

Heavy Flavor Kinetics in HIC

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Motivation

☞ Charm quarks probe the properties of the QGP

Task: study the $c\bar{c}$ production in equilibrated QGP, $N_f = 2 + 1$, two evolution scenarios:

- 1D ideal fluid (Bjorken flow)

-(2+1)D viscous QGP + $(\eta/s)(T)$

☞ Quasiparticle model

☞ Rate equation: cross sections, charm production rate

☞ Number of $c\bar{c}$ pairs in hot deconfined matter

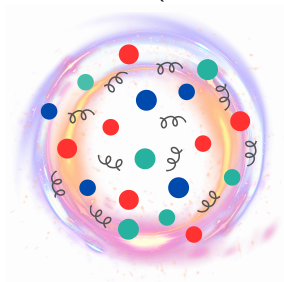
Quasiparticle Model - Effective Approach to QCD

☞ similar to massive quasielectron moving freely in solid states

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Real QGP:

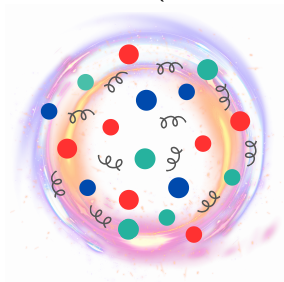


strongly-interacting particles,
constant (bare) masses m_j

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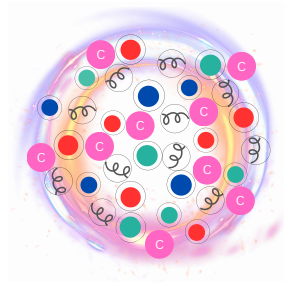
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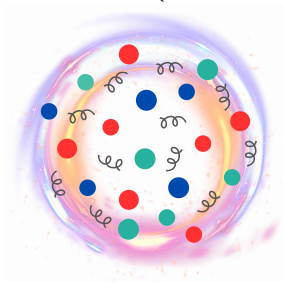


weakly-interacting **quasi**particles,
dynamical $m_i^{eff} [T, G(T)]$

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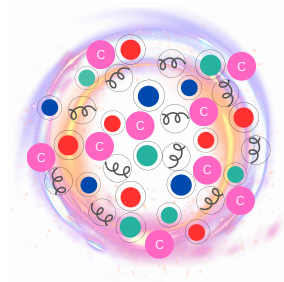
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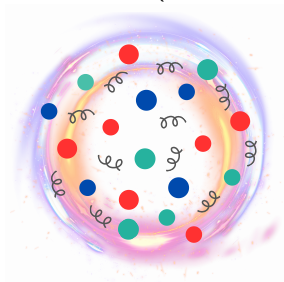
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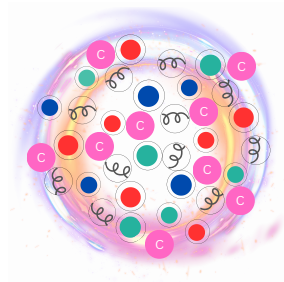
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$G(T)$ from lattice QCD EoS

Quasiparticle Model

Quasiparticles are „dressed” with effective masses $m_i[G(T), T]$:

$$m_i[G(T), T] = \sqrt{(m_i^0)^2 + \Pi_i[G(T), T]} \quad (1)$$

self-energies Π_i from pQCD (Hard Thermal Loops):

$$\text{gluons: } \Pi_g[G(T), T] = \left(3 + \frac{N_f}{2}\right) \frac{G^2(T)}{6} T^2 \quad (2)$$

$$\text{quarks: } \Pi_{l,s}[G(T), T] = 2 \left[m_{l,s}^0 \sqrt{\frac{G^2(T) T^2}{6}} + \frac{G^2(T) T^2}{6} \right] \quad (3)$$

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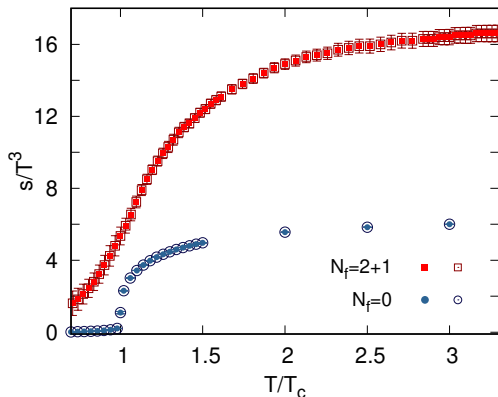
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➡ effective coupling $G(T)$ – reliable thermodynamics – lattice QCD

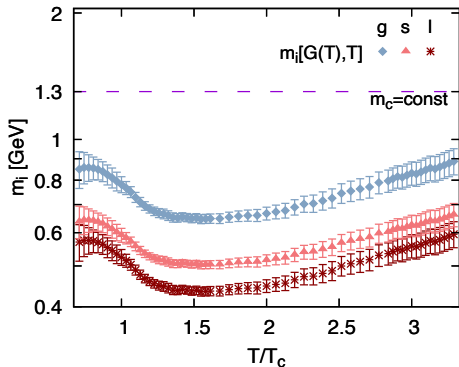
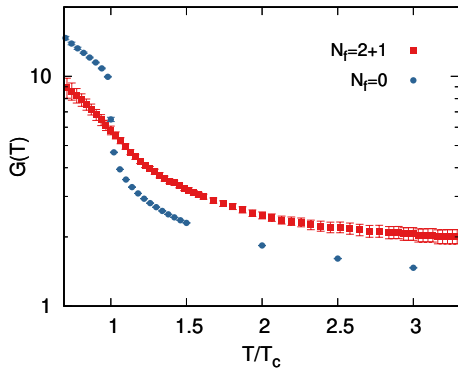
Quasiparticle Model

$$s(T) \simeq \sum_{i=g,l,s,\dots} \int d^3p ([1 \pm f_i^0] \ln[1 \pm f_i^0] \mp f_i^0 \ln f_i^0) = \text{lattice data} \rightarrow G(T)$$

$$f_i^0(E_i) : E_i[G(T), T] = \sqrt{p^2 + m_i^2[G(T), T]} \quad (4)$$



Effective Coupling and Masses



$$m_i[G(T), T] \gg m_i^0 = 5 \text{ MeV}, m_s^0 = 95 \text{ MeV}$$

[V.M., M. Bluhm, K. Redlich, C. Sasaki, PRD100 '19]

Charm Quark Evolution

Rate equation [Biro et al., PRC 48 '93; Zhang et al., PRC 77 '08]:

describes time/temperature evolution of the number density function

$$\partial_\mu (n_c u^\mu) = R_{l\bar{l} \rightarrow c\bar{c}} + R_{s\bar{s} \rightarrow c\bar{c}} + R_{gg \rightarrow c\bar{c}} - R_{c\bar{c} \rightarrow l\bar{l}} - R_{c\bar{c} \rightarrow s\bar{s}} - R_{c\bar{c} \rightarrow gg} \quad (5)$$

Applying the detailed balance:

$$\partial_\mu (n_c u^\mu) = [\bar{\sigma}_{l\bar{l} \rightarrow c\bar{c}} (n_l^0)^2 + \bar{\sigma}_{s\bar{s} \rightarrow c\bar{c}} (n_s^0)^2 + \frac{1}{2} \bar{\sigma}_{gg \rightarrow c\bar{c}} (n_g^0)^2] \left(1 - \frac{n_c^2}{(n_c^0)^2}\right) \quad (6)$$

$$n_i^0 = d_i \int d^3 p f_i^0 [E_i(T)] \quad (7)$$

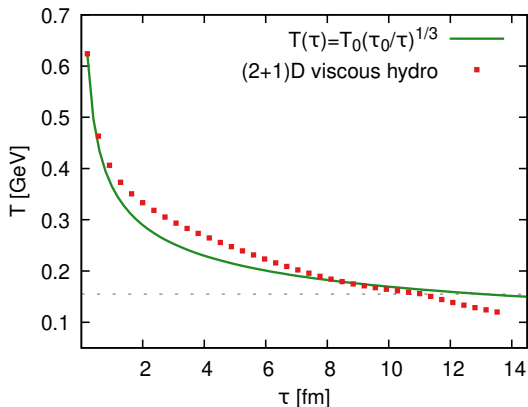
$$n_c = \lambda_c(\tau) n_c^0 \quad (8)$$

* LHS depends on the QGP evolution:

- i) 1D Bjorken flow, ideal fluid;
- ii) (2+1)D expansion $+(\eta/s)(T)$

[V.M., M. Bluhm, K. Redlich, C. Sasaki, PRD100 '19; Auvinen, Eskola, Huovinen, Niemi, Paatelainen, Petreczky, PRC 102 '20]

QGP Evolution



$$T_0 = 0.624 \text{ GeV}, \tau_0 = 0.2 \text{ fm}$$

[Auvinen, Eskola, Huovinen, Niemi, Paatelainen, Petreczky, PRC 102 '20]

Thermal-Averaged Cross Sections

$$\sigma(\sqrt{s}) \rightarrow \bar{\sigma} = \langle \sigma v \rangle \quad (9)$$

For Boltzmann statistics:

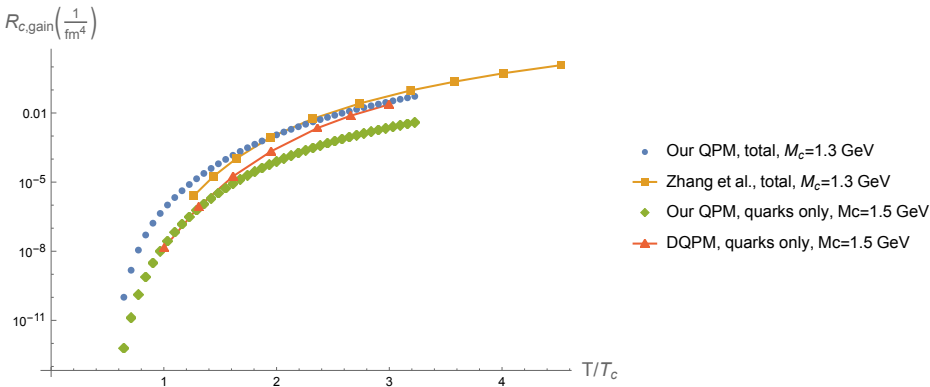
$$\bar{\sigma}_{ab \rightarrow cd} = \frac{\int d^3 p_a d^3 p_b f_a f_b \sigma_{ab \rightarrow cd} v_{ab}}{\int d^3 p_a d^3 p_b f_a f_b} = \quad (10)$$

$$\left[4 \frac{M_a^2}{T^2} \frac{M_b^2}{T^2} K_2\left(\frac{M_a}{T}\right) K_2\left(\frac{M_b}{T}\right) \right]^{-1} \times$$

$$\int_{\sqrt{s_0}}^{\infty} d(\sqrt{s}) K_1\left(\frac{\sqrt{s}}{T}\right) \sigma_{ab \rightarrow cd} \left[\frac{s}{T^2} - \left(\frac{M_a^2}{T^2} + \frac{M_b^2}{T^2} \right)^2 \right] \left[\frac{s}{T^2} - \left(\frac{M_c^2}{T^2} + \frac{M_d^2}{T^2} \right)^2 \right];$$

$$\sqrt{s_0} = \max[M_a + M_b, M_c + M_d]$$

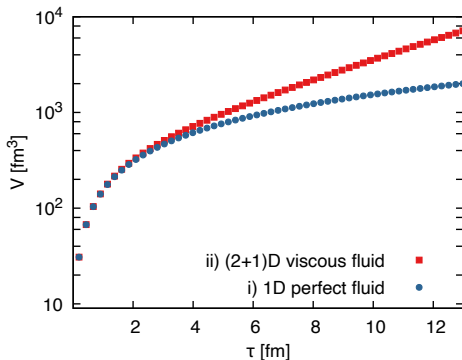
Charm Quark Production Rate



DQPM: T. Song, I. Grishmanovskii, O. Soloveva, E. Bratkovskaya, arXiv:2404.00425 (2024);
Zhang et al., Phys. Rev. C 77 (2008)

Volume of the QGP

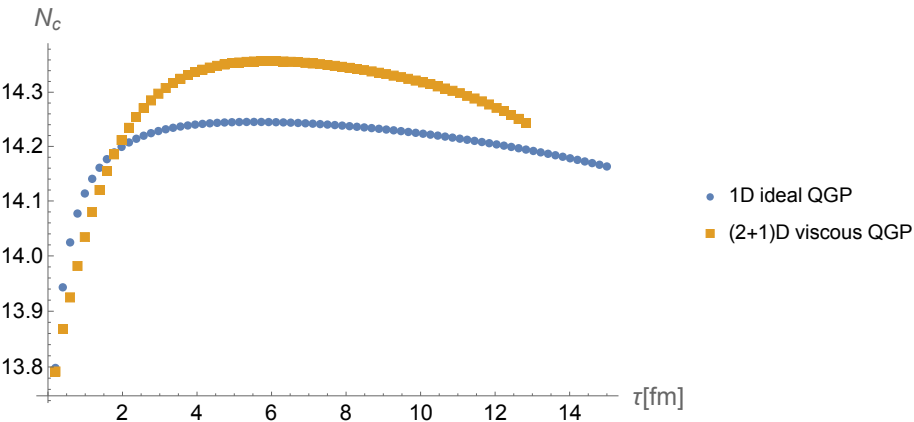
$$V(\tau) = \pi R^2 \tau \quad (11)$$



i) 1D ideal Bjorken dynamics: $R = R_0 = 7 \text{ fm}$

ii) (2+1)D viscous expansion: $R(\tau) = R_0 + (\tau - \tau_0)^2 a/2$, $a = 0.01 \text{ fm}$

Charm Quark Production



Initial charm quark number: $\frac{dN_c}{dy} = 13.8$ (12)

Summary

- ☞ **Quasiparticle model** – effective well-established tool connecting non-perturbative and perturbative QCD regimes.
- ☞ **Charm quarks** – minor thermal production in both ideal 1D- and viscous (2+1)D-expanding plasma.
- ☞ **Possibilities** – quasihadrons out of chemical equilibrium, finite μ ,
 $N_f = 2 + 1 + 1\dots$