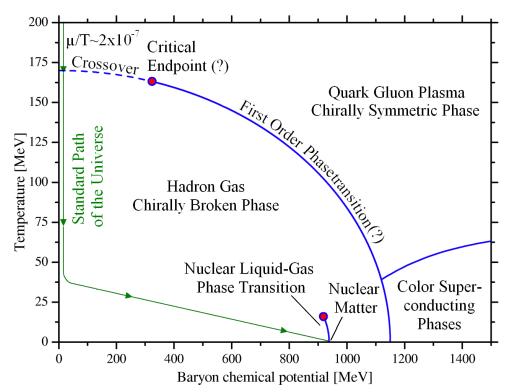


Equation of state of strongly interacting matter for hydrodynamical simulations of heavy-ion collisions

Oleksandr Vitiuk June 21st, 2022 58 Karpacz Winter School of Theoretical Physics, Poland



Introduction

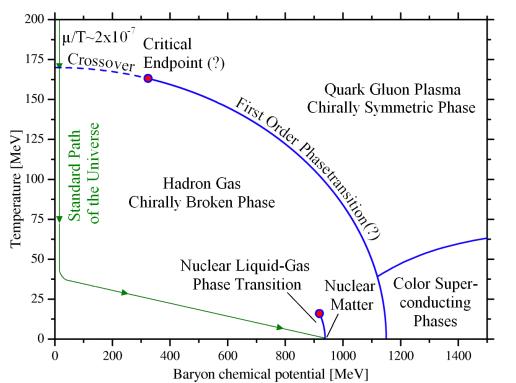


The main goal of heavy-ion collisions experiments is the understanding theory of strong interactions - QCD.

Exploring of the QCD phase diagram:

- Detect signals of deconfinement PT
- Detect signals of (partial) chiral symmetry restoration
- Locate (tri)critical endpoint(s) if such exists

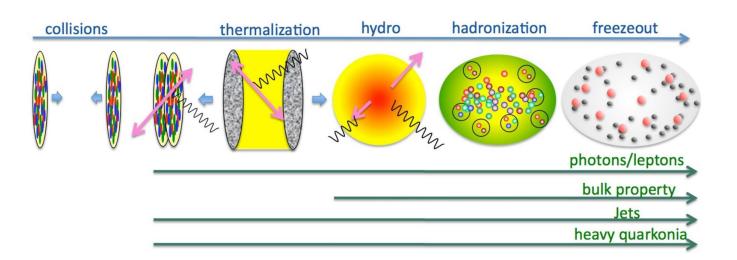
Introduction



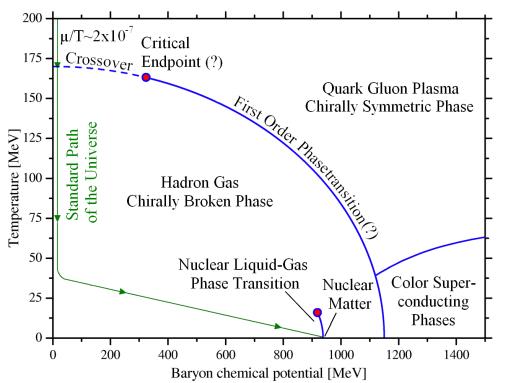
The main goal of heavy-ion collisions experiments is the understanding theory of strong interactions - QCD.

Exploring of the QCD phase diagram:

- Detect signals of deconfinement PT
- Detect signals of (partial) chiral symmetry restoration
- Locate (tri)critical endpoint(s) if such exists



Introduction

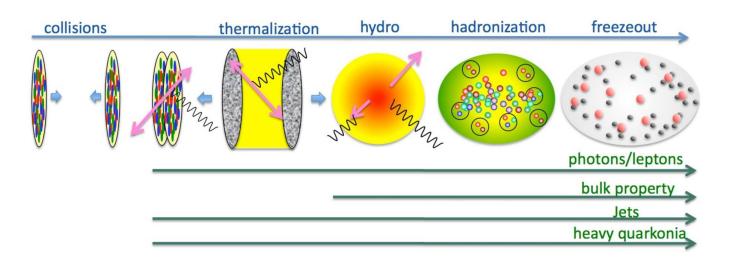


In order to resolve these tasks we need a very good observables and tools to analyse the data!

The main goal of heavy-ion collisions experiments is the understanding theory of strong interactions - QCD.

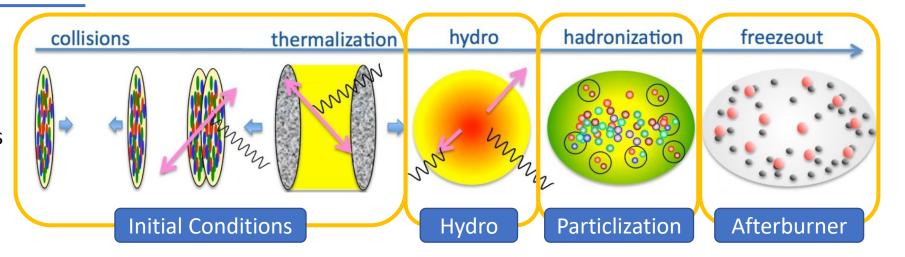
Exploring of the QCD phase diagram:

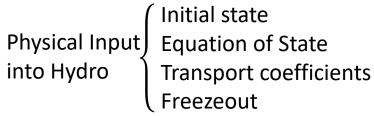
- Detect signals of deconfinement PT
- Detect signals of (partial) chiral symmetry restoration
- Locate (tri)critical endpoint(s) if such exists



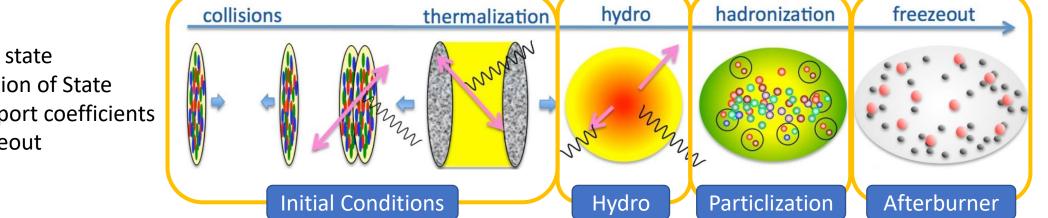
Physical Input into Hydro Freezeout

Physical Input into Hydro Freezeout





EoS



Is needed to build closed system of equations

Contains all the information about the thermodynamic properties of the system that affect its evolution Modeling heavy-ion collisions with different EoS, one can perform a-posteriori analyses of the QCD matter properties in the unsolvable regions of the phase diagram

Physical Input into Hydro Freezeout

EoS

Is needed to build closed system of equations

collisions

Contains all the information about the thermodynamic properties of the system that affect its evolution Modeling heavy-ion collisions with different EoS, one can perform a-posteriori analyses of the QCD matter properties in the unsolvable regions of the phase diagram

Initial Conditions

Construction of the family of equation of states with and without 1st order PT Heavy-ion collision simulations; Computation of observables and comparison with existing data; Model calibration & constraints on EoS

hydro

Hydro

thermalization

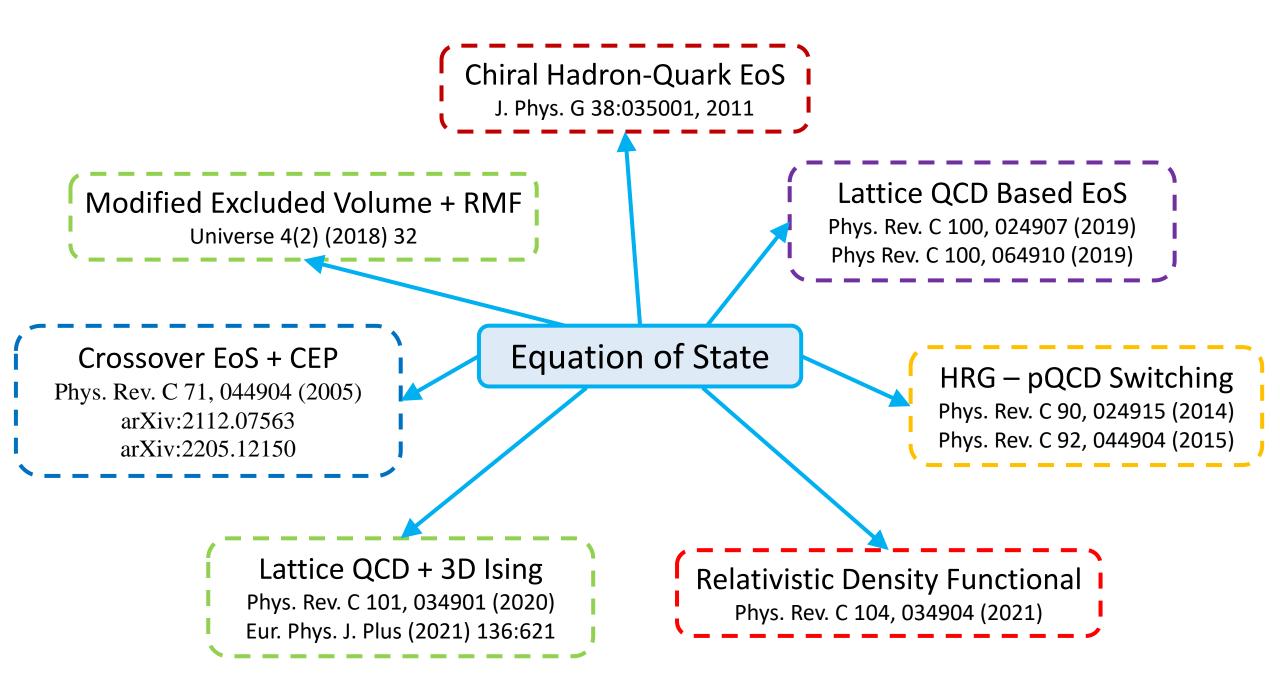
hadronization

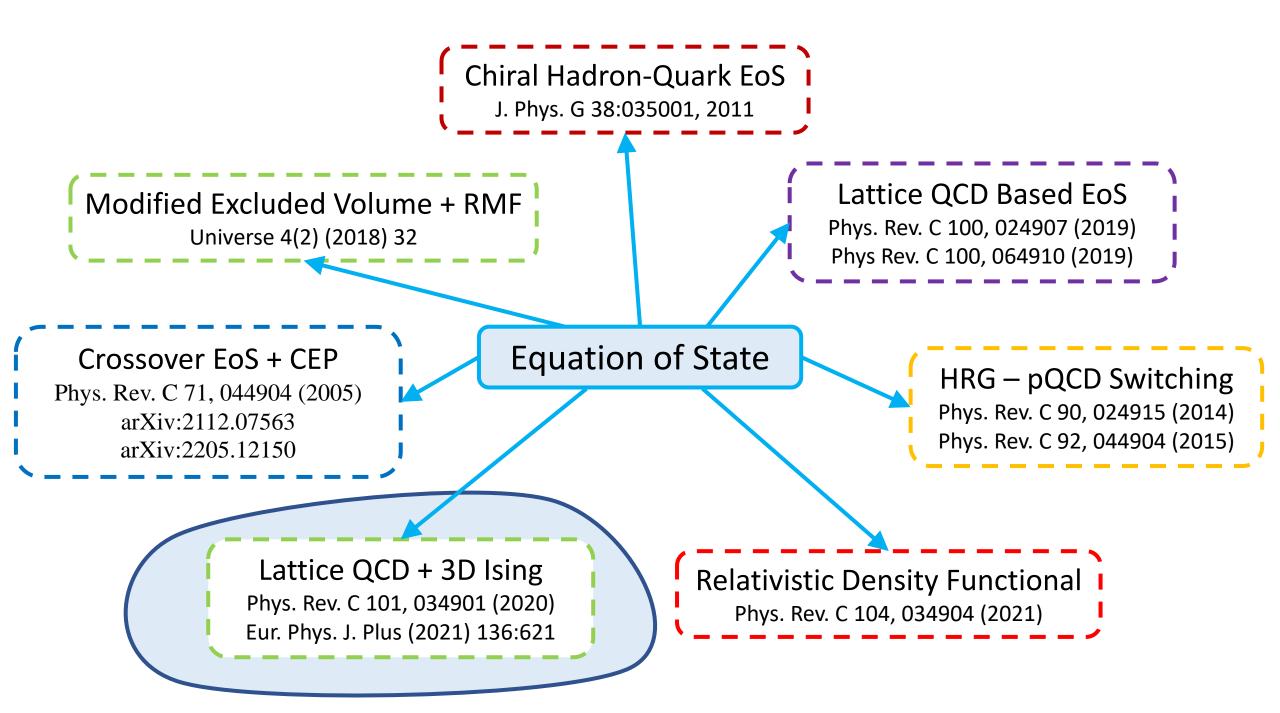
Particlization

freezeout

Afterburner

Prediction for basic observables for future experiments



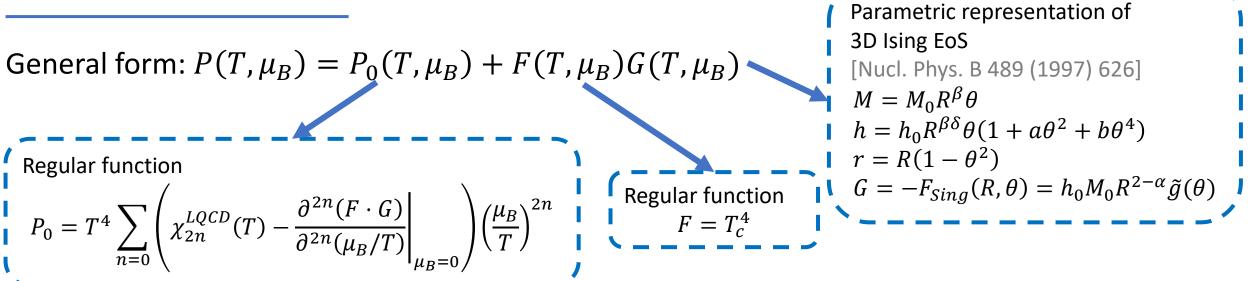


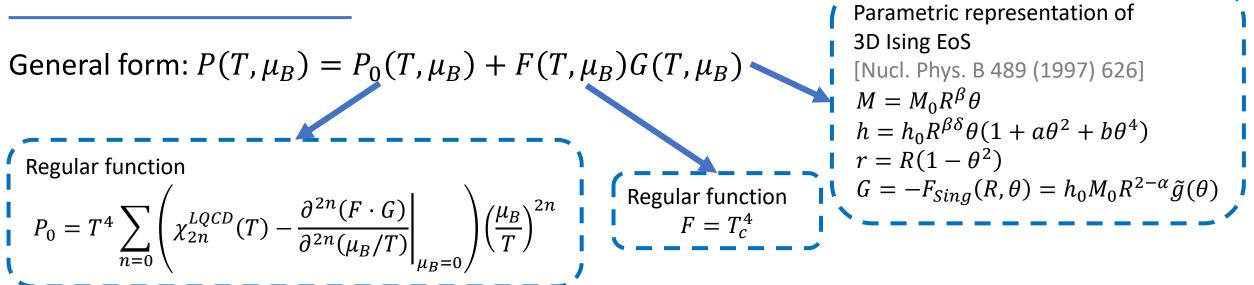
General form: $P(T, \mu_B) = P_0(T, \mu_B) + F(T, \mu_B)G(T, \mu_B)$

General form: $P(T, \mu_B) = P_0(T, \mu_B) + F(T, \mu_B)G(T, \mu_B)$

Parametric representation of 3D Ising EoS [Nucl. Phys. B 489 (1997) 626] $M = M_0 R^{\beta} \theta$ $h = h_0 R^{\beta\delta} \theta (1 + a\theta^2 + b\theta^4)$ $r = R(1 - \theta^2)$ $G = -F_{Sing}(R, \theta) = h_0 M_0 R^{2-\alpha} \tilde{g}(\theta)$

General form: $P(T, \mu_B) = P_0(T, \mu_B) + F(T, \mu_B)G(T, \mu_B)$ $M = M_0 R^{\beta} \theta$ $h = h_0 R^{\beta\delta} \theta(1 + a\theta^2 + b\theta^4)$ $r = R(1 - \theta^2)$ $G = -F_{sing}(R, \theta) = h_0 M_0 R^{2-\alpha} \tilde{g}(\theta)$





One more mapping is needed! $(R, \theta) \rightarrow (r, h) \leftarrow (T, \mu_B)$

 $T = T_c + a_{11}r + a_{12}h$ $\mu_B = \mu_c + a_{21}r + a_{22}h$

Linear

General form:
$$P(T, \mu_B) = P_0(T, \mu_B) + F(T, \mu_B)G(T, \mu_B)$$

Regular function
 $P_0 = T^4 \sum_{n=0}^{\infty} \left(\chi_{2n}^{LQCD}(T) - \frac{\partial^{2n}(F \cdot G)}{\partial^{2n}(\mu_B/T)} \Big|_{\mu_B = 0} \right) \left(\frac{\mu_B}{T} \right)^{2n}$
Regular function
 $P_0 = T^4 \sum_{n=0}^{\infty} \left(\chi_{2n}^{LQCD}(T) - \frac{\partial^{2n}(F \cdot G)}{\partial^{2n}(\mu_B/T)} \Big|_{\mu_B = 0} \right) \left(\frac{\mu_B}{T} \right)^{2n}$
Regular function

One more mapping is needed!

$$(R, \theta) \rightarrow (r, h) \leftarrow (T, \mu_B)$$

$$T = \frac{1}{2} \left(2T_c(h+1) - \mu_c(1-r) \right)$$

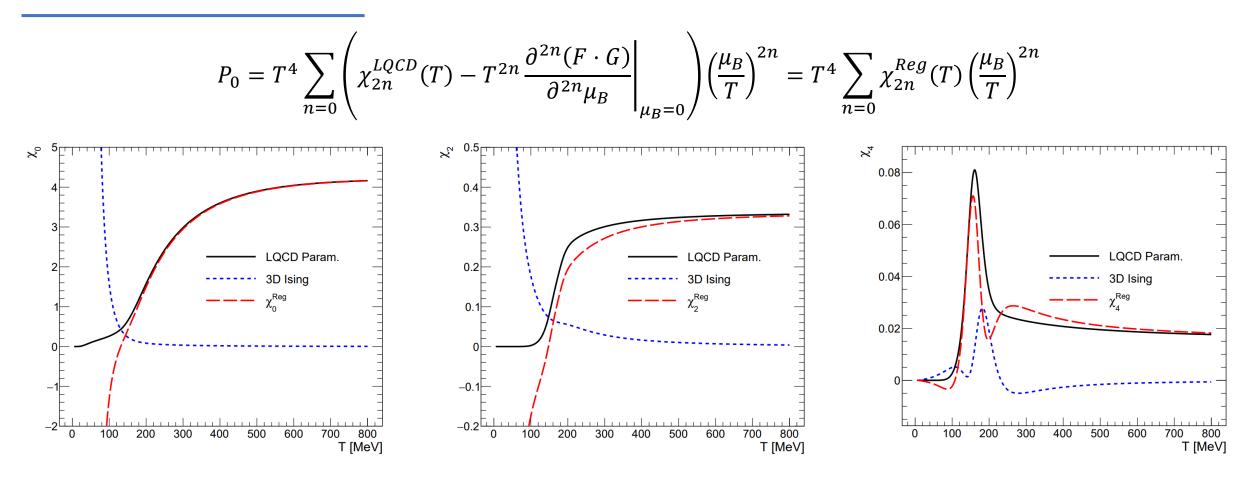
$$\mu_B = a \sqrt{2T_c \mu_c(h+1)(1-r)}$$
Parabolic

$$\mu_B = a \sqrt{2T_c \mu_c(h+1)(1-r)}$$

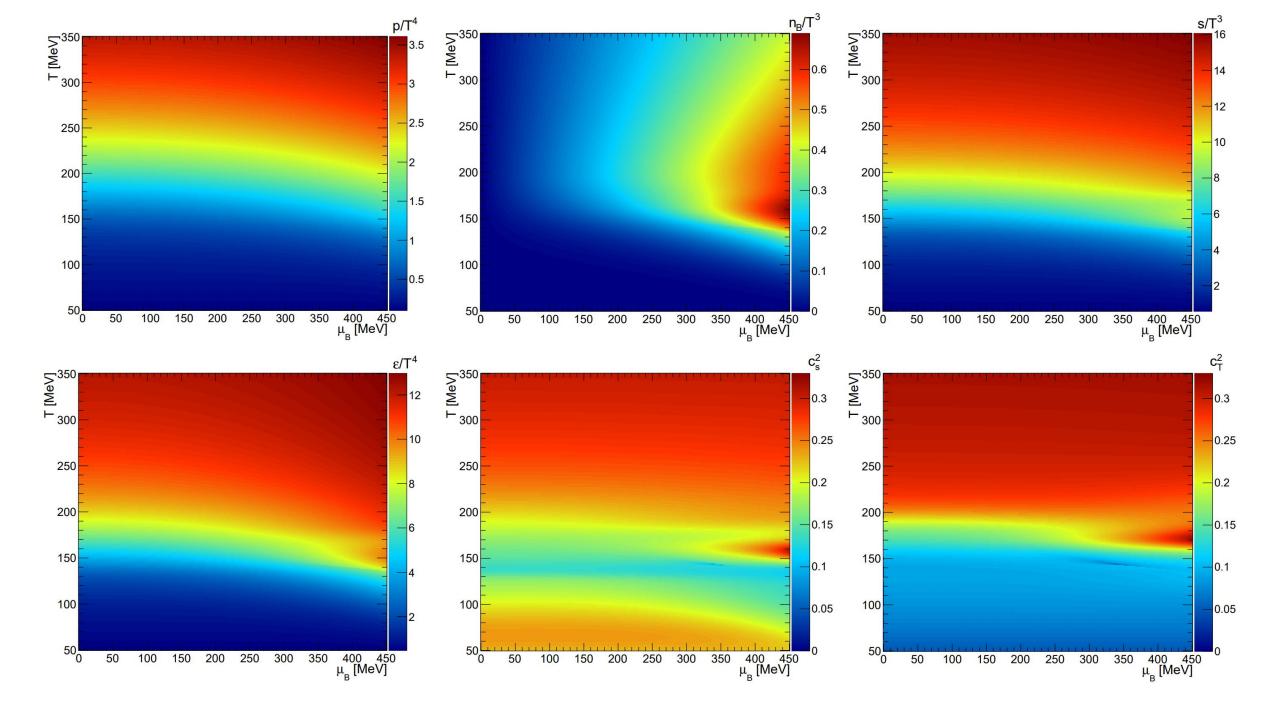
$$T = \frac{T_{pc}}{c} (h+1) \sqrt{c^2 - \mu_{pc}^2(1-r)^2}$$

$$\mu_B = \frac{\mu_{pc}}{c} (1-r) \sqrt{c^2 + T_{pc}^2(1+h)^2}$$
Elliptic

EoS Construction Parametric representation of 3D Ising EoS General form: $P(T, \mu_B) = P_0(T, \mu_B) + F(T, \mu_B)G(T, \mu_B)$ [Nucl. Phys. B 489 (1997) 626] $M = M_0 R^\beta \theta$ $h = h_0 R^{\beta \delta} \theta (1 + a\theta^2 + b\theta^4)$ $r = R(1 - \theta^2)$ **Regular function** $G = -F_{Sing}(R,\theta) = h_0 M_0 R^{2-\alpha} \tilde{g}(\theta)$ **Regular function** $P_0 = T^4 \sum_{n=1}^{\infty} \left(\chi_{2n}^{LQCD}(T) - \frac{\partial^{2n}(F \cdot G)}{\partial^{2n}(\mu_R/T)} \right) \left(\frac{\mu_B}{T} \right)^{2n} \qquad \text{Regular functions} F = T_c^4$ [Me/ One more mapping is needed! $(R,\theta) \rightarrow (r,h) \leftarrow (T,\mu_R)$ 120 $T = \frac{1}{2} \left(2T_c(h+1) - \mu_c(1-r) \right)$ Parabolic 100 Linear $\mu_B = a \sqrt{2T_c \mu_c (h+1)(1-r)}$ Parabolic $T = T_c + a_{11}r + a_{12}h$ 80 Elliptic $\mu_B = \mu_c + a_{21}r + a_{22}h$ $T = \frac{T_{pc}}{c}(h+1)\sqrt{c^2 - \mu_{pc}^2(1-r)^2}$ Linear Elliptic $\mu_B = \frac{\mu_{pc}}{c} (1-r) \sqrt{c^2 + T_{pc}^2 (1+h)^2}$ 20 1000 μ_B [MeV] 600 200 400 800



$$P \Rightarrow fP + (1 - f)P^{HRG}$$
, where $f(T, \mu_B) = \frac{1}{2} \left[1 + \tanh\left(\frac{T - T_0(\mu_B)}{\Delta T}\right) \right]$



Conclusions & Outlook

Conclusions:

- We examined effective EoS construction with the incorporated 3D Ising model singularity
- Such class of EoS constructions allows for the full control on the shape of the phase transition line and the critical endpoint location
- However, it is hard to construct causal and stable EoS
- Other approach is needed

Outlook:

- Try another constructions & provide own one
- Start hydrodynamic simulations to determine observables which are sensitive to different EoS features
- Use Bayesian analyses to constrain parameter space