zablocki, Annals of Physics, 348 (2014)
- D. Blaschke, D. Zablocki, Phys. Part. Nucl. 39.1010 (2008)

NJL → PNJL → CONFINEMENT

- M.A. Kaltenborn, N.U.F. Bastian, D. Blaschke, Phys. Rev. D, 96, 056024, (2017)

Relativistic Density Functional (RDF) → String-flip potential model (**2nd order**) → CONFINEMENT

- M. Albright, J. Kapusta, C. Young, Phys. Rev. C 90, 024915 (2014)
- J. Kapusta, T. Welle, Phys. Rev. C 104, L012801 (2021)
- Including better npEoS, eg. density-dependent couplings by S. Typel

Unified EoS

**THANK YOU FOR YOUR
ATTENTION!**

BACKUP SLIDES

3DFF NJL

