







# Probing the QCD phase transition within various astrophysical scenarios

Noshad Khosravi Largani Brainstorming workshop





Weber, F. & Negreiros, Rodrigo & Rosenfield, Philip. (2007)











Khosravi Largani et al., A&A 687, A245 (2024)



Khosravi Largani et al., A&A 687, A245 (2024)

 $M_{\rm NS} = 1.55~M_{\odot}$ 



Khosravi Largani et al., A&A 687, A245 (2024)

Weak process 1  $e^- + p \rightleftharpoons n + \nu_e$  $e^+ + n \rightleftharpoons p + \bar{\nu}_e$ 2 3  $n \rightleftharpoons p + e^- + \bar{\nu}_e$  $\nu_e + (A, Z - 1) \rightleftharpoons (A, Z) + e^-$ 4  $\nu + N \rightleftharpoons \nu' + N$ 5  $\nu + (A, Z) \rightleftharpoons \nu' + (A, Z)$ 6  $\nu + e^{\pm} \rightleftharpoons \nu' + e^{\pm}$ 7  $e^- + e^+ \rightleftharpoons \nu + \bar{\nu}$ 8 9  $N + N \rightleftharpoons \nu + \bar{\nu} + N + N$  $\nu_e + \bar{\nu}_e \rightleftharpoons \nu_{\mu/\tau} + \bar{\nu}_{\mu/\tau}$ 10  $\nu = \{\nu_e, \bar{\nu}_e, \nu_{\mu/\tau}, \bar{\nu}_{\mu/\tau}\} \text{ and } N = \{n, p\}$ 



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$$Q_{ij} = \int \rho(r)(3r_ir_j - |\vec{r}|^2\delta_{ij})d^3r$$
$$h_{ij} = \frac{2G}{c^4r}\frac{d^2}{dt^2}Q_{ij}$$

$$\partial_r \eta_r = A \eta_r + B \eta_\perp,$$

$$\partial_r \eta_\perp = C \eta_r + D \eta_\perp.$$

g-modes:

$$\partial_r \eta_r + \left[\frac{2}{r} + 6\frac{\partial_r \psi}{\psi}\right] \eta_r - \frac{l(l+1)}{r^2} \eta_\perp = 0,$$
  
 $\partial_r \eta_\perp - \left(1 - \frac{\mathcal{N}^2}{\sigma^2}\right) \eta_r + \left[\partial_r \ln q - G\right] \eta_\perp = 0.$ 

Torres-Forne et al. (2019), Mon. Not. Roy. Astron. Soc. Torres-Forne et al. (2018), Mon. Not. Roy. Astron. Soc.

$$\begin{split} \rho &\to \rho + \delta \rho \\ \delta \rho &= \delta \rho + \xi^{i} \partial_{i} \rho \\ \partial_{t} \xi^{i} &= \delta v^{*i} \\ \delta P &= \delta \hat{P} Y_{lm} e^{-i\sigma t}, \\ \xi^{r} &= \eta_{r} Y_{lm} e^{-i\sigma t}, \\ \xi^{\theta} &= \eta_{\perp} \frac{1}{r^{2}} \partial_{\theta} Y_{lm} e^{-i\sigma t}, \\ \xi^{\varphi} &= \eta_{\perp} \frac{1}{r^{2} \sin^{2} \theta} \partial_{\varphi} Y_{lm} e^{-i\sigma t}, \end{split}$$

## Types of fluid modes (non-radial oscillations of non-rotating stars)

#### p-modes

- Pressure modes ( $l \ge 0$ )
- Restoring force = pressure
- Standing sound waves
- Lamb frequency (*L*):

$$L^2 \approx c_s^2 \frac{l(l+1)}{r^2}$$

- Mode frequency

$$f \propto c_s \propto \sqrt{\overline{\rho}} \propto \sqrt{M/R^3}$$

#### f-mode

- Fundamental mode ( $l \ge 2$ )
- Node-less mode (simple case)
- Lowest-order p-mode?

$$L^2 \approx c_s^2 \frac{l(l+1)}{r^2}$$

Mode frequency

$$f \propto c_s \propto \sqrt{\overline{\rho}} \propto \sqrt{M/R^3}$$

#### g-modes

- Gravity modes ( $l \ge 1$ )
- Restoring force = buoyancy
- Brunt-Väisälä frequency (*N*)

$$N^{2} \approx \frac{\partial \Phi}{\partial r} \frac{1}{\rho} \left( \frac{1}{c_{s}^{2}} \frac{\partial P}{\partial r} - \frac{\partial \rho}{\partial r} \right)$$

Mode frequency

$$f \propto \frac{M}{R^2} \times \sqrt{\frac{(\Gamma - 1)m_n}{\Gamma k_b T}}$$

 $M_{\rm NS} = 1.55~M_{\odot}$ 





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Credit: NASA images from Hubble telescope









## Black-hole formation: Two distinct scenarios



Jakobus et al., MNRAS, 516, 2554 (2022) Khosravi Largani et al., ApJ 946, 143 (2024)

### v – signal @ Super-Kamiokande (d ~ 10 kpc)



Fischer et al., Nature Astron. 2, 980 (2018)



Kuroda et al., ApJ 924, 38 (2022)



Kuroda et al., ApJ 924, 38 (2022)





Khosravi Largani et al., ApJ 946, 143 (2024)

Progenitor	EOS	$t_{\mathrm{burst}}$	$L_{\overline{\nu}_e, \mathrm{peak}}$	$\langle E_{\overline{\nu}_e} \rangle$	$E_{\mathrm{expl}}$
	RDF	$[\mathbf{S}]$	$[10^{53} \text{ erg s}^{-1}]$	[MeV]	$[10^{51} \text{ erg}]$
s25a28	1.9	0.345	6.36	38.59	4.21
s30a28	1.2	1.056	4.80	56.21	1.93
s30a28	1.8	0.833	5.64	42.21	2.66
s30a28	1.9	0.580	8.30	43.49	3.28
s40a28	1.2	0.895	4.15	38.60	1.59
s40a28	1.8	0.717	2.06	35.77	1.23
s40a28	1.9	0.491	4.28	39.94	3.31
s40.0	1.8	0.694	5.61	43.03	2.32
s40.0	1.9	0.443	8.52	48.69	3.79
u50	1.1	1.227	3.90	26.55	2.3
u50	1.2	0.819	5.37	36.19	3.8
s75.0	1.2	1.803	3.06	34.35	1.0

















### Thank you