

Future studies of photon-photon scattering in ultraperipheral heavy ion collisions

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in collaboration with P. Jucha, M. Klusek-Gawenda
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Various faces of QCD, April 26-28, 2024, Wrocław, Poland

Introduction

- $\gamma\gamma \rightarrow \gamma\gamma$ scattering is fully quantal effect.
- Several mechanisms can be present in principle.
- It was/is a **dream for the laser community**.
Some works in this direction are going on.
- It was proposed by: **D. d'Enterria, G. Silveira** to do it in heavy ion UPC in 2013
(box diagrams only)
- **M. Klusek-Gawenda, P. Lebiedowicz and A. Szczurek** made an update in 2016.
Different result numerically and a **double-hadronic excitation** (virtual vector mesons),
called by us VDM Regge, was discussed.
- **M. Klusek-Gawenda, W. Schafer and A. Szczurek**,
two-gluon exchange mechanism.
- **ATLAS** and **CMS** performed such experiments and got results consistent with our
predictions.
- **M. Klusek-Gawenda and A. Szczurek** considered a background contribution:
 $\gamma\gamma \rightarrow \pi^0(\rightarrow 2\gamma)\pi^0(\rightarrow 2\gamma)$, also in $AA \rightarrow AA\pi^0\pi^0$.
- **M. Klusek-Gawenda, R. McNulty, R. Schicker and A. Szczurek** studied what **ALICE**
and **LHCb** could do.
- There are at least two experimental projects that could provide new insights into
photon-photon scattering: **FoCal** and **ALICE 3**.

Sketch of the formalism, EPA

The nuclear cross section is calculated using equivalent photon approximation (EPA) in the b -space. In this approach, the di-photon cross section can be written as:

$$\frac{d\sigma(PbPb \rightarrow PbPb\gamma\gamma)}{dy_{\gamma 1} dy_{\gamma 2} dp_{t,\gamma}} = \int \frac{d\sigma_{\gamma\gamma \rightarrow \gamma\gamma}(W_{\gamma\gamma})}{dz} N(\omega_1, b_1) N(\omega_2, b_2) S_{abs}^2(b) \\ \times d^2 b d\bar{b}_x d\bar{b}_y \frac{W_{\gamma\gamma}}{2} \frac{dW_{\gamma\gamma} dY_{\gamma\gamma}}{dy_{\gamma 1} dy_{\gamma 2} dp_{t,\gamma}} dz, \quad (1)$$

where $\bar{b}_x = (b_{1x} + b_{2x})/2$ and $\bar{b}_y = (b_{1y} + b_{2y})/2$. The relation between $\vec{b}_1$, $\vec{b}_2$ and impact parameter: $b = |\vec{b}| = \sqrt{|\vec{b}_1|^2 + |\vec{b}_2|^2 - 2|\vec{b}_1||\vec{b}_2|\cos\phi}$. Absorption factor $S_{abs}^2(b)$ is calculated as:

$$S_{abs}^2(b) = \Theta(b - b_{max}) \quad (2)$$

or

$$S_{abs}^2(b) = \exp(-\sigma_{NN} T_{AA}(b)), \quad (3)$$

where σ_{NN} is the nucleon-nucleon interaction cross section.

Sketch of the formalism, EPA

$T_{AA}(b)$ above is related to the so-called nuclear thickness, $T_A(b)$,

$$T_{AA}(|\vec{b}|) = \int d^2\rho T_A(\vec{\rho} - \vec{b}) T_A(\vec{\rho}), \quad (4)$$

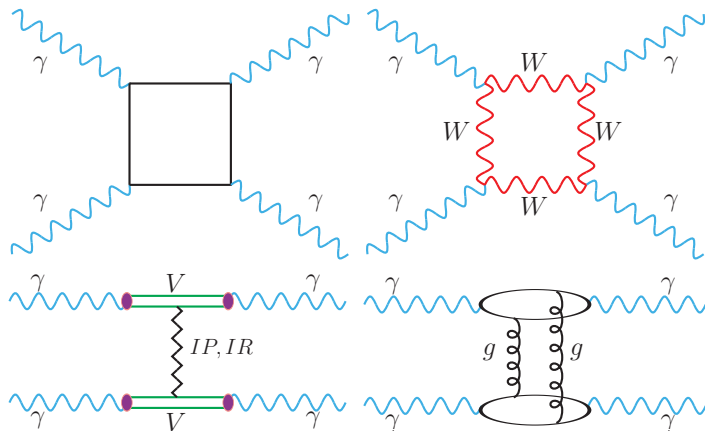
and the nuclear thickness is obtained by integrating the nuclear density

$$T_A(\vec{\rho}) = \int \rho_A(\vec{r}) dz, \quad \vec{r} = (\vec{\rho}, z), \quad (5)$$

where ρ_A is the nuclear charge distribution.

The nuclear photon fluxes $N(\omega_1, b_1)$ and $N(\omega_2, b_2)$ are calculated with realistic charge distribution.

Mechanisms of photon-photon scattering



Elementary cross section

The elementary cross section here is calculated for unpolarised photons. To do this, all **16 photon helicity combinations** of the cross section must be added up. In fact, due to symmetries, it is sufficient to count **only five combinations** and then add them up with the corresponding weights:

$$\begin{aligned} \Sigma_{\text{helicities}} |\mathcal{A}_{\gamma\gamma\rightarrow\gamma\gamma}|^2 &= 2|\mathcal{A}_{++++}^{\text{fermions}}|^2 + 2|\mathcal{A}_{+---}^{\text{fermions}}|^2 \\ &+ 2|\mathcal{A}_{+-+-}^{\text{fermions}}|^2 + 2|\mathcal{A}_{++--}^{\text{fermions}}|^2 + 8|\mathcal{A}_{+-+-}^{\text{fermions}}|^2 . \end{aligned}$$

(6)

Elementary cross section calculations for the box contribution were carried out using **FormCalc** and **LoopTools libraries** based on **Mathematica** software.

Elementary cross section, VDM

This component was calculated for the first time by Klusek-Gawenda, Lebedowicz, Szczurek 2016 assuming vector dominance model. In this approach, the amplitude is:

$$\begin{aligned} \mathcal{M} = & \sum_{i,j} C_i^2 C_j^2 \left(\mathcal{C}_{\mathbb{P}} \left(\frac{s}{s_0} \right)^{\alpha_{\mathbb{P}}(t)-1} F(t) + \mathcal{C}_{\mathbb{R}} \left(\frac{s}{s_0} \right)^{\alpha_{\mathbb{R}}(t)-1} F(t) \right), \\ & + \sum_{i,j} C_i^2 C_j^2 \left(\mathcal{C}_{\mathbb{P}} \left(\frac{s}{s_0} \right)^{\alpha_{\mathbb{P}}(u)-1} F(u) + \mathcal{C}_{\mathbb{R}} \left(\frac{s}{s_0} \right)^{\alpha_{\mathbb{R}}(u)-1} F(u) \right). \end{aligned} \quad (7)$$

In the simplest version of the model $i, j = \rho^0, \omega, \phi$ (only light vector mesons are included). The couplings C_i, C_j describe the $\gamma \rightarrow V_{i/j}$ transitions that are calculated based on vector meson dilepton width. $\mathcal{C}_{\mathbb{P}}$ and $\mathcal{C}_{\mathbb{R}}$ are extracted from the Regge factorization hypothesis (Szczurek, Nikolaev, Speth).

It was shown that the component is concentrated mainly at small photon transverse momenta which at not too small subsystem energies corresponds to $z \approx \pm 1$. The Regge trajectories are usually written in a linear form:

$$\begin{aligned} \alpha_{\mathbb{P}}(t/u) &= \alpha_{\mathbb{P}}(0) + \alpha'_{\mathbb{P}} t/u, \\ \alpha_{\mathbb{R}}(t/u) &= \alpha_{\mathbb{R}}(0) + \alpha'_{\mathbb{R}} t/u. \end{aligned} \quad (8)$$

These linear forms are valid at not too large $|t|$ or $|u|$. At large $|t|$ or $|u|$ the energy dependent factors are artificially small. Therefore here where we explore it more, we propose to smoothly switch off the t/u dependent terms in (8) at $t \sim -0.5 \text{ GeV}^2$ and $u \sim -0.5 \text{ GeV}^2$.

Elementary cross section, VDM

We also wish to analyze whether the more heavy vector mesons such as J/ψ can give a sizeable contribution.

For example, for the double J/ψ fluctuations (both photons fluctuate into virtual J/ψ mesons) we take the following Ansatz for the helicity conserving amplitude:

$$\mathcal{M}_{VDM}^{J/\psi J/\psi} = g_{J/\psi}^2 C_{\mathbb{P}}^{J/\psi} \left(\frac{s}{s_0} \right)^{\alpha_{\mathbb{P}}^{J/\psi J/\psi}(t)-1} F_{J/\psi J/\psi \mathbb{P}}^H(t) F_{J/\psi J/\psi \mathbb{P}}^H(t) \quad (9)$$

$$+ g_{J/\psi}^2 C_{\mathbb{P}}^{J/\psi} \left(\frac{s}{s_0} \right)^{\alpha_{\mathbb{P}}^{J/\psi J/\psi}(u)-1} F_{J/\psi J/\psi \mathbb{P}}^H(u) F_{J/\psi J/\psi \mathbb{P}}^H(u) . \quad (10)$$

In this case (double J/ψ fluctuations) only pomeron can be exchanged (no subleading regions are possible due to the $c\bar{c}$ structure of J/ψ mesons). In this case, for simplicity, we take the simplified trajectories as

$$\alpha_{\mathbb{P}}^{J/\psi J/\psi}(t) = \alpha_{\mathbb{P}}^{J/\psi J/\psi}(u) = \alpha_{\mathbb{P}}^{J/\psi J/\psi}(0) . \quad (11)$$

Here the t/u dependencies are totally ignored. In numerical calculations we take $\alpha_{\mathbb{P}}^{J/\psi J/\psi}(0) = 1.3 - 1.4$ (typical hard pomeron).

Elementary cross section, VDM

Since the J/ψ mesons are far off-mass-shell and more compact than light vector mesons also the form factors must be modified:

$$F_{J/\psi J/\psi \mathbb{P}}^H(t) = \exp\left(\frac{t - m_{J/\psi}^2}{\Lambda_{J/\psi}^2}\right), \quad (12)$$

$$F_{J/\psi J/\psi \mathbb{P}}^H(u) = \exp\left(\frac{u - m_{J/\psi}^2}{\Lambda_{J/\psi}^2}\right). \quad (13)$$

Please note that the **form factors are normalized to 1 on the meson (J/ψ) mass shell**. One could also use monopole-like form factors. These form factors drastically reduce the $J/\psi J/\psi$ component of the amplitude in comparison to light vector meson components. However, due to compactness of J/ψ we expect $\Lambda_{J/\psi}$ to be large. In the calculations presented here, we take $\Lambda_{J/\psi} = 2 \text{ GeV}$ for illustration. In a similar fashion, one could include **one J/ψ fluctuation and one light vector meson fluctuation**. However, **there the choice of trajectories is unclear**.

Elementary cross section, VDM

Finally, let us discuss the helicity structure of the double photon hadronic fluctuation amplitude. We write:

$$\mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{(t)} = A(t) \delta_{\lambda_1 \lambda_3} \delta_{\lambda_2 \lambda_4} , \quad (14)$$

$$\mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{(u)} = A(u) \delta_{\lambda_1 \lambda_4} \delta_{\lambda_2 \lambda_3} . \quad (15)$$

$A(t)$ and $A(u)$ are given explicitly in (7). Then the total double VDM amplitude, including t and u processes, reads:

$$\mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{VDM} = \frac{1}{\sqrt{2}} \left(\mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{VDM,(t)} + \mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{VDM,(u)} \right) . \quad (16)$$

Having the double VDM helicity amplitudes, we can **add different mechanisms coherently**:

$$\mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4} = \mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{boxes} + \mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{VDM} + \mathcal{M}_{\lambda_1 \lambda_2 \rightarrow \lambda_3 \lambda_4}^{\pi^0} + \dots . \quad (17)$$

In the following, we shall discuss the coherent sum of the larger two components (boxes and VDM) and quantify their interference effects.

Background contribution

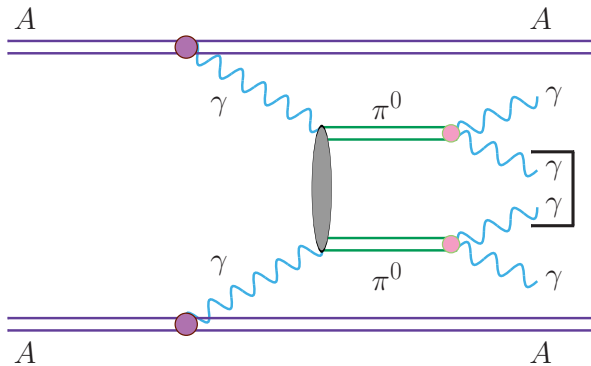


Figure: The dominant background contribution.

Background contribution

The $\gamma\gamma \rightarrow \pi^0(\rightarrow 2\gamma)\pi^0(\rightarrow 2\gamma)$ reaction constitutes a difficult background for the $\gamma\gamma \rightarrow \gamma\gamma$ measurements at intermediate $M_{\gamma\gamma}$. How to calculate the cross section for $\gamma\gamma \rightarrow \pi^0\pi^0$ reaction was discussed by Klusek-Gawenda, Szczurek 2013.

The calculation of the background proceeds in three steps.

- First the cross section for $\gamma\gamma \rightarrow \pi^0\pi^0$ is calculated.
- Next the cross section for $AA \rightarrow AA\pi^0\pi^0$ is computed in the equivalent photon approximation in an analogous way as described in the previous subsection.
- Finally the simulation of both π^0 decays is performed and combined distributions of one photon from the first π^0 and one photon from the second π^0 are constructed.

Results, elementary cross section

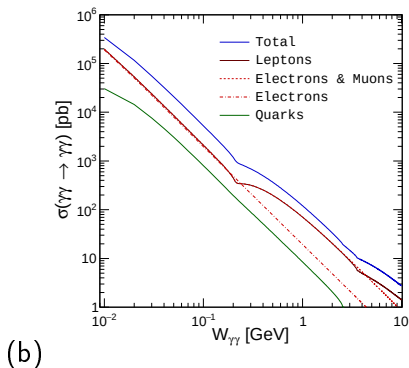
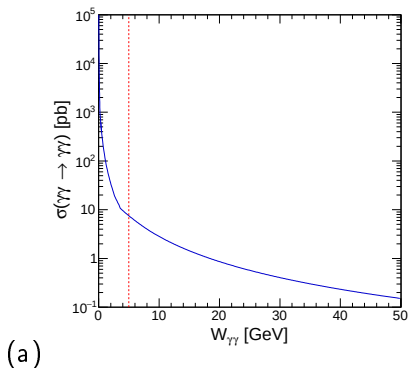


Figure: Elementary cross section (in pb) as a function of energy. Results for fermionic boxes are shown over a wide range of energies (a) MeV range (b). The total cross section (blue solid line) is split into quarks (green solid line), electrons (red dashed line), electrons and muons (red dotted line), and leptons (red solid line) contributions.

Results, elementary cross section

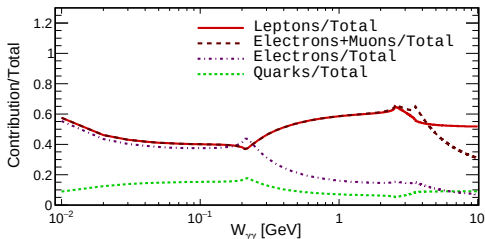
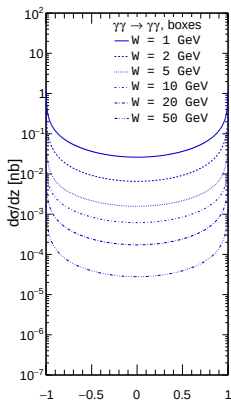


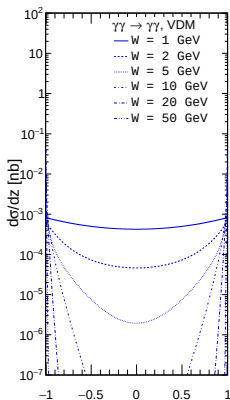
Figure: A ratio of each contribution to a coherent sum of them. The red solid line represents leptonic cross section divided by total cross section, the green dotted line refers to quarks, magenta dash-dotted line to electrons and dashed dark-red line applies to a sum of electron and muon contributions.

strong interference effects

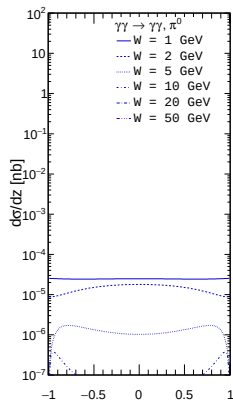
Results, elementary cross section



(a)



(b)



(c)

Figure: $\cos(\theta)$ distributions for (a) boxes, (b) double hadronic fluctuations and (c) π^0 exchange for different energies $W = 1, 2, 5, 10, 20, 50$ GeV.

Results, elementary cross section, heavier V in VDM

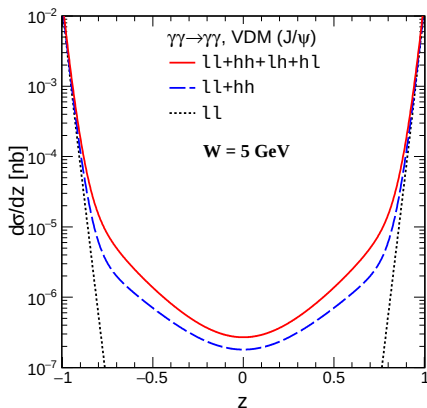
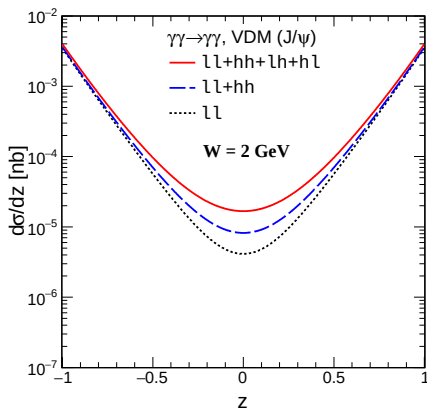


Figure: Including different photon fluctuations: $VV, J/\psi J/\psi, VJ/\psi, J/\psi V$.

The effect disappears at $z \pm 1$. Here $C_{J/\psi J/\psi} = C_{VV}/4$ and

$C_{J/\psi V} = C_{V J/\psi}/2$.

Results, elementary cross section, box+VDM interference

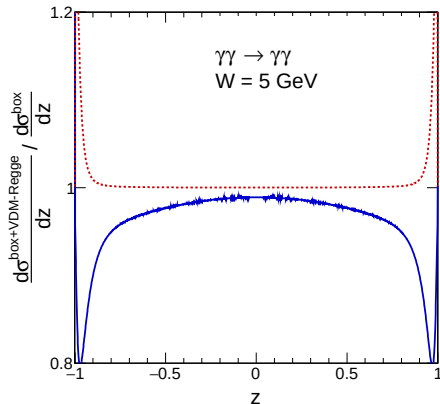


Figure: The ratio of the coherent (blue) and incoherent (red) sum of the box and VDM-Regge contributions divided by the cross section for the box contribution alone for $W = 5 \text{ GeV}$.

Results - nuclear cross section

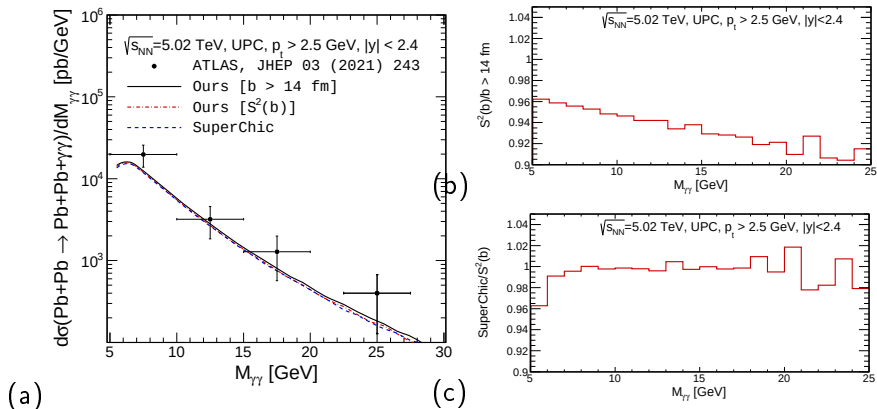


Figure: Differential cross section as a function of two-photon invariant mass at $\sqrt{s_{NN}} = 5.02$ TeV.

(a) The ATLAS experimental data are collected with theoretical results including a sharp cut on impact parameter ($b > 14$ fm - solid black line) and the nuclear modification factor $S^2(b)$ (dashed red line). (b) The ratio of the cross section for $b > 14$ fm to the cross section for $S^2(b)$ (red line). (c) The ratio of the SuperChic results to the $S^2(b)$ results (red line).

Results - nuclear cross section

Experiment	$p_{t,min}$ [GeV]	UPC condition	σ_{tot} [nb]
ATLAS	2.5	$b > 14$ fm	81.062 ± 0.05
		$S^2(b)$ Eq. (3)	77.084 ± 0.005
		SuperChick	76.421 ± 0.074
CMS	2	$b > 14$ fm	105.986 ± 0.067
		$S^2(b)$	102.104 ± 0.057
		SuperChick	100.101 ± 0.144

Table: Total cross section for $PbPb \rightarrow PbPb\gamma\gamma$ in nb obtained in different approaches for experimental ATLAS/CMS kinematics: collision energy $\sqrt{s_{NN}} = 5.02$ TeV, di-photon invariant mass $M_{\gamma\gamma} > 5$ GeV, photon rapidity $|y| < 2.4$. ATLAS and CMS detected photons in different range of transverse momenta.

Results - nuclear cross section

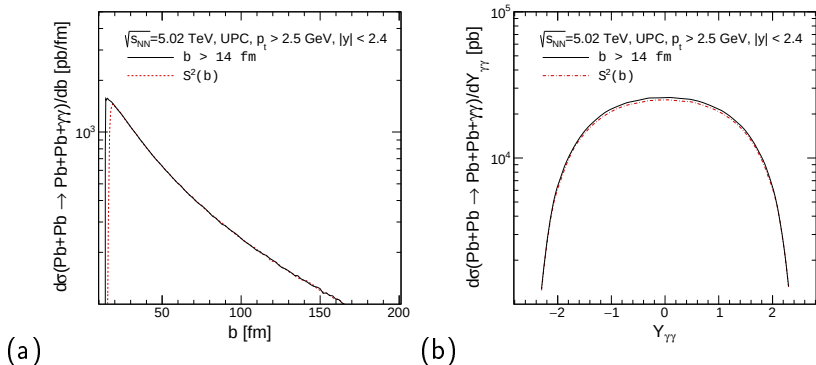
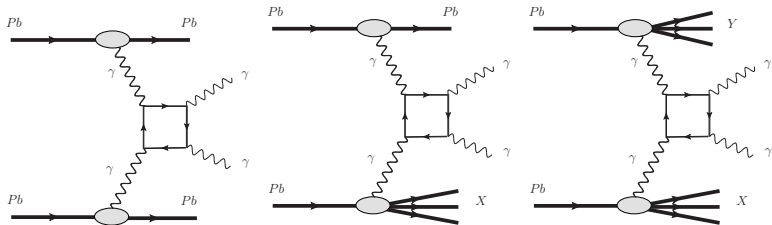


Figure: (a) Impact parameter and (b) rapidity distribution for UPC of lead-lead at collision energy $\sqrt{s_{NN}} = 5.02$ TeV. The solid line corresponds to a sharp cut on the impact parameter and the red dash-dotted line is for the absorption factor given by Eq. (3).

Something missing?



Inelastic contributions

The inclusion of the inelastic component in the calculation of LbL scattering in $PbPb$ collisions, implies that the cross-section will be given schematically by

$$\begin{aligned}\sigma^{LbL}(\sqrt{s_{NN}}) &\propto f_{\gamma/Pb}^{el}(x_1) \otimes f_{\gamma/Pb}^{el}(x_2) \otimes \hat{\sigma}[\gamma\gamma \rightarrow \gamma\gamma; W_{\gamma\gamma}] \\ &+ f_{\gamma/Pb}^{el}(x_1) \otimes f_{\gamma/Pb}^{inel}(x_2) \otimes \hat{\sigma}[\gamma\gamma \rightarrow \gamma\gamma; W_{\gamma\gamma}] + \\ &+ f_{\gamma/Pb}^{inel}(x_1) \otimes f_{\gamma/Pb}^{el}(x_2) \otimes \hat{\sigma}[\gamma\gamma \rightarrow \gamma\gamma; W_{\gamma\gamma}] + \\ &+ f_{\gamma/Pb}^{inel}(x_1) \otimes f_{\gamma/Pb}^{inel}(x_2) \otimes \hat{\sigma}[\gamma\gamma \rightarrow \gamma\gamma; W_{\gamma\gamma}] , \quad (18)\end{aligned}$$

$$f_{\gamma/Pb}^{el}(x) = \frac{\alpha Z^2}{\pi x} \{2\xi K_0(\xi)K_1(\xi) - \xi^2[K_1^2(\xi) - K_0^2(\xi)]\} , \quad (19)$$

and

$$f_{\gamma/Pb}^{inel}(x, \mu^2) = Z \times f_{\gamma/p}(x, \mu^2) + (A - Z) \times f_{\gamma/n}(x, \mu^2) , \quad (20)$$

Photon distributions in nucleons

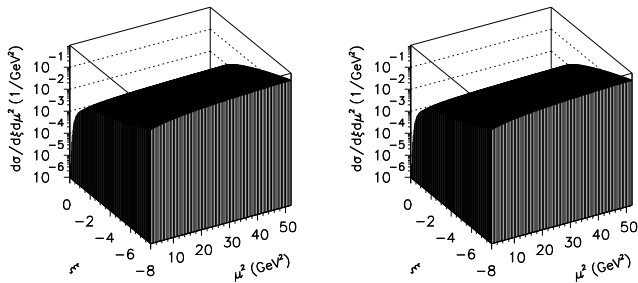


Figure: $xf_{\gamma/p}(x, \mu^2)$ (left) and $xf_{\gamma/n}(x, \mu^2)$ (right)

There are two components in $f_{\gamma/p}$ and $f_{\gamma/n}$.
Here I will use **CTEQ18QED fluxes**

Inelastic photon distribution in Pb

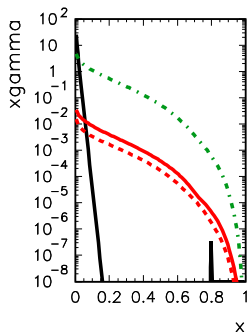


Figure: $xf_{\gamma/Pb}(x, \mu = 5\text{GeV})$ elastic and inelastic contributions.

Inelastic contribution:

$$xf_{\gamma/A}^{ine}(x) = Zxf_{\gamma/p}(x) + Nxf_{\gamma/n}(x) . \quad (21)$$

Modification of results due to the inelastic contributions

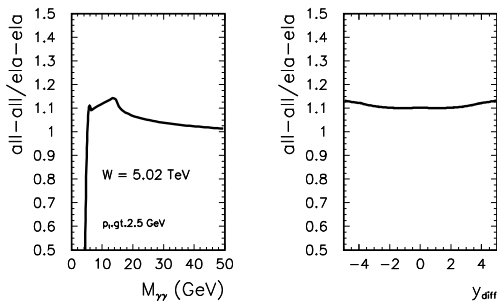
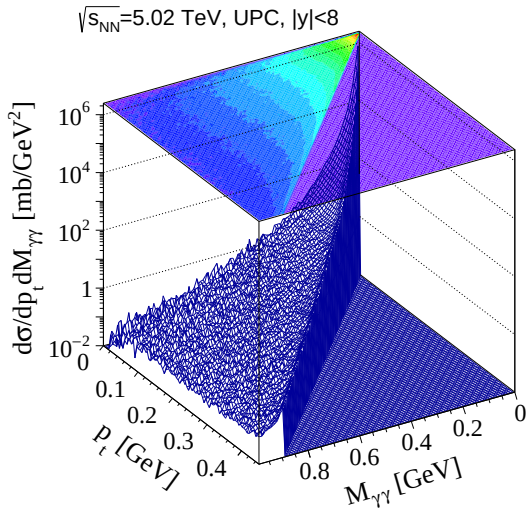


Figure: The ratio of the modified cross section to the coherent-coherent contribution.

END OF NEW RESULTS

Let us return to the main line

Broad range of rapidity, full phase space



Broad range of rapidity, full phase space

	σ_{tot} [mb]
Total	91.675 ± 0.023
electrons+muons	41.597 ± 0.010
electrons	39.163 ± 0.010
quarks	12.483 ± 0.003

Table: Total cross section in mb for $\text{PbPb} \rightarrow \text{PbPb} \gamma\gamma$ for different fermionic contributions artificially separated. Here collision energy $\sqrt{s_{NN}} = 5.02$ TeV, di-photon invariant mass $M_{\gamma\gamma} \in (0.01 - 1)$ GeV, photon transverse momentum $p_t > 5$ MeV and photon rapidity $|y| < 8$.

Broad range of rapidity, full phase space

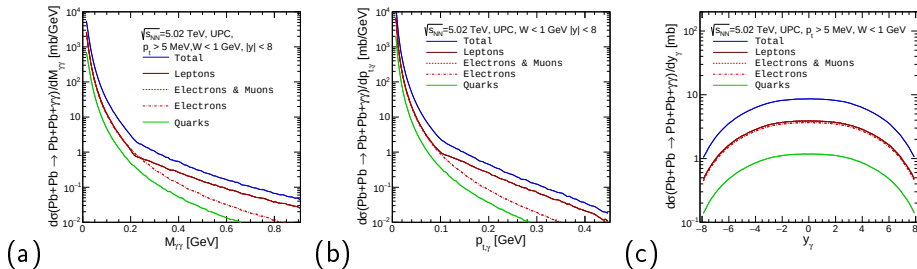


Figure: (a) Di-photon invariant mass, (b) photon transverse momentum and (c) rapidity distribution in mb for different fermionic contributions in ultraperipheral lead-lead collisions at energy $\sqrt{s_{NN}} = 5.02$ TeV. The blue solid line corresponds to a sum of all contributions, the red solid line is a sum for leptons, the dashed red line relates to electrons, and the green solid line represents quarks contribution.

Broad range of rapidity, full phase space

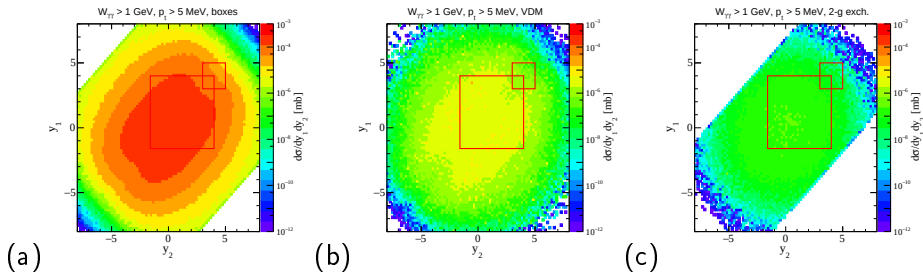


Figure: Distribution in (y_1, y_2) in mb for photon transverse momentum $p_t > 5$ MeV, di-photon invariant mass $M_{\gamma\gamma} < 1$ GeV
(a) Boxes, (b) VDM-Regge and (c) two-gluon exchange.

Broad range of rapidity, full phase space

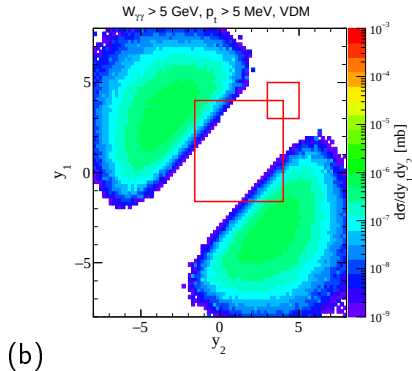
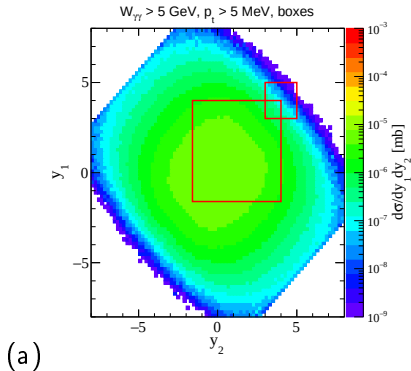


Figure: Distribution in (y_1, y_2) in mb for transverse momentum $p_t > 5 \text{ MeV}$, di-photon invariant mass $M_{\gamma\gamma} > 5 \text{ GeV}$.

(a) Boxes, (b) VDM-Regge mechanism.

Broad range of rapidity, full phase space

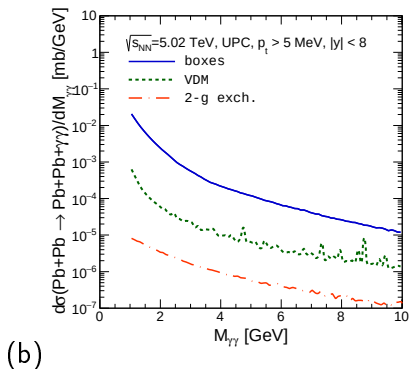
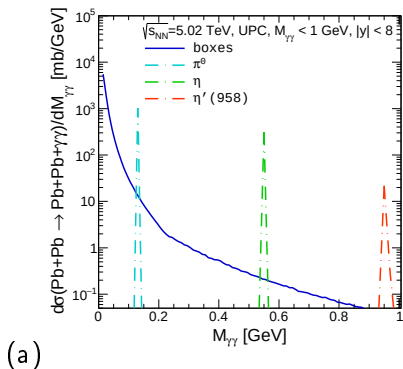


Figure: Di-photon invariant mass distribution for photon transverse momentum $p_t > 5$ MeV and photon rapidities $y_{1/2} \in (-8, 8)$.
 (a) Boxes vs resonances, $M_{\gamma\gamma} < 1$ GeV (b) Boxes vs VDM-Regge and vs two-gluon exchange contribution, $M_{\gamma\gamma} > 1$ GeV.

Broad range of rapidity, full phase space

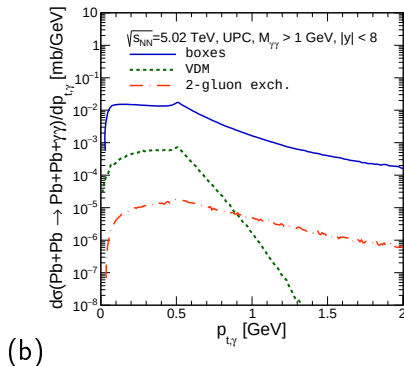
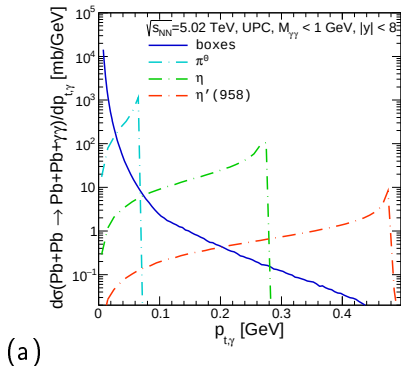


Figure: Transverse momentum distribution for photon transverse momentum $p_t > 5$ MeV, di-photon invariant mass $M_{\gamma\gamma} < 1$ GeV (a) or $M_{\gamma\gamma} > 1$ GeV (b) and photon rapidities $y_{1/2} \in (-8, 8)$.

(a) Boxes vs resonances; (b) boxes vs VDM-Regge vs two-gluon exchange.

Broad range of rapidity, full phase space

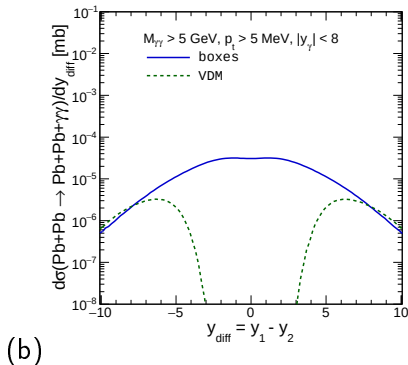
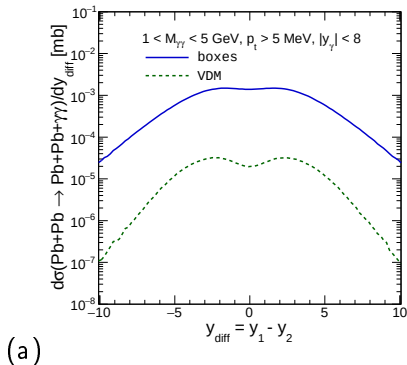


Figure: Distribution in y_{diff} for light-by-light scattering in $PbPb \rightarrow PbPb\gamma\gamma$. Here the transverse momentum cut is equal to **5 MeV**. The blue solid line relates to the boxes, and the green dotted line to the VDM-Regge contribution. Here the range of measured di-photon invariant mass is: (a) **(1 – 5) GeV**, (b) **> 5 GeV**.

ALICE and FoCal

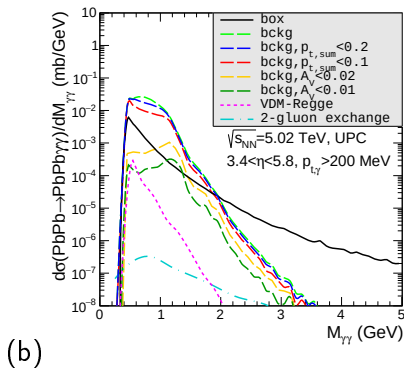
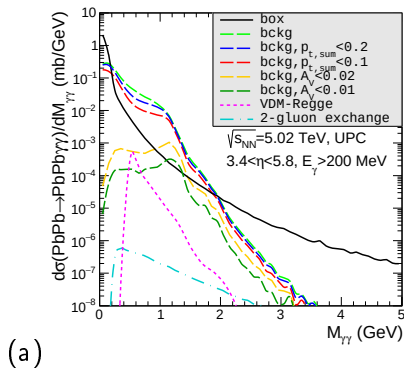


Figure: Invariant mass distribution for the nuclear process for the future FoCal:

(a) $E_{t,\gamma} > 200$ MeV and $3.4 < y_{\gamma_{1/2}} < 5.8$,

(b) $p_{t,\gamma} > 200$ MeV and $3.4 < y_{\gamma_{1/2}} < 5.8$.

Here, both photons are "measured" in FoCal. The background contribution is presented for different elimination cuts

ALICE and FoCal

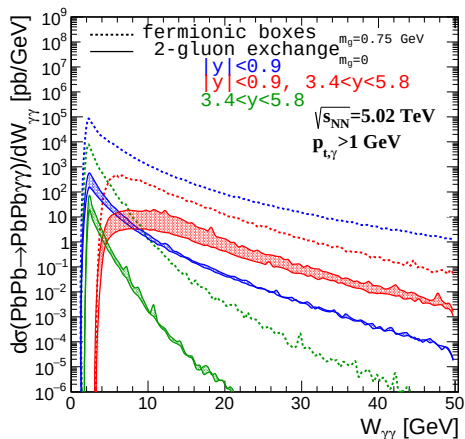


Figure: Di-photon invariant mass distribution for $\text{PbPb} \rightarrow \text{PbPb}\gamma\gamma$ process for ALICE, FoCal and their combination.

ALICE and FoCal

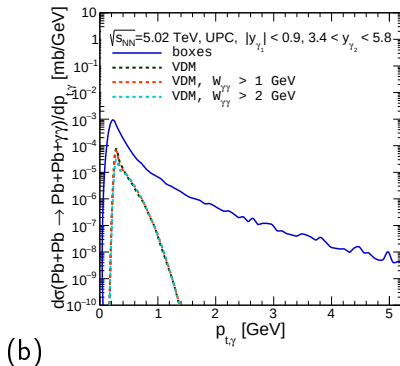
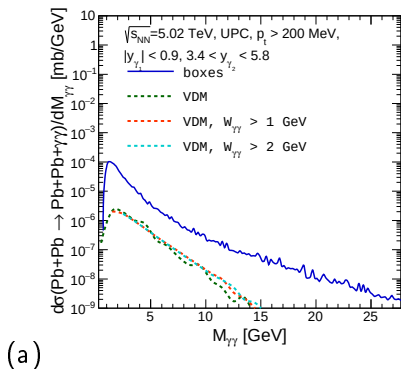


Figure: Prediction for the FoCal detector in association with mid-rapidity ALICE detector for photons: $p_t > 200$ MeV, di-photon mass $M_{\gamma\gamma} > 400$ MeV and photon rapidities $|y_1| < 0.9$ and $y_2 \in (3.4, 5.8)$. The blue line corresponds to fermionic loops and the green lines to the VDM-Regge contribution.

ALICE and FoCal

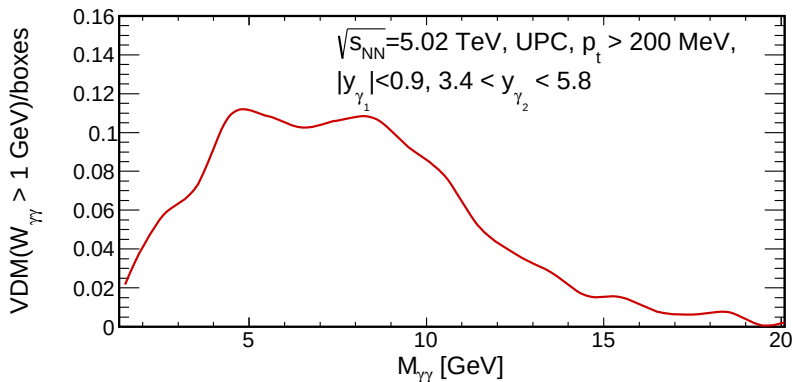


Figure: Ratio of the VDM-Regge ($W_{\gamma\gamma} > 1$ GeV) and box contributions as a function of di-photon invariant mass. **No interference effects are included here.**

ALICE 3 kinematics

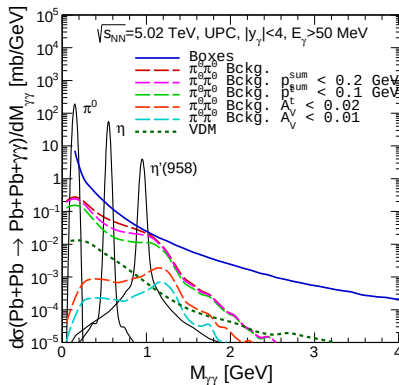


Figure: Di-photon invariant mass distribution for ALICE 3, i.e. rapidity $y_\gamma \in (-4, 4)$ and photon energy $E_\gamma > 50 \text{ MeV}$.

Here the blue solid line relates to the box contribution, the dotted line to the VDM-Regge component and the dashed lines are for double- π^0 background

ALICE 3 kinematics

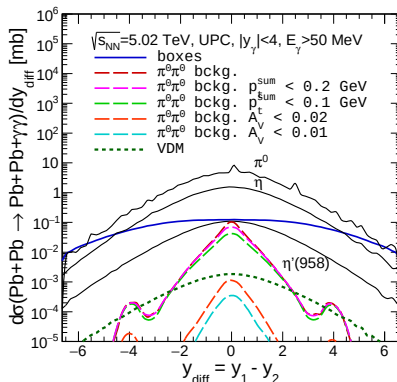


Figure: Differential cross section as a function of $y_{diff} = y_1 - y_2$ for extended ALICE 3 kinematics: $|y_\gamma| < 4$ and $E_\gamma > 50$ MeV. Results are presented for boxes, resonances, VDM-Regge and double- π^0 background.

ALICE 3 kinematics

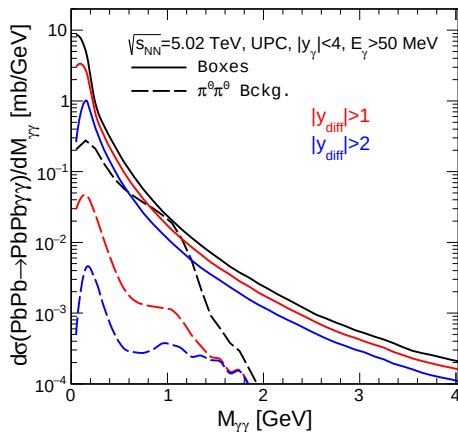


Figure: Influence of the $|y_{\text{diff}}|$ cut on the invariant $M_{\gamma\gamma}$ distribution.

ALICE 3 kinematics

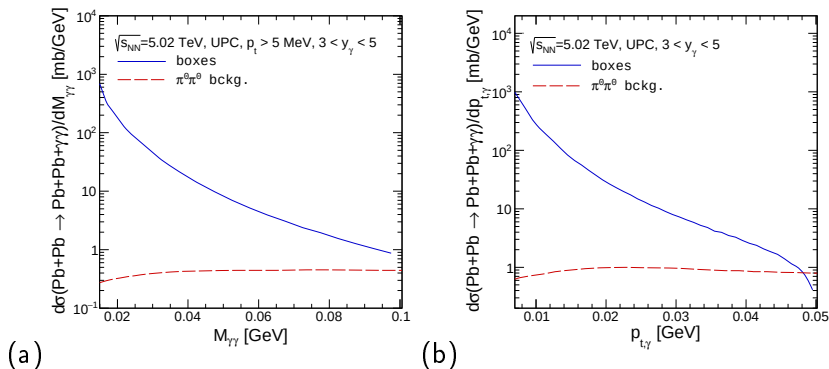
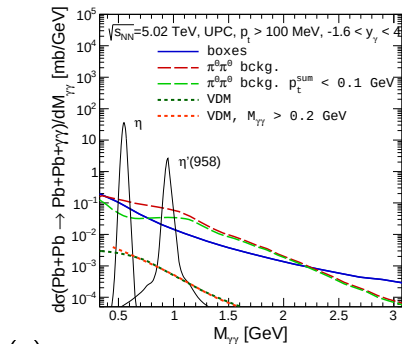
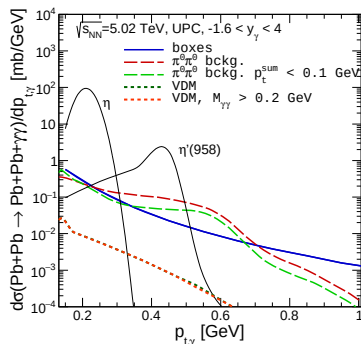


Figure: Prediction for the ALICE 3 experiment for soft photons: $p_t \in (5 - 50)$ MeV and photon rapidities $y_i \in (3, 5)$. The blue line corresponds to fermionic loops, the red line relates to the double- π^0 background.

ALICE 3 kinematics



(a)



(b)

Figure: Prediction for the ALICE 3 experiment for photons: $p_t > 100$ MeV and photon rapidities $y_i \in (-1.6, 4)$. The blue line corresponds to fermionic loops, the red line relates to the double- π^0 background.

ALICE 3 kinematics

kinematical limitation			$\sigma(\pi^0\pi^0 \text{ bckg}) \text{ } \mu b$	$\sigma(\text{boxes}) \text{ } \mu b$
$p_t > 5 \text{ MeV}$	$E_\gamma > 50 \text{ MeV}$	$-4 < \eta_\gamma < 4$	112.824	952.590
$p_t > 100 \text{ MeV}$	$E_\gamma > 50 \text{ MeV}$	$-1.6 < \eta_\gamma < 4$	159.231	82.682
$p_t = (1 - 50) \text{ MeV}$	$E_\gamma > 50 \text{ MeV}$	$3 < \eta_\gamma < 5$	105.301	3095.795

Table: Fermionic box signal contribution versus double- π^0 background given in μb at PbPb collision energy $\sqrt{s_{NN}} = 5.02 \text{ TeV}$, with three considered scenarios of ALICE 3 detector limitations.

Conclusions

- We have investigated how to extend future studies of light-by-light scattering using future infrastructure (FoCal and ALICE 3).
- Several mechanisms of $\gamma\gamma \rightarrow \gamma\gamma$ scattering have been discussed.
- We have discussed shortly double photon fluctuation into vector mesons and Regge-type of interaction (VDM-Regge) as well as its interference with boxes. Negative interference was found.
- Combination of ALICE and FoCal can extend the range of $\gamma\gamma \rightarrow \gamma\gamma$ scattering to smaller energies.
- The special photon detector at $y \in (3,5)$ would allow to study the LbL scattering at very low energies and/or transverse momenta.
- The ways to eliminate unwanted double- π^0 background were discussed.
(condition on $p_{t,sum}$, vector asymmetry, y_{diff})
- We have found that the background can be drastically reduced at large $|y_{diff}|$ which automatically means larger energies. This condition could be easily imposed using ALICE 3 detector.