Systematic 3D MHD simulations of core-collapse supernovae

Ko Nakamura (Fukuoka Univ.)

T. Takiwaki (NAOJ), J. Matsumoto (Keio Univ.), S. Horiuchi (Virginia Tech.), K. Kotake (Fukuoka Univ.)

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Summary

From a massive MS star to a CCSN

✓ The standard scenario toward explosion

A massive star forms iron core.

Gravitational collapse

Fe

~0.1 s BB

- \rightarrow The core gravitationally collapses.
- \rightarrow Shock stalls and revives via neutrino heating.
- \rightarrow Finally the shock breaks out the stellar surface.

Note that the time scale of stellar evolution depends on its mass. Shown is the case of a ~ 10 solar-mass star.

Neutrino trapping

Fe core

v sphere

t = 0



 $\sim 50 \text{ ms PB}$

min. ~ a day

"Standard" CCSNe

We focus on "standard" CCSNe coming from non-extreme initial conditions (mass, rotation, B-field), resulting in normal explosion (not GRB-HNe, not SLSNe).

What is the "standard" explosion? $E_{exp} = 10^{51} \text{ erg}?$ $M_{Ni} = 0.07 \text{ solar mass}?$

Stripped envelope via stellar winds/ binary interaction Ic Ic-BL II II-87A Shivvers+17



✓ O'Connor & Ott (2011)

1D simulations, M = 10-120 Msun, Z = 0-Zsun.

$$Q_{\nu_i}^{\text{heat}}(r) = f_{\text{heat}} \frac{L_{\nu_i}(r)}{4\pi r^2} \sigma_{\text{heat},\nu_i} \frac{\rho}{m_u} X_i \left\langle \frac{1}{F_{\nu_i}} \right\rangle e^{-2\tau_{\nu_i}}$$

Compactness parameter

$$\xi_M = \frac{M/M_{\odot}}{R(M_{\text{bary}} = M)/1000 \,\text{km}}$$

Explosion/non-explosion is non-monotonic and divided by ξ 2.5 ~ 0.45.



✓ Ugliano et al. (2012)

1D simulations, M = 10-40 Msun, Z = Zsun.

$$R_{\rm c}(t) = R_{\rm c,f} + (R_{\rm c,i} - R_{\rm c,f})/(1+t)^n \to L_{\nu}(R_{\rm c}(t))$$

Explosion/non-explosion is non-monotonic, but not well correlated with the compactness.

Explosion properties are also non-monotonic.



✓ Ertl et al. (2016)

Two-parameter criterion for the explodability of 1D models.

$$M_4 \equiv m(s=4)/M_{\odot} \sim M_{\rm Si}$$
$$\mu_4 \equiv \frac{\Delta m/M_{\odot}}{\Delta r/1000 \text{ km}} \bigg|_{s=4} \sim dM/dt$$



✓ Pejcha & Thompson (2015); Pejcha & Prieto (2015)

Observations suggest a linear relation between E_{exp} and M_{Ni} .



Note: It's NOT easy to estimate Eexp!

✓ Popov (1993)

Time scale and luminosity at the plateau phase (Type IIP)

$$t_{\rm SN} = 100 \,\mathrm{d} \, \left(\frac{\kappa}{0.34 \mathrm{cm}^2/\mathrm{g}}\right)^{1/6} \left(\frac{M_{\rm ej}}{10M_{\odot}}\right)^{1/2} \left(\frac{E}{10^{51}\mathrm{erg}}\right)^{-1/6} \left(\frac{R_0}{500R_{\odot}}\right)^{1/6} \left(\frac{T}{5000\mathrm{K}}\right)^{-2/3}$$
$$L = 1.5 \times 10^{42} \,\mathrm{erg/s} \left(\frac{\kappa}{0.34 \mathrm{cm}^2/\mathrm{g}}\right)^{-1/3} \left(\frac{M_{\rm ej}}{10M_{\odot}}\right)^{-1/2} \left(\frac{E}{10^{51}\mathrm{erg}}\right)^{5/6} \left(\frac{R_0}{500R_{\odot}}\right)^{2/3} \left(\frac{T}{5000\mathrm{K}}\right)^{4/3}$$

+ spectrum

$$E = rac{1}{2}Mv^2
ightarrow v = \left(rac{2E}{M}
ight)^{1/2}$$

Luminosity at the tail

$$\begin{split} L(t) &= E_{\rm Ni} \left(-\frac{dN_{\rm Ni}}{dt} \right) + E_{\rm Co} \left(-\frac{dN_{\rm Co}}{dt} \right) \\ &= \frac{E_{\rm Ni} N_{\rm Ni,0}}{\tau_{\rm Ni}} e^{-t/\tau_{\rm Ni}} + \frac{E_{\rm Co} N_{\rm Ni,0}}{\tau_{\rm Co} - \tau_{\rm Ni}} e^{-t/\tau_{\rm Co}} - \frac{E_{\rm Co} N_{\rm Ni,0}}{\tau_{\rm Co} - \tau_{\rm Ni}} \frac{\tau_{\rm Co}}{\tau_{\rm Ni}} e^{-t/\tau_{\rm Ni}} \\ &= \frac{N_{\rm Ni,0}}{\tau_{\rm Co} - \tau_{\rm Ni}} \left(\frac{(\tau_{\rm Co} - \tau_{\rm Ni}) E_{\rm Ni} - \tau_{\rm Co} E_{\rm Co}}{\tau_{\rm Ni}} e^{-t/\tau_{\rm Ni}} + E_{\rm Co} e^{-t/\tau_{\rm Co}} \right) \\ &= 10^{42} \rm erg/s \, \left(6.45 \, e^{-t/\tau_{\rm Ni}} + 1.45 \, e^{-t/\tau_{\rm Co}} \right) \left(\frac{M_{\rm Ni}}{0.1 \, M_{\odot}} \right) \end{split}$$





✓ KN et al. (2015)

2D self-consistent simulations, M = 10.8-75 Msun, Z = 0-Zsun.

Linear relations between some explosion properties and $\,\xi\,.\,$

Simulations ended at $t_{pb} < 1s$.



✓ KN et al. (2015)

2D self-consistent simulations, M = 10.8-75 Msun, Z = 0-Zsun. Linear relations between some

explosion properties and ξ .

Simulations ended at $t_{pb} < 1s$.

✓ KN et al. (2019)

2D long-term simulations.

One model (s17) obtains $E_{exp} > 10^{51}$ erg.

← Long-lasting accretion supplies energy source & targets of neutrinos.

1.2

1.0

0.8

0.6

0.4

0.2

0.0

s13.8 s16.0

s17.0 s19.6

s19.8 -----s20.0 ------

2

diagnostic energy [10⁵¹erg]



✓ Burrows et al. (2020)

3D self-consistent simulations, M = 9-20, 25, 60 Msun, Z = Zsun. Some models (s13,14,15) failed in explosion. Explosion energy is ~ 10^{50} erg.



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✓ Wang et al. (2022)

Compactness does NOT represent the explodability.

$$\frac{P_{\mathrm{ram}}(t) - P_{\mathrm{ram}}(t + \Delta t)}{P_{\mathrm{ram}}(t)}$$

Ram pressure jump determines the fate of a stalled shock.



Summary of the current understandings from the systematic studies

- ✓ Explosion properties are NOT a monotonic function of progenitor mass.
 *t*_revival (explodability), *E*_exp, *M*_NS/BH, *L*_*v*, *M*_Ni, …
- Mass accretion onto the central core plays a dominant role.
 Parameter(s) characterizing the accretion rate (e.x., ξ) and some of the explosion properties are well correlated.
- ✓ Explodability cannot be expressed by the compactness. Two-parameter criterion? Degree of P_ram jump?

Systematic CCSN simulations

	Spatial dim.	Model #	ν heating	ZAMS <i>M</i> [Msun]	Z	sim. time	Summary
O'connor & Ott '11	1D	~100	× factor	10-120	0-solar	~1s	Non-monotonic explosion/BH formation.
Ugliano+'13	1D	~100	$L v \left(R_{\rm NS}, t ight)$	10-40	solar	~10s+	Non-monotonic explosion properties.
KN+'15	2D	~400	Self-consistent	10-75	0-solar	<1s	Explosion properties depend on $ \xi .$
KN+'19	2D	10	Self-consistent	10-20	solar	~10s	Long-term accretion produces $E_{exp} > 10^{51}$ erg.
Burrows+'20	3D	14	Self-consistent	9-20,25,60	solar	<1s	Eexp ~ 0.1x10 ⁵¹ erg
KN+, in prep.	3D	16	Self-consistent	9-24	solar	0.5s	Independent 3D study, including B .

3D CCSN simulations with magnetic field

✓ Mueller & Varma (2020)

3D simulation with weak B-field. The weak seed field is amplified and **supports** a development of a neutrino-driven explosion. ✓ Matsumoto et al. (2022)

3D simulations with weak/strong B-fields. The strong B-filed model presents **early shock revival**.



Time after bounce [ms]

3D CCSN simulations with magnetic field

✓ Matsumoto et al. (2022)





3D CCSN simulations with magnetic field

✓ Matsumoto et al. (2022)





Systematic 3D MHD simulations - Numerical scheme

✓ <u>3DnSNe_MHD code</u> (*Matsumoto+'20*) based on 3DnSNe code (*Takiwaki+'16,'18*).

- ✓ Initial 2D simulation:
 - No rotation, $A_{\phi} = \frac{B_0}{2} \frac{r_0^3}{r^3 + r_0^3} r \sin \theta$ with $B_0 = 10^{10} [G]$ and $r_0 = 10^3$ km.
 - $600(r)x128(\theta)$ grids for $0 \leq R \leq 10^4$ km and $0 \leq \theta \leq \pi$.

✓ Subsequent 3D simulation:

- $2D \rightarrow 3D$ at 10ms after bounce.
- Random density perturbation (\leq 1%) is imposed in R > 100 km.
- $600(r)x64(\theta)x128(\phi)$ grids for $0 \leq R \leq 10^4$ km, $0 \leq \theta \leq \pi$, and $0 \leq \phi \leq 2\pi$.

Systematic 3D MHD simulations - Progenitor models

✓ 9-24 Msun progenitors from Sukhbold et al. (2016) ApJ, 821, 38

Density profiles of some models (e.x., **s10**, **s11**, **s12**) show a jump at Si/O interface.

This progenitor series have a peak (island) of <u>compactness</u> around $M_{ZMAS} \sim 23$ solar masses.

$$\xi = M(R)/R$$



Systematic 3D MHD simulations - Overview

✓ Movies show iso-surface of entropy for 10 models.

Note that spatial scale (box size) is changing as the shock waves successfully revive and expand. The model s11 presents outstanding burst at 140 ms after bounce.



120ms after bounce.

Systematic 3D MHD simulations - Overview

✓ Movies show iso-surface of entropy for 10 models.

Note that spatial scale (box size) is changing as the shock waves successfully revive and expand. The model s11 presents outstanding burst at 140 ms after bounce.



150ms after bounce.

Systematic 3D MHD simulations - Shock revival

✓ (Top panel) Mass accretion rate @ r = 500km.

Roughly in order of ZAMS mass (or compactness) in the early phase (< 100 ms).

Some models show sudden drop when the Si/O interface passes through.

✓ (Bottom panel) Angle-averaged shock radius.

In some models the shock jumps when the Si/O interface falls onto the shock and ram pressure from the accreting matter is suppressed. \rightarrow Shock revival time is not in order of ZAMS mass.

Finally all examined models successfully revive their shock wave.



Systematic 3D MHD simulations - Explosion energy

- ✓ Diagnostic explosion energy
 - = Ekin + Eint + Egrv of ejecta (unbound & $v_r > 0$).

Here overburden of stellar envelope is not taken into account.

Most models show Eexp < 0.2 x 10⁵¹ erg @500ms except (s23 &) s24 model.



Multi-messenger signals from CCSN

✓ KN+ 2016, MNRAS, 461, 3296



Systematic 3D MHD simulations - Neutrino LC

✓ (Top panel) Mass accretion rate @ r = 500km.

Roughly in order of ZAMS mass (or compactness) in the early phase (< 100 ms).

Some models show sudden drop when the Si/O interface passes through.

✓ (Bottom panel) Anti-e neutrino luminosity.

Neutrino luminosity \sim accretion rate \sim compactness.

→ Neutrino detection event number could restrict the progenitor compactness.

The sudden drop in mass accretion rate causes luminosity jump.

→ It tells us the density structure of the progenitor star.



Systematic 3D MHD simulations - Neutrino detection event

- ✓ Assuming D = 10 kpc.
- ✓ SK: fiducial volume = 22.5 kton, threshold energy = 7 MeV.
- ✓ Only IBD is considered. No oscillation.



Systematic 3D MHD simulations - Neutrino (3)

✓ Total neutrino energy emitted between 10-500 ms after bounce.

It well correlates to the gravitational potential energy released via the core collapse, which can be represented by the Fe core mass (shown below), compactness parameter, or PNS mass.



Systematic 3D MHD simulations - GW (1)

✓ GW waveforms for 9 of 12 models.



Systematic 3D MHD simulations - GW (2)

✓ GW spectrogram of s9 (left) and s20 (right).

Common features: silent phase (< 100 ms) \rightarrow nonlinear phase \rightarrow explosion phase. Amplitude is correlated to the mass (~ compactness) of the progenitor stars.



Systematic 3D MHD simulations - GW (3)

✓ GW spectra of s9 (small mass & small ξ) and s20 (high mass & high ξ), compared with sensitivity curves of some current and future GW detectors (D = 10 kpc).

Common features:

High-frequency components coming from central region are dominant.

s9: Only $f > \sim 200$ Hz is detectable.

s20: low-f signal (>~70 Hz) including PNS convection signal could be detectable.



NS properties - mass



1.4 - 2 Msun.

High accretion rate produces massive PNSs.

The time evolution is not (directly) observable. We could guess from a ν luminosity curve. "Observed" data from Table 1 in *Lattimer (2012) Annu. Rev. Nucl. Part. Sci., 62, 485.* "Model prediction" is IMF (Salpeter) weighted.

1.2-1.7 Msun well reproduced!Low mass NSs from binary interaction?High mass ones supported by rapid rotation?

NS properties - kick velocity



NS properties - spin

Spiral mode of SASI in s24 model.



(Note: our simulations start from non-rotating progenitors.)



CCSN properties - what determines?







Summary

- ✓ Systematic study of 3D CCSN models is still challenging but now it's a feasible idea.
- ✓ We demonstrate 3D MHD simulations for 9-24 solar mass progenitors (Sukhbold+'16).
- ✓ All the examined models show successful shock revival in 300 ms.
- ✓ The MM signals predicted from our 3D CCSN models confirm the previous findings in 1D/2D simulations:
 - The v luminosity and average energy differs between the models by a factor of $<\sim 2$.
 - The total neutrino energy emitted during our simulations well correlates to the Fe core mass of the progenitor stars.
 - GW signals have silent phase (< 100 ms) \rightarrow nonlinear phase \rightarrow explosion phase.
 - High-*f* strong signal and low-*f* weak signal exist.
 - Amplitude is correlated to the mass (compactness) of progenitor stars.
 - \rightarrow They will provide us with fruitful information on the structure of the CCSN core!